Hi!

Simplifying life with simplified models

Suchita Kulkarni (LPSC, Grenoble)



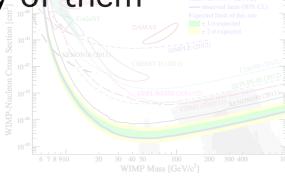
based on: arXiv: 1312.4175 S. Kraml, U. Laa, A. Lessa, W. Magerl, D. Proschofsky, W. Waltenberger and arXiv: 1308.3735 G. Belanger, G. Drieu La Rochelle, B. Dumont, R. Godbole, S. Kraml

The scene

Hunt for BSM physics is strong from the smallest to largest scales

Many new and interesting results from astrophysics and collider searches exist and they must be taken into account to test a BSM theory

Many BSM theories and no conclusive evidence for any of them





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ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

ATLA	S Preliminary
$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$	$\sqrt{s} = 7, 8 \text{ TeV}$

$\int \mathcal{L} dt =$	(4.6 - 22.9) fb ⁻¹	$\sqrt{s} =$
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010	201 0001 2010							$\int \mathcal{L} dt = (4.6 - 22.9) 10^{-2}$	$\gamma s = 7, o \text{ rev}$
	Model	e, μ, τ, γ	Jets	E ^{miss} T	∫£ dt[fb	p ⁻¹]	Mass limit		Reference
Inclusive Searches	$\begin{array}{l} MSUGRA/CMSSM \\ MSUGRA/CMSSM \\ MSUGRA/CMSSM \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \bar{q} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{\chi}_{1}^{1} \rightarrow q q W^{\pm} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\ell/(\ell v / v v) \tilde{\chi}_{1}^{0} \\ GMSB (\tilde{\ell} \; NLSP) \\ GMSB (\tilde{\ell} \; NLSP) \\ GGM (bino \; NLSP) \\ GGM (mino \; NLSP) \\ GGM (higgsino \; bino \; NLSP) \\ GGM (higgsino \; NLSP) \\ Gravitino \; LSP \end{array}$	$\begin{array}{c} 0 \\ 1 \ e, \mu \\ 0 \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 1 - 2 \ \tau \\ 2 \ \gamma \\ 1 \ e, \mu + \gamma \\ \gamma \\ 2 \ e, \mu \left(Z \right) \\ 0 \end{array}$	2-6 jets 3-6 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 0-2 jets 1 b 0-3 jets mono-jet	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 4.7 20.7 4.8 4.8 4.8 4.8 5.8 10.5	q. g g	1.7 TeV 1.2 TeV 1.2 TeV 1.1 TeV 740 GeV 1.3 TeV 1.3 TeV 1.18 TeV 1.12 TeV 1.12 TeV 1.24 TeV 1.24 TeV 1.4 TeV 1.07 TeV 619 GeV 900 GeV 690 GeV	$\begin{array}{l} m(\tilde{q}) = m(\tilde{g}) \\ any \ m(\tilde{q}) \\ any \ m(\tilde{q}) \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) < 200 \ \text{GeV}, \ m(\tilde{k}^{\pm}) = 0.5(m(\tilde{k}_{1}^{0}) + m(\tilde{g})) \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ tan\beta < 15 \\ tan\beta > 18 \\ m(\tilde{k}_{1}^{0}) > 50 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) > 50 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) > 220 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) > 220 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) > 10^{-4} \ \text{eV} \end{array}$	ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 1308.1841 ATLAS-CONF-2013-047 ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 ATLAS-CONF-2013-089 1208.4688 ATLAS-CONF-2013-026 1209.0753 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-147
3 rd gen. ἒ med.	$\widetilde{g} \rightarrow b \overline{b} \widetilde{\chi}_{1}^{0}$ $\widetilde{g} \rightarrow t \overline{t} \widetilde{\chi}_{1}^{0}$ $\widetilde{g} \rightarrow t \overline{t} \widetilde{\chi}_{1}^{0}$ $\widetilde{g} \rightarrow b \overline{t} \widetilde{\chi}_{1}^{+}$	0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 b 7-10 jets 3 b 3 b	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	90 700 700	1.2 TeV 1.1 TeV 1.34 TeV 1.3 TeV	$m(\tilde{k}_{1}^{0}) < 600 \text{ GeV}$ $m(\tilde{k}_{1}^{0}) < 350 \text{ GeV}$ $m(\tilde{k}_{1}^{0}) < 400 \text{ GeV}$ $m(\tilde{k}_{1}^{0}) < 300 \text{ GeV}$	ATLAS-CONF-2013-061 1308.1841 ATLAS-CONF-2013-061 ATLAS-CONF-2013-061
3rd gen. squarks direct production	$ \begin{array}{l} \tilde{b}_{1}\tilde{b}_{1}, \ \tilde{b}_{1} \rightarrow b \tilde{\chi}_{1}^{0} \\ \tilde{b}_{1}\tilde{b}_{1}, \ \tilde{b}_{1} \rightarrow t \tilde{\chi}_{1}^{\pm} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{light}), \ \tilde{t}_{1} \rightarrow b \tilde{\chi}_{1}^{\pm} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{light}), \ \tilde{t}_{1} \rightarrow W b \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{medium}), \ \tilde{t}_{1} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{medium}), \ \tilde{t}_{1} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{heavy}), \ \tilde{t}_{1} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{heavy}), \ \tilde{t}_{1} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}(\text{heavy}), \ \tilde{t}_{1} \rightarrow t \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}, \ \tilde{t}_{1} \rightarrow c \tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{2}, \ \tilde{t}_{2} \rightarrow \tilde{t}_{1} + Z \end{array} $	$\begin{array}{c} 0 \\ 2 \ e, \mu \ (\text{SS}) \\ 1-2 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 3 \ e, \mu \ (Z) \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b tono-jet/c-ta 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.7 20.7	\tilde{b}_1 \tilde{b}_1 \tilde{t}_2	100-620 GeV 275-430 GeV 110 <mark>-167 GeV</mark> 130-220 GeV 225-525 GeV 225-525 GeV 200-610 GeV 320-660 GeV 90-200 GeV 500 GeV 271-520 GeV	$\begin{array}{l} m(\tilde{k}_{1}^{0}) < 90 \ \text{GeV} \\ m(\tilde{k}_{1}^{+}) = 2 \ m(\tilde{k}_{1}^{0}) \\ m(\tilde{k}_{1}^{0}) = 55 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) = m(\tilde{t}_{1}) - m(\mathcal{W}) - 50 \ \text{GeV}, \ m(\tilde{t}_{1}) < < m(\tilde{k}_{1}^{+}) \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) < 200 \ \text{GeV}, \ m(\tilde{k}_{1}^{+}) - m(\tilde{k}_{1}^{0}) = 5 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ m(\tilde{k}_{1}^{0}) = 0 \ \text{GeV} \\ m(\tilde{t}_{1}) = 0 \ \text{GeV} \\ m(\tilde{t}_{1}) = 0 \ \text{GeV} \\ m(\tilde{t}_{1}) = m(\tilde{k}_{1}^{0}) < 85 \ \text{GeV} \\ m(\tilde{t}_{1}) = m(\tilde{k}_{1}^{0}) + 180 \ \text{GeV} \\ \end{array}$	1308.2631 ATLAS-CONF-2013-007 1208.4305, 1209.2102 ATLAS-CONF-2013-048 ATLAS-CONF-2013-065 1308.2631 ATLAS-CONF-2013-037 ATLAS-CONF-2013-024 ATLAS-CONF-2013-025 ATLAS-CONF-2013-025
EW direct	$\begin{array}{l} \tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\ell} \nu (\ell \tilde{\nu}) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\tau} \nu (\tau \tilde{\nu}) \\ \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{2}^{0} \rightarrow \tilde{\ell}_{L} \nu \tilde{\ell}_{L} \ell (\tilde{\nu} \nu), \ell \tilde{\nu} \tilde{\ell}_{L} \ell (\tilde{\nu} \nu) \\ \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} Z \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} h \tilde{\chi}_{1}^{0} \end{array}$	2 e,μ 2 e,μ 2 τ 3 e,μ 3 e,μ 1 e,μ	0 0 0 0 2 b	Yes Yes Yes Yes Yes Yes	20.3 20.3 20.7 20.7 20.7 20.7 20.3	$\vec{\ell}$ $\vec{\chi}_{1}^{\pm}$ $\vec{\chi}_{1}^{\pm}$ $\vec{\chi}_{1}^{\pm}$ $\vec{\chi}_{1}^{\pm}$ $\vec{\chi}_{2}^{\pm}$ $\vec{\chi}_{1}^{\pm}$ $\vec{\chi}_{2}^{0}$	85-315 GeV 125-450 GeV 180-330 GeV 600 GeV m($\tilde{\chi}_1^{\pm}$)= 315 GeV 285 GeV	$\begin{array}{l} m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV}, \; m(\tilde{\ell}, \bar{\nu}) = 0.5 (m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV}, \; m(\tilde{\tau}, \bar{\nu}) = 0.5 (m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; m(\tilde{\ell}, \bar{\nu}) = 0.5 (m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; sleptons \; decoupled \\ m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; sleptons \; decoupled \end{array}$	ATLAS-CONF-2013-049 ATLAS-CONF-2013-049 ATLAS-CONF-2013-028 ATLAS-CONF-2013-035 ATLAS-CONF-2013-035 ATLAS-CONF-2013-093
Long-lived particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q q \mu$ (RPV)	Disapp. trk 0 e, μ) 1-2 μ 2 γ 1 μ, displ. vtx	1 jet 1-5 jets - -	Yes Yes Yes -	20.3 22.9 15.9 4.7 20.3		270 GeV 832 GeV 475 GeV 230 GeV 1.0 TeV	$\begin{array}{l} m(\vec{k}_1^{+}) \cdot m(\vec{k}_1^{0}) {=} 160 \; \mathrm{MeV}, \; r(\vec{k}_1^{+}) {=} 0.2 \; \mathrm{ns} \\ m(\vec{k}_1^{0}) {=} 100 \; \mathrm{GeV}, \; 10 \; \mu \mathrm{s} {<} r(\vec{g}) {<} 1000 \; \mathrm{s} \\ 10 {<} \mathrm{tan}\beta {<} 50 \\ 0.4 {<} r(\vec{k}_1^{0}) {<} 2 \; \mathrm{ns} \\ 1.5 {<} c r {<} 156 \; \mathrm{mm}, \; \mathrm{BR}(\mu) {=} 1, \; m(\vec{k}_1^{0}) {=} 108 \; \mathrm{GeV} \end{array}$	ATLAS-CONF-2013-069 ATLAS-CONF-2013-057 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$ \begin{array}{l} LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e + \mu \\ LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow ee\tilde{v}_{\mu}, e\mu \tilde{v}_{e} \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow \tau \tau \tilde{v}_{e}, e\tau \tilde{v}_{\tau} \\ \tilde{g} \rightarrow qqq \\ \tilde{g} \rightarrow \tilde{t}_{1} t, \tilde{t}_{1} \rightarrow bs \end{array} $		- 7 jets - 6-7 jets 0-3 b	- Yes Yes - Yes	4.6 4.7 20.7 20.7 20.3 20.3	\tilde{v}_r \tilde{v}_r $\tilde{q}_r \tilde{g}$ $\tilde{\chi}_1^{\pm}$ $\tilde{\chi}_1^{\pm}$ \tilde{g} \tilde{g}	1.61 TeV 1.1 TeV 1.2 TeV 760 GeV 350 GeV 916 GeV 880 GeV	$\begin{array}{l} \lambda_{311}'=0.10,\lambda_{132}=0.05\\ \lambda_{311}'=0.10,\lambda_{1(2)33}=0.05\\ m(\tilde{q})=m(\tilde{g}),c\tau_{4.5P}<1mm\\ m(\tilde{k}_1^0)>300{\rm GeV},\lambda_{121}>0\\ m(\tilde{k}_1^0)>80{\rm GeV},\lambda_{133}>0\\ {\rm BR}(t)={\rm BR}(b)={\rm BR}(c)=0\% \end{array}$	1212.1272 1212.1272 ATLAS-CONF-2012-140 ATLAS-CONF-2013-036 ATLAS-CONF-2013-036 ATLAS-CONF-2013-091 ATLAS-CONF-2013-007
Other		2 e, µ (SS) 0 √s = 8 TeV		- Yes Yes 8 TeV	4.6 14.3 10.5	sgluon sgluon M* scale	100-287 GeV 800 GeV 704 GeV 0 ⁻¹ 1	incl. limit from 1110.2693 $m(\chi)$ <80 GeV, limit of<687 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
	full data	artial data	full c	lata	l	'	· ·	Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 or theoretical signal cross section uncertainty.

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

ATLAS Preliminary

 $\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$

	Model	e, μ, τ, γ	Jets	E_{T}^{miss}	∫£ dt[fb	-1] Mass limit	9	Reference
Inclusive Searches	MSUGRA/CMSSM MSUGRA/CMSSM MSUGRA/CMSSM $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q \tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q \bar{q} \tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{\chi}_{1}^{1} \rightarrow q q W^{\pm} \tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\ell \ell / \ell \nu / \nu \nu) \tilde{\chi}_{1}^{0}$ GMSB ($\tilde{\ell}$ NLSP) GGM (bino NLSP) GGM (bino NLSP) GGM (higgsino-bino NLSP) GGM (higgsino NLSP) GGM (higgsino NLSP) GGM (higgsino NLSP) Gravitino LSP	$\begin{array}{c} 0 \\ 1 \ e, \mu \\ 0 \\ 0 \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 1 - 2 \ \tau \\ 2 \ \gamma \\ 1 \ e, \mu + \gamma \\ \gamma \\ 2 \ e, \mu (Z) \\ 0 \end{array}$	2-6 jets 3-6 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 0-2 jets 1 <i>b</i> 0-3 jets mono-jet	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 4.7 20.7 4.8 4.8 4.8 4.8 5.8 10.5	q.ğ 1.1 r 1.7 rev ğ 1.2 rev ğ 1.1 r q 740 GeV ğ 1.1 r ğ 1.2 r ğ 1.0 r ğ 619 GeV ğ 690 GeV ğ 690 GeV ğ 690 GeV ğ 645 GeV	$ \begin{array}{l} m(\tilde{q}) = m(\tilde{g}) \\ any \ m(\tilde{q}) \\ any \ m(\tilde{q}) \\ m(\tilde{k}_1^0) = 0 \ {\rm GeV} \\ m(\tilde{k}_1^0) = 0 \ {\rm GeV} \\ m(\tilde{k}_1^0) < 200 \ {\rm GeV}, \ m(\tilde{\chi}^{\pm}) = 0.5(m(\tilde{\chi}_1^0) + m(\tilde{g})) \\ m(\tilde{\chi}_1^0) = 0 \ {\rm GeV} \\ tan \beta < 15 \\ tan \beta < 18 \\ m(\tilde{k}_1^0) > 50 \ {\rm GeV} \\ m(\tilde{k}_1^0) > 50 \ {\rm GeV} \\ m(\tilde{k}_1^0) > 220 \ {\rm GeV} \\ m(\tilde{k}_1^0) > 200 \ {\rm GeV} \\ m(\tilde{k}_1^0) > 10^{-4} \ {\rm eV} \end{array} $	ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 1308.1841 ATLAS-CONF-2013-047 ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 ATLAS-CONF-2013-089 1208.4688 ATLAS-CONF-2013-026 1209.0753 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152 ATLAS-CONF-2012-147
3 rd gen. <i>ἒ</i> med.	$\begin{array}{l} \tilde{g} \rightarrow b \tilde{b} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \tilde{t} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \tilde{t} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow b \tilde{t} \tilde{\chi}_{1}^{+} \end{array}$	0 0 0-1 e,μ 0-1 e,μ	3 b 7-10 jets 3 b 3 b	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	ğ 1.2 TeV ğ 1.1 TeV ğ 1.84 TeV ğ 1 3 TeV	$m(\tilde{k}_{1}^{0})<600 \text{ GeV}$ $m(\tilde{k}_{1}^{0})<350 \text{ GeV}$ $m(\tilde{k}_{1}^{0})<400 \text{ GeV}$ $m(\tilde{k}_{1}^{0})<300 \text{ GeV}$	ATLAS-CONF-2013-061 1308.1841 ATLAS-CONF-2013-061 ATLAS-CONF-2013-061
3rd gen. squarks direct production	$ \begin{array}{l} \tilde{b}_1 \tilde{b}_1, \ \tilde{b}_1 \rightarrow b \tilde{\chi}_1^0 \\ \tilde{b}_1 \tilde{b}_1, \ \tilde{b}_1 \rightarrow t \tilde{\chi}_1^\pm \\ \tilde{t}_1 \tilde{t}_1 (\text{light}), \ \tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm \\ \tilde{t}_1 \tilde{t}_1 (\text{light}), \ \tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0 \\ \tilde{t}_1 \tilde{t}_1 (\text{medium}), \ \tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0 \\ \tilde{t}_1 \tilde{t}_1 (\text{medium}), \ \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{t}_1 \tilde{t}_1 (\text{heavy}), \ \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{t}_1 \tilde{t}_1 (\text{heavy}), \ \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0 \\ \tilde{t}_1 \tilde{t}_1 (\text{netural GMSB}) \\ \tilde{t}_2 \tilde{t}_2, \ \tilde{t}_2 \rightarrow \tilde{t}_1 + Z \end{array} $	0 2 e, µ (SS) 1-2 e, µ 2 e, µ 2 e, µ 0 1 e, µ 0 2 e, µ (Z) 3 e, µ (Z)	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b ono-jet/c-t 1 b 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.7 20.7	$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{array}{l} m(\tilde{k}_{1}^{0}){<}90\text{GeV} \\ m(\tilde{k}_{1}^{0}){=}2m(\tilde{k}_{1}^{0}) \\ n(\tilde{k}_{1}^{0}){=}55\text{GeV} \\ (\tilde{k}_{1}^{0}){=}55\text{GeV} \\ (\tilde{k}_{1}^{0}){=}0\text{GeV} \\ ({>}200\text{GeV}, m(\tilde{k}_{1}^{1}){-}m(\tilde{k}_{1}^{0}){=}5\text{GeV} \\ ({>}200\text{GeV}, m(\tilde{k}_{1}^{1}){-}m(\tilde{k}_{1}^{0}){=}5\text{GeV} \\ ({>}0\text{GeV} \\ m(\tilde{t}_{1}){-}m(\tilde{k}_{1}^{0}){<}85\text{GeV} \\ m(\tilde{t}_{1}){-}m(\tilde{k}_{1}^{0}){+}150\text{GeV} \\ m(\tilde{t}_{1}){=}m(\tilde{k}_{1}^{0}){+}180\text{GeV} \\ \end{array}$	1308.2631 ATLAS-CONF-2013-007 1208.4305, 1209.2102 ATLAS-CONF-2013-048 ATLAS-CONF-2013-065 1308.2631 ATLAS-CONF-2013-037 ATLAS-CONF-2013-024 ATLAS-CONF-2013-068 ATLAS-CONF-2013-025 ATLAS-CONF-2013-025
EW direct	$\begin{array}{l} \tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\ell}\nu(\ell \tilde{\nu}) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\tau}\nu(\tau \tilde{\nu}) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{2}^{0} \rightarrow \ell_{L}\nu \tilde{\ell}_{L}\ell(\tilde{\nu}\nu), \ell \tilde{\nu} \tilde{\ell}_{L}\ell(\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} Z \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} h \tilde{\chi}_{1}^{0} \end{array}$	2 e, μ 2 e, μ 2 τ 3 e, μ 3 e, μ 1 e, μ	0 - 0 2 b	Yes Yes Yes Yes Yes Yes	20.3 20.3 20.7 20.7 20.7 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{array}{l} m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV} \\ m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV}, \; m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{0}) = 0 \; \mathrm{GeV}, \; m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ l = m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; sleptons \; decoupled \\ m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}), \; m(\tilde{\chi}_{1}^{0}) = 0, \; sleptons \; decoupled \end{array}$	ATLAS-CONF-2013-049 ATLAS-CONF-2013-049 ATLAS-CONF-2013-028 ATLAS-CONF-2013-035 ATLAS-CONF-2013-035 ATLAS-CONF-2013-093
Long-lived particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu})_+ \tau(q)$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)	0	1 jet 1-5 jets - -	Yes Yes Yes -	20.3 22.9 15.9 4.7 20.3	x ₁ [±] 270 GeV g 832 GeV x ₁ ⁰ 475 GeV x ₁ ⁰ 230 GeV q 1.0 TeV	$\begin{array}{l} m(\tilde{k}_1^{\pm})\text{-}m(\tilde{k}_1^{0}) {=} 160 \; \mathrm{MeV}, \; r(\tilde{k}_1^{\pm}) {=} 0.2 \; \mathrm{ns} \\ m(\tilde{k}_1^{0}) {=} 100 \; \mathrm{GeV}, \; 10 \; \mu \mathrm{s} {<} r(\tilde{g}) {<} 1000 \; \mathrm{s} \\ 10 {<} \tan \! \beta {<} 50 \\ 0.4 {<} r(\tilde{k}_1^{0}) {<} 2 \; \mathrm{ns} \\ 1.5 < c \tau {<} 156 \; \mathrm{mm}, \; \mathrm{BR}(\mu) {=} 1, \; m(\tilde{k}_1^{0}) {=} 108 \; \mathrm{GeV} \end{array}$	ATLAS-CONF-2013-069 ATLAS-CONF-2013-057 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$ \begin{array}{l} LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \ \tilde{v}_{\tau} \rightarrow e + \mu \\ LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \ \tilde{v}_{\tau} \rightarrow e(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \ \tilde{\chi}_{1}^{+} \rightarrow W \tilde{\chi}_{1}^{0}, \ \tilde{\chi}_{1}^{0} \rightarrow ee \tilde{v}_{\mu}, e\mu \tilde{v} \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \ \tilde{\chi}_{1}^{+} \rightarrow W \tilde{\chi}_{1}^{0}, \ \tilde{\chi}_{1}^{0} \rightarrow \tau \tau \tilde{v}_{e}, e \tau \tilde{v} \\ \tilde{g} \rightarrow q q q \\ \tilde{g} \rightarrow \tilde{t}_{1} t, \ \tilde{t}_{1} \rightarrow bs \end{array} $		- 7 jets - 6-7 jets 0-3 b	- Yes Yes - Yes	4.6 4.7 20.7 20.7 20.3 20.7	$\begin{array}{c c} \tilde{v}_r & 1.61 \text{ TeV} \\ \tilde{v}_r & 1.1 \text{ TeV} \\ \bar{v}_r & 1.2 \text{ TeV} \\ \bar{v}_1 & 760 \text{ GeV} \\ \tilde{x}_1^{\pm} & 760 \text{ GeV} \\ \tilde{x}_1^{\pm} & 350 \text{ GeV} \\ \tilde{s} & 880 \text{ GeV} \end{array}$	$\begin{array}{l} \lambda_{311}'=0.10,\lambda_{132}=0.05\\ \lambda_{311}'=0.10,\lambda_{1(2)33}=0.05\\ m(\tilde{q})=m(\tilde{g}),cr_{4,SP}<1\ mm\\ m(\tilde{k}_1^0)>300\ {\rm GeV},\lambda_{121}>0\\ m(\tilde{k}_1^0)>80\ {\rm GeV},\lambda_{133}>0\\ {\rm BR}(t)={\rm BR}(b)={\rm BR}(c)=0\% \end{array}$	1212.1272 1212.1272 ATLAS-CONF-2012-140 ATLAS-CONF-2013-036 ATLAS-CONF-2013-036 ATLAS-CONF-2013-091 ATLAS-CONF-2013-007
Other	Scalar gluon pair, sgluon $\rightarrow q\bar{q}$ Scalar gluon pair, sgluon $\rightarrow t\bar{t}$ WIMP interaction (D5, Dirac χ)	0 2 e, µ (SS) 0	4 jets 1 <i>b</i> mono-jet		4.6 14.3 10.5	sgluon 100-287 GeV sgluon 800 GeV M* scale 704 GeV	incl. limit from 1110.2693 $m(\chi)$ <80 GeV, limit of<687 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
		√s = 8 TeV partial data		8 TeV data		10 ⁻¹	Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 or theoretical signal cross section uncertainty.

The LHC frontier

- Theoretical model development is influenced by conclusions at 8 TeV, e.g. GUT scale SUSY models like cMSSM are pushed to higher scales and we are thinking of more non-minimal models
- The strategies for 13 TeV results depend on the conclusions at 8 TeV
- It is necessary to interpret the results in the most generic fashion and test as many models as possible

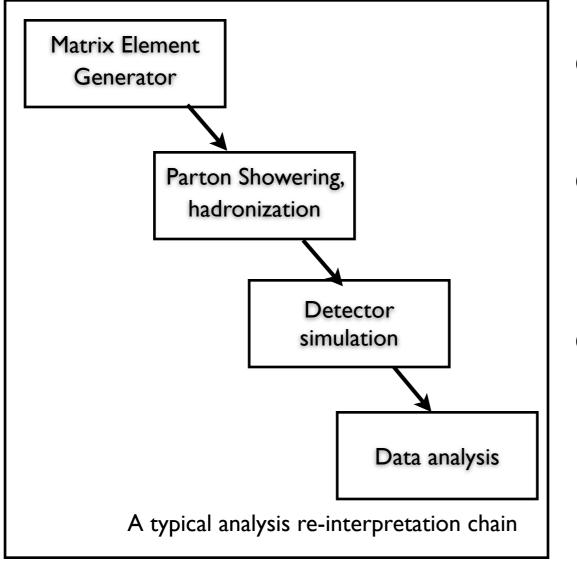
Analyses statistics (SUSY searches only)

	8 TeV (Any luminosity)	8 TeV (20 fb-1)
ATLAS	39	23
CMS	28	20

Huge number of searches Easier said than done!

A way to test our favorite BSM model against LHC results should exist

Traditional approach

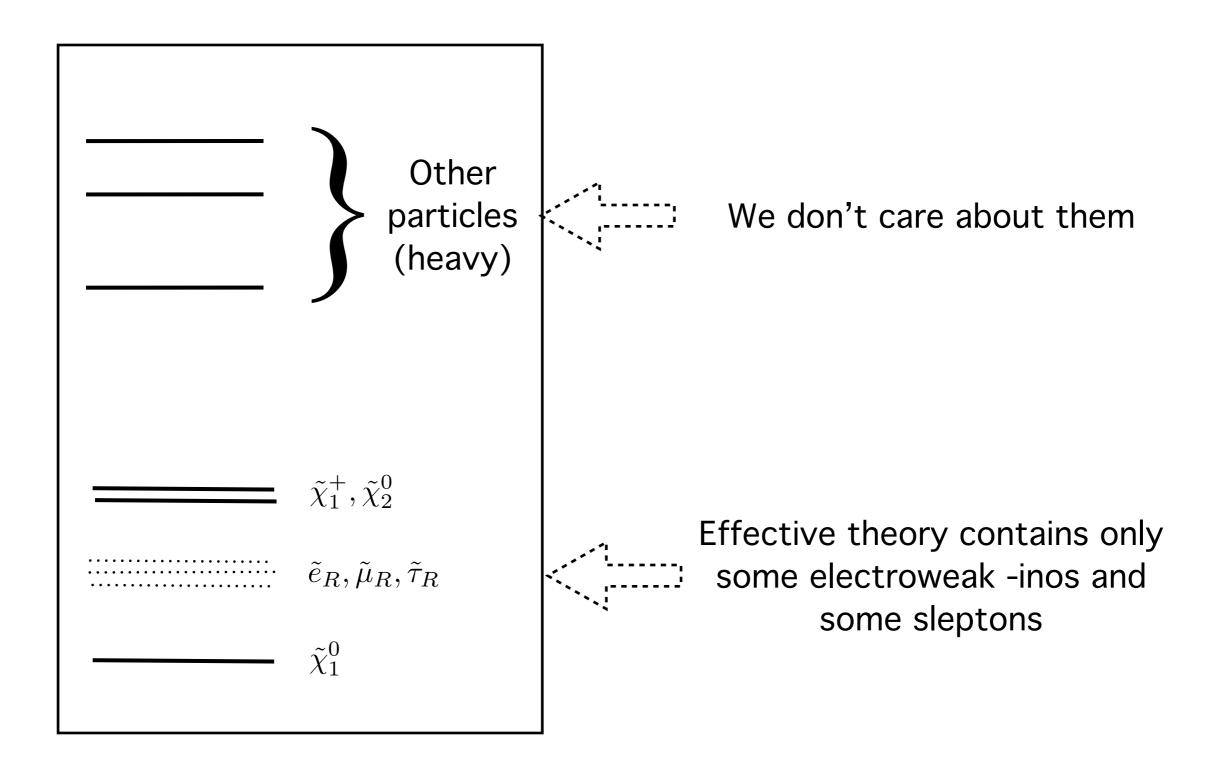


- Interpretation of LHC searches are model dependent
- Model dependence comes while converting the number of events observed to a limit on particle masses
- For a more generic case:
 - 1. Re-interpret the results yourself
 - 2. Use simplified model spectra results
- Re-interpreting the results yourself involves re-implementing the analysis, requires expertise, large computing power, time consuming
- We stick to simplified models results

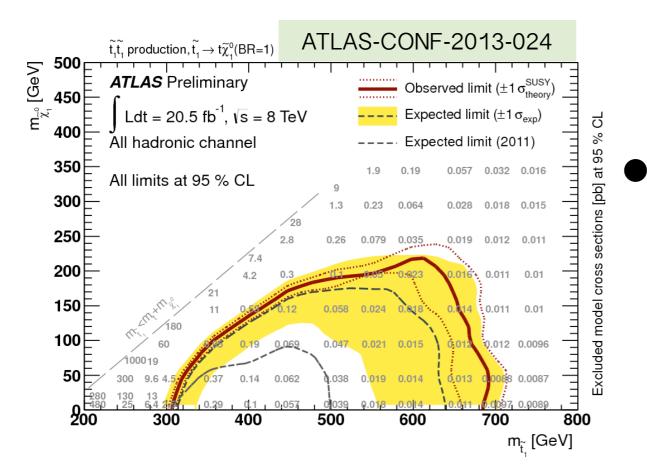
What is an SMS result?

• SMS are an effective-Lagrangian description of BSM involving a limited set of new particles.

What is an SMS result?



What is an SMS result?

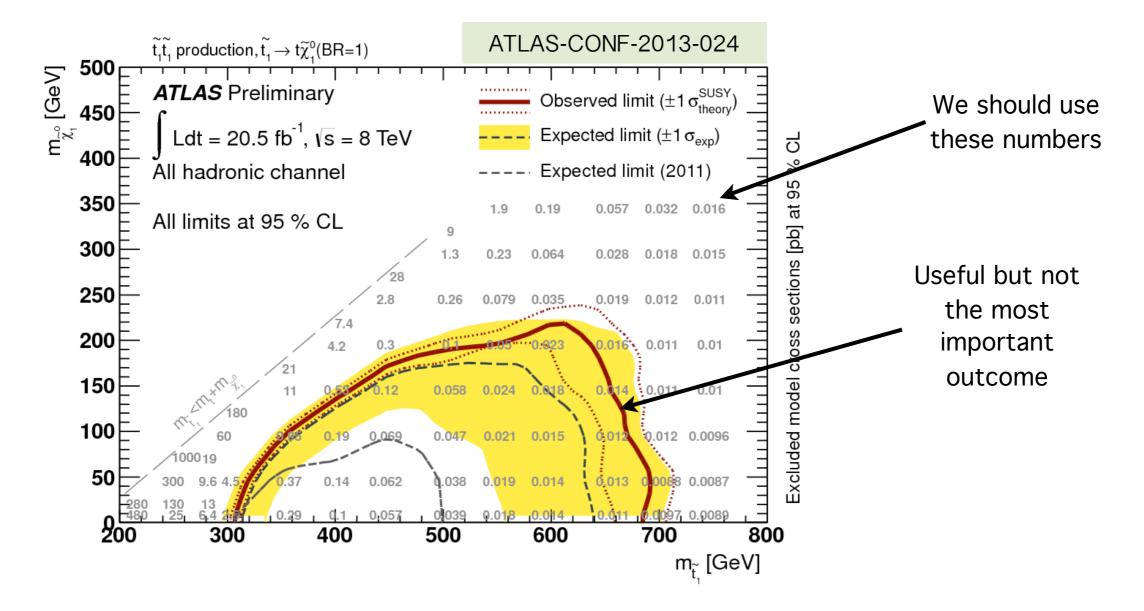


Note: the grid numbers on the plot are more important than the exclusion lines

Every SMS interpretation is based on a set of assumptions and is applicable for specific topologies e.g. ttbar + MET

 A generic point in e.g. SUSY parameter space contains many topologies and is sensitive to more than one SMS interpretation e.g. ttbar + MET, bbar + MET

How to read an SMS result



95% CL UL is the unfolded maximum amount of cross-section allowed for a specific decay chain and a mass combination

Is sigmaXBR(ttbar + MET, Mother mass, LSP mass) of your model > the number on the plot? -- Yes, point excluded; No, point allowed

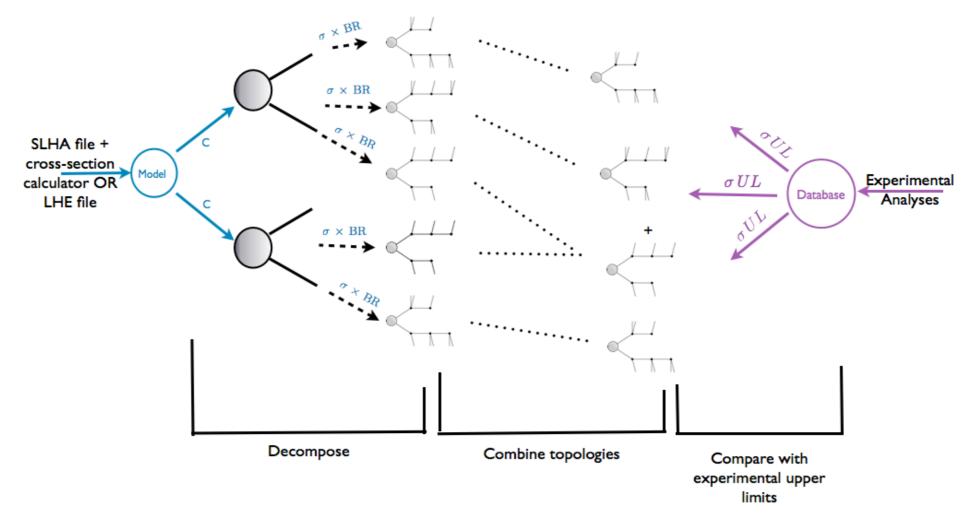
Can we have a centralized database of all the SMS results to check a given SUSY point in parameter space by decomposing it into SMS topologies?

Central concept of

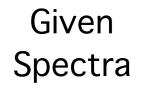


SModelS framework

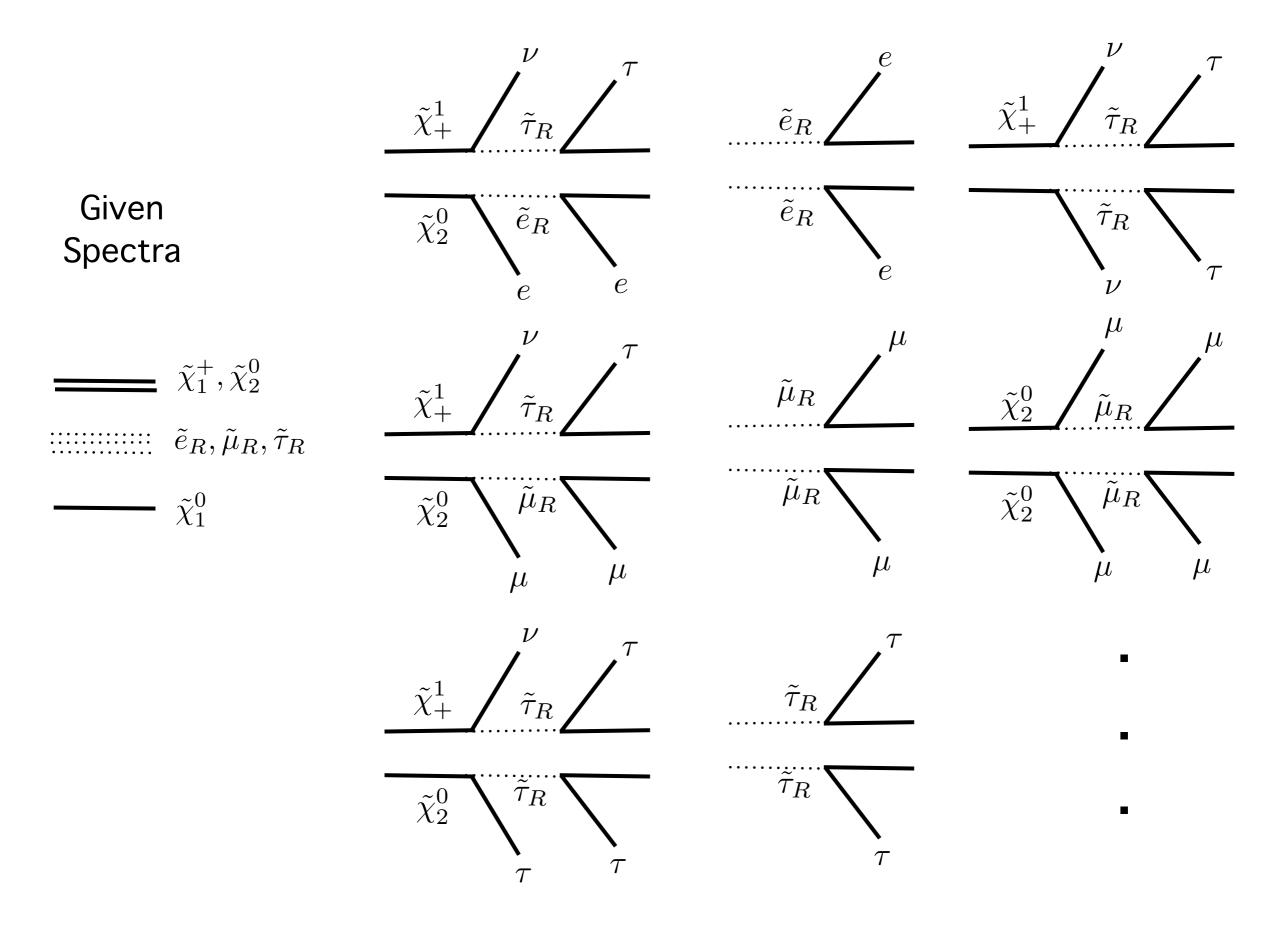
 It assumes, for most experimental searches, the BSM model can be approximated by a sum over effective simplified models

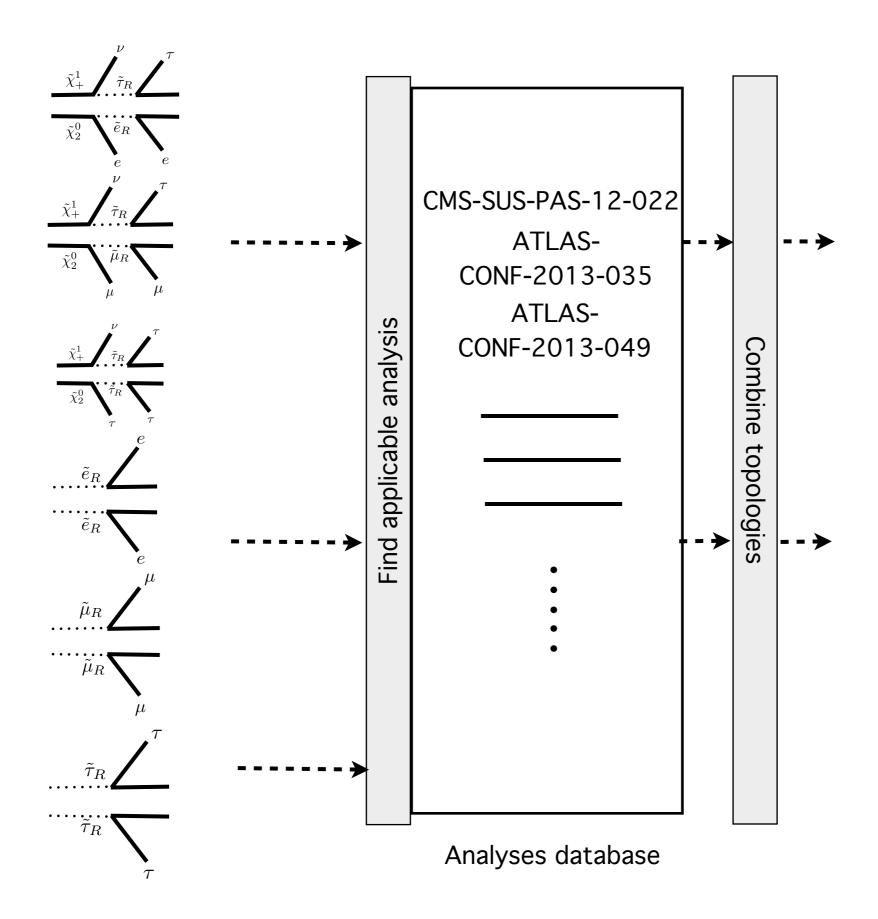


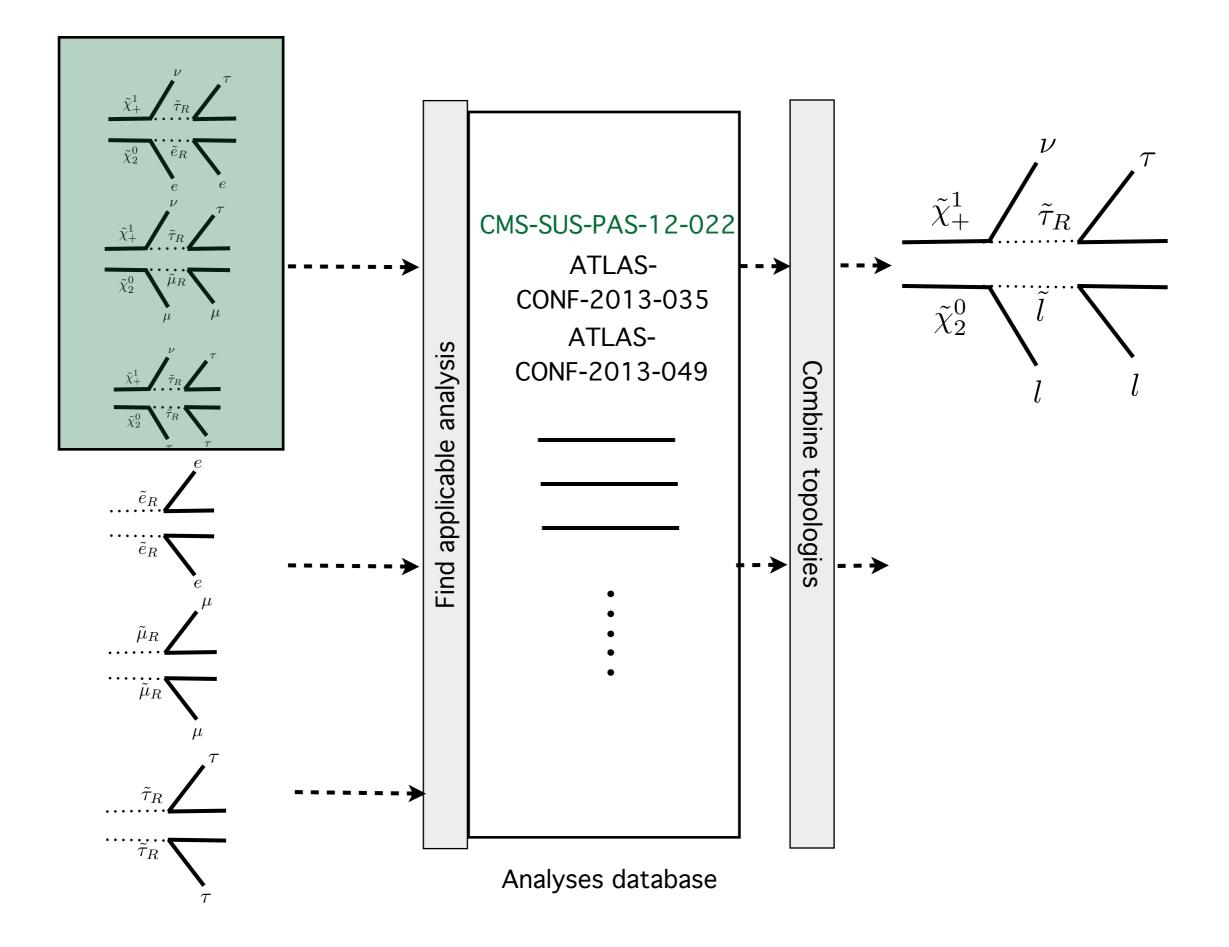
• Current implementation assumes R-parity is conserved

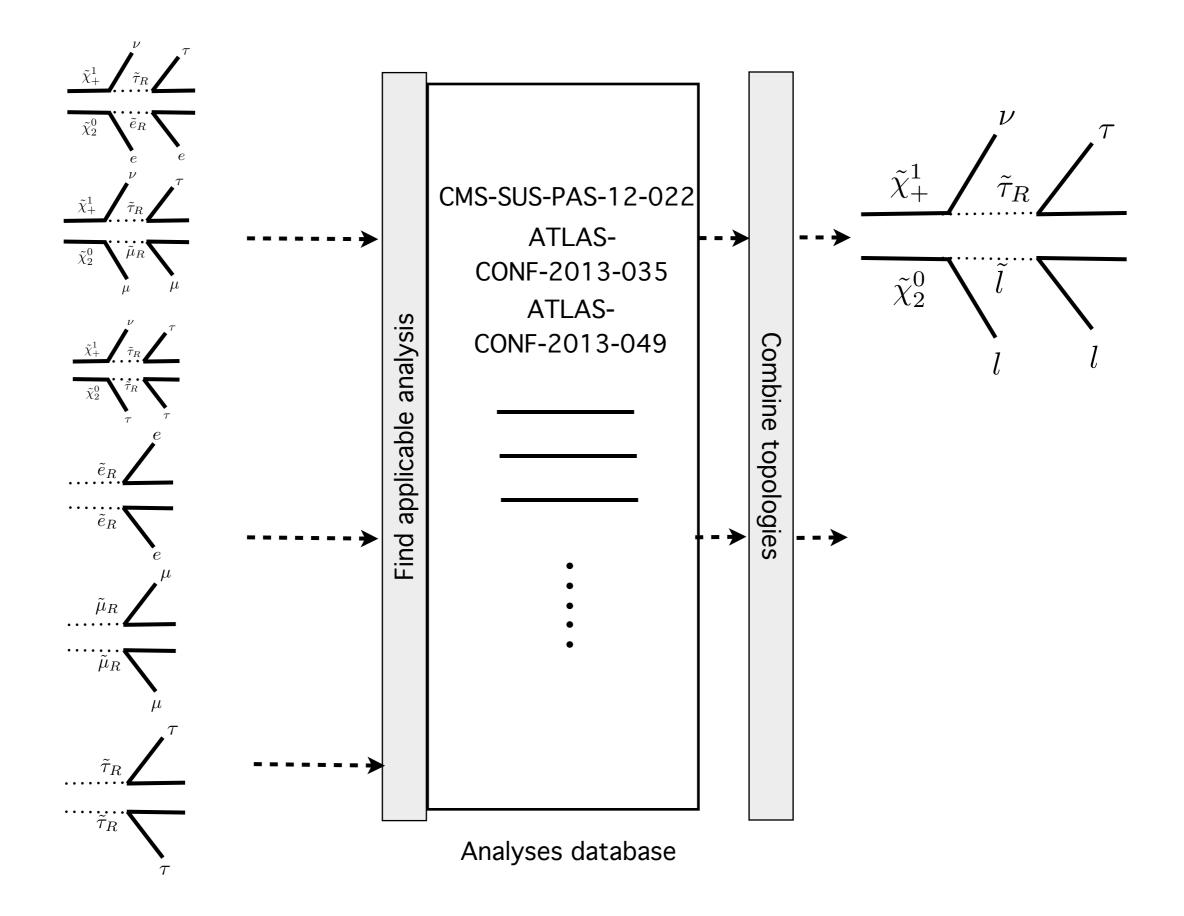


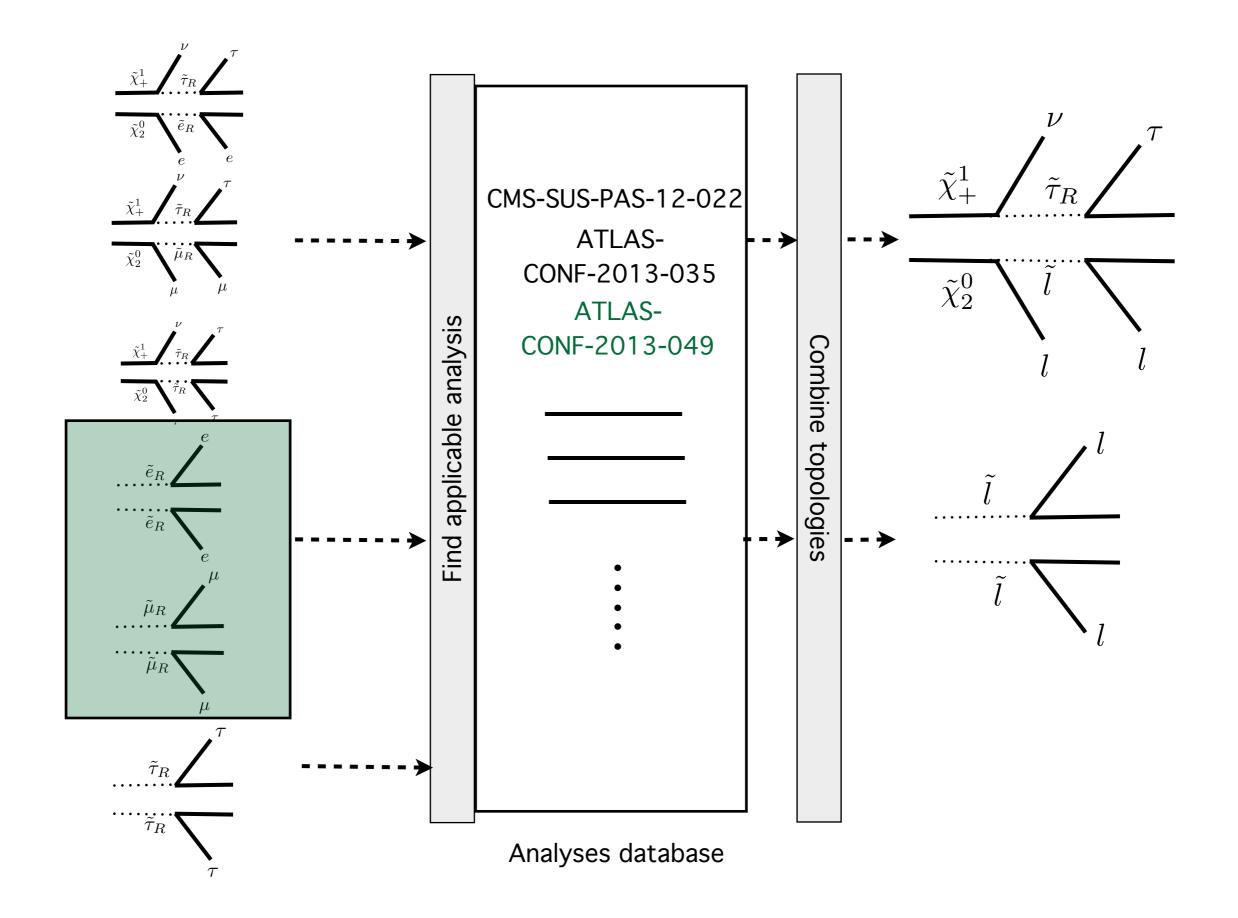
Decomposition

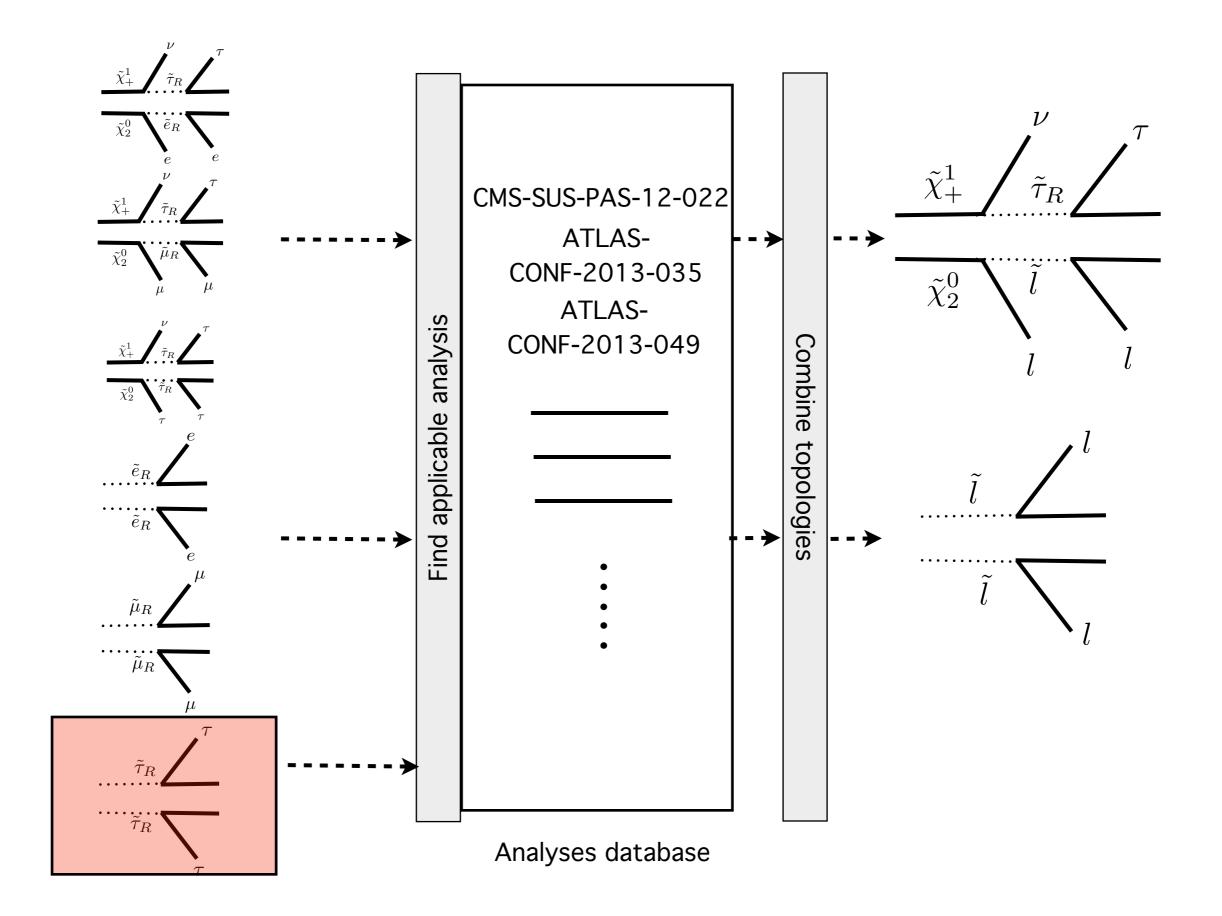


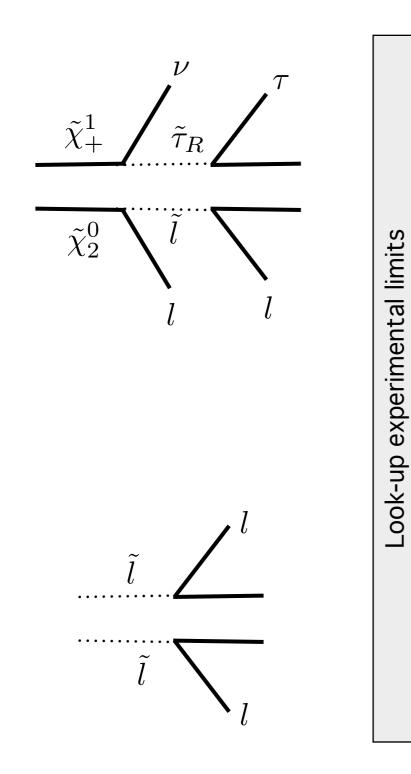


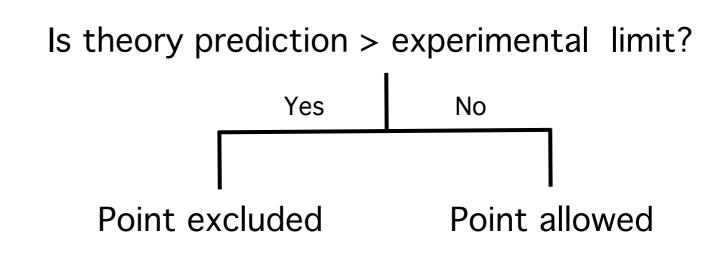






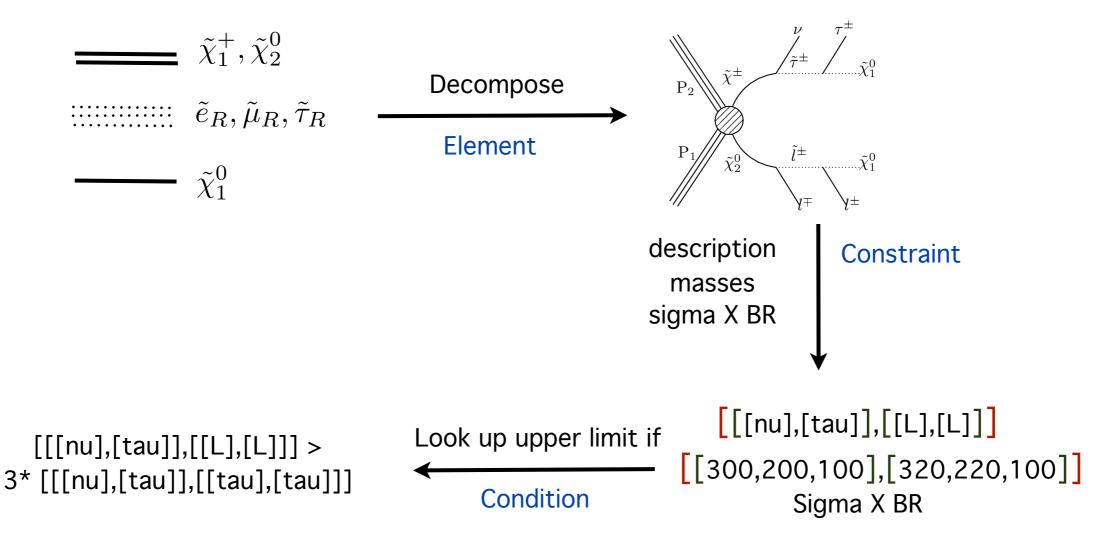






SModelS framework

Consider:

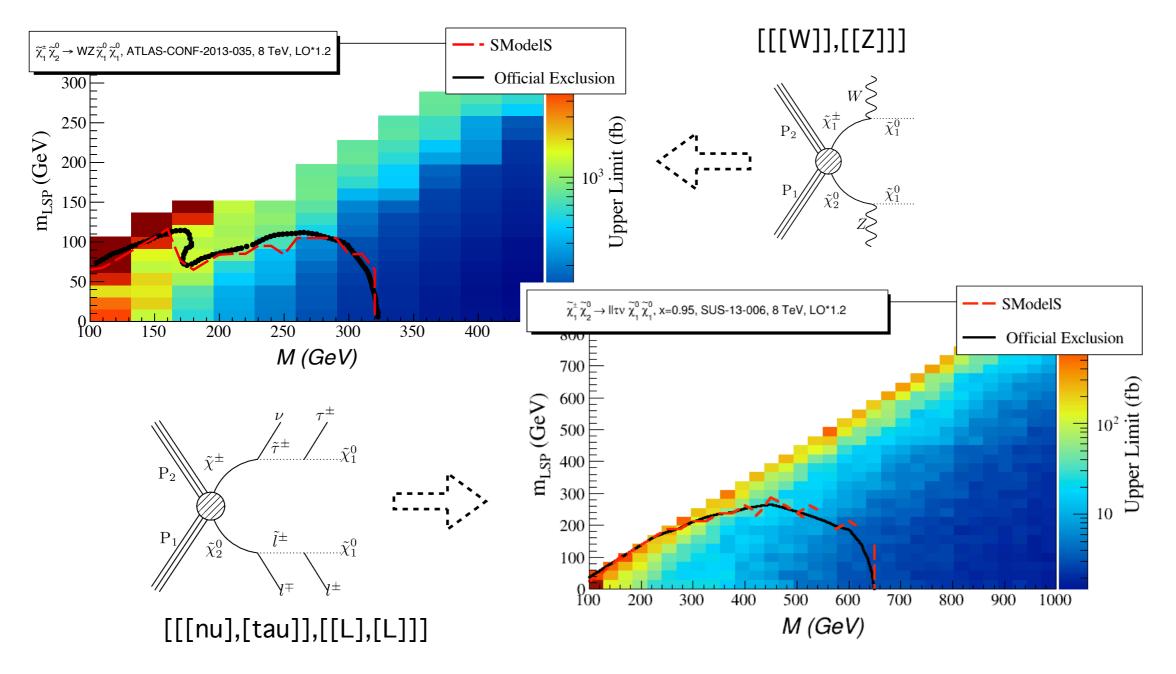


 The framework does not depend on characteristics of SUSY particles, can also be applied to decompose any BSM spectra of arbitrary complexity

SModelS language

How do we know it works?

 The code has been validated through the reproduction of various SMS exclusion curves



Typical examples of validation plot

Salient features

- Code is equipped to decompose any BSM model with a Z2 symmetry
- It can handle compressed topologies
- It can take care of invisible decays
- It has the most comprehensive database of simplified model results, 22 CMS, 24 ATLAS (7 + 8 TeV)
- Now a web SLHA interface is available to check your point

Good, so what do you learn out of it?

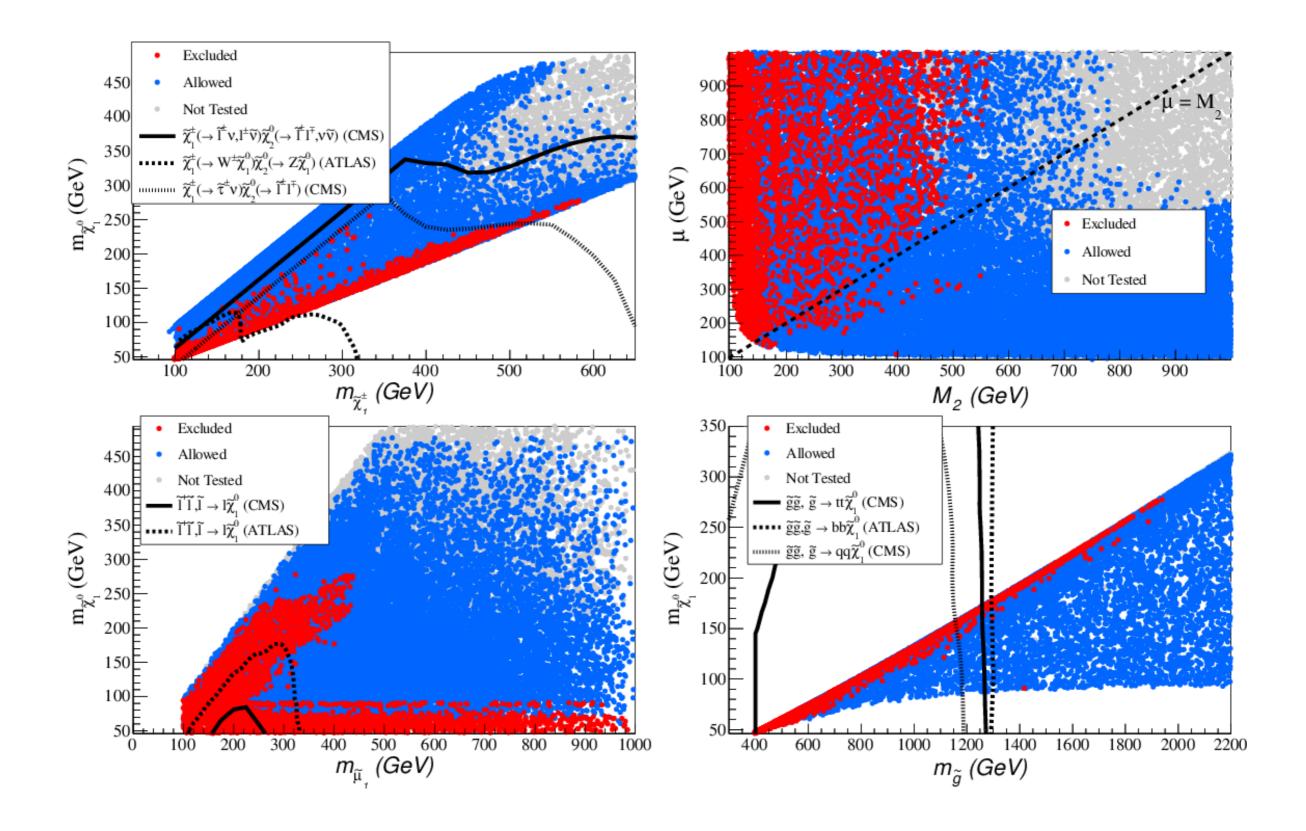
based on: arXiv: 1312.4175 [hep-ph]

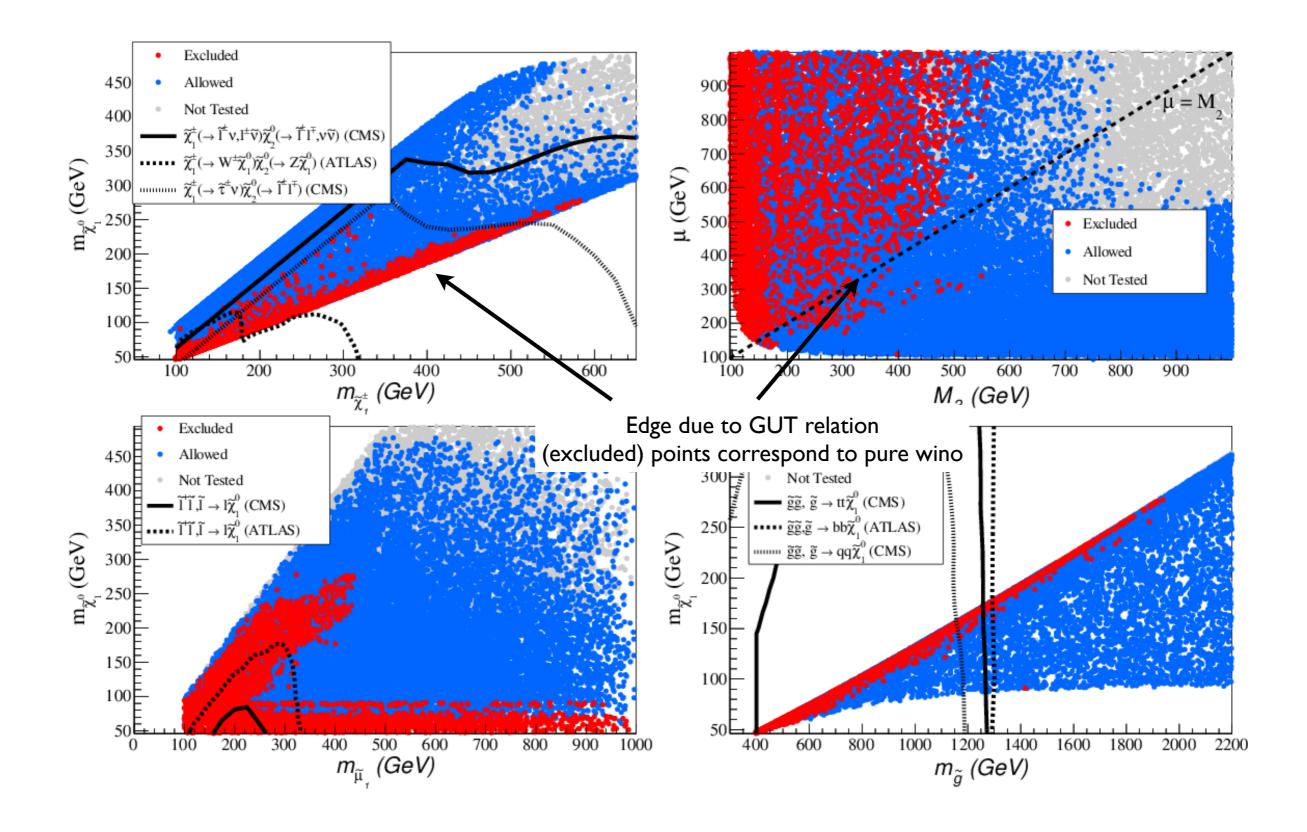
• pMSSM scan over 6 parameters

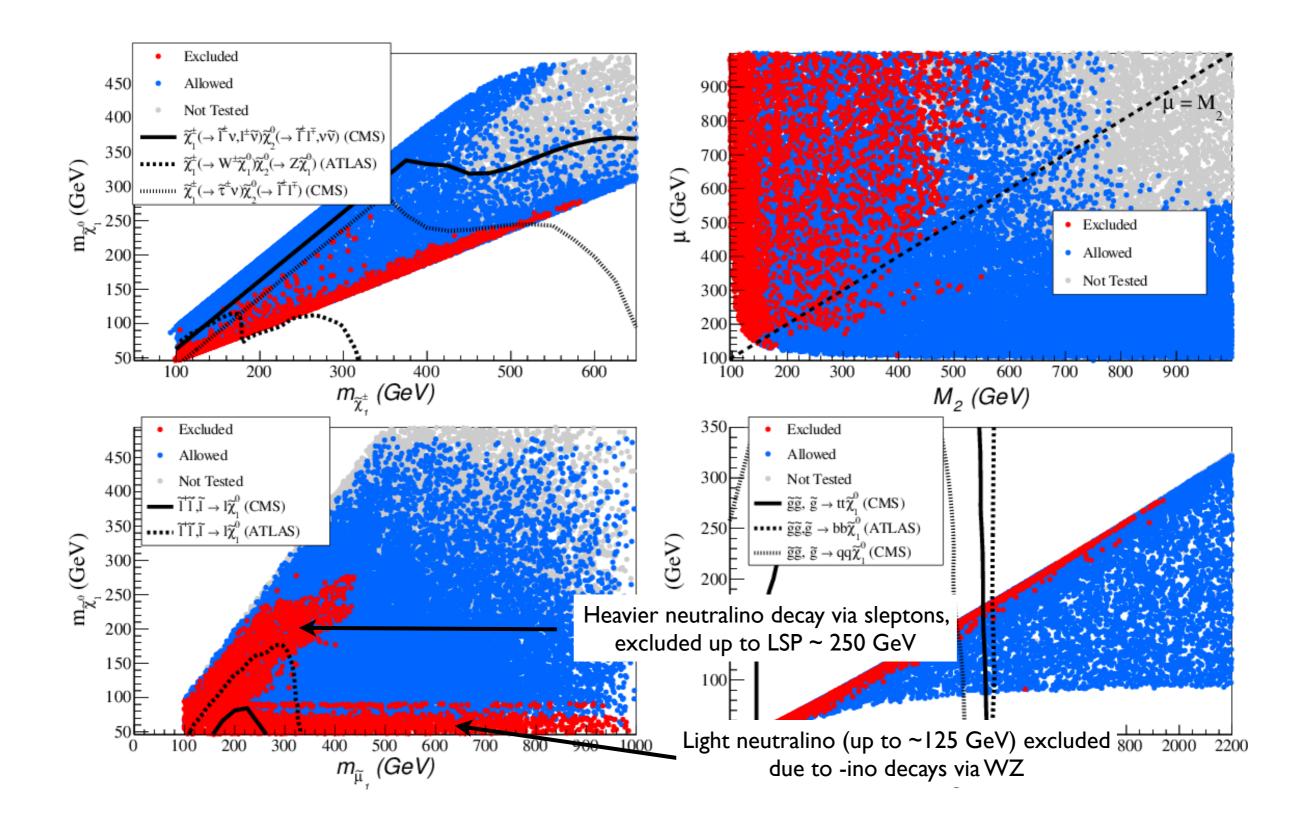
_	M_2	μ	aneta	$M_{ ilde{L}}$	$M_{\tilde{E}}$	A_{τ}
	0.1 - 1	0.1 - 1	3-60	0.1 - 1	0.1 - 1	± 1

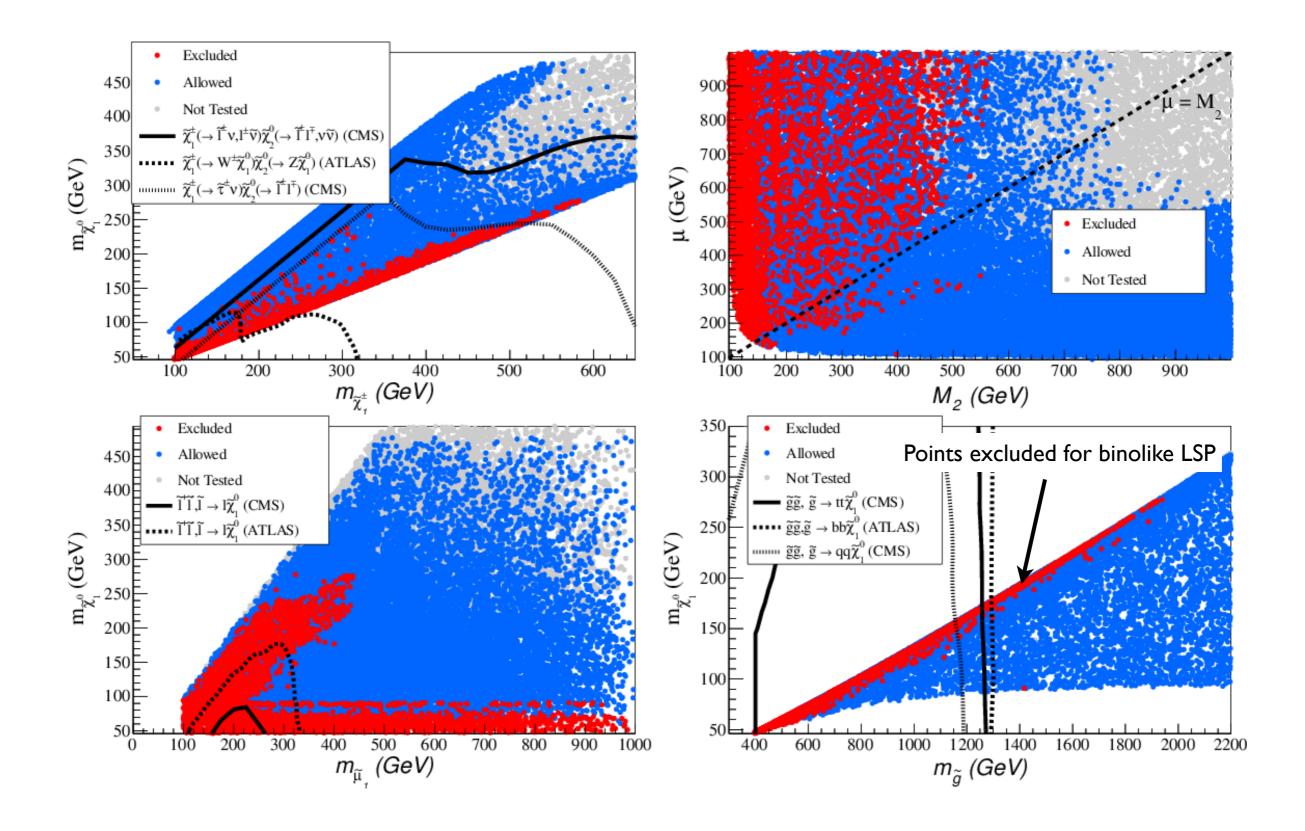
- Gaugino masses obey GUT relation
- Flavor constraints, invisible Z width, Higgs mass, LEP limits imposed
- Limits obtained will always be conservative
- We probe electroweak -ino decays via WZ and sleptons, direct slepton production and gluino decays

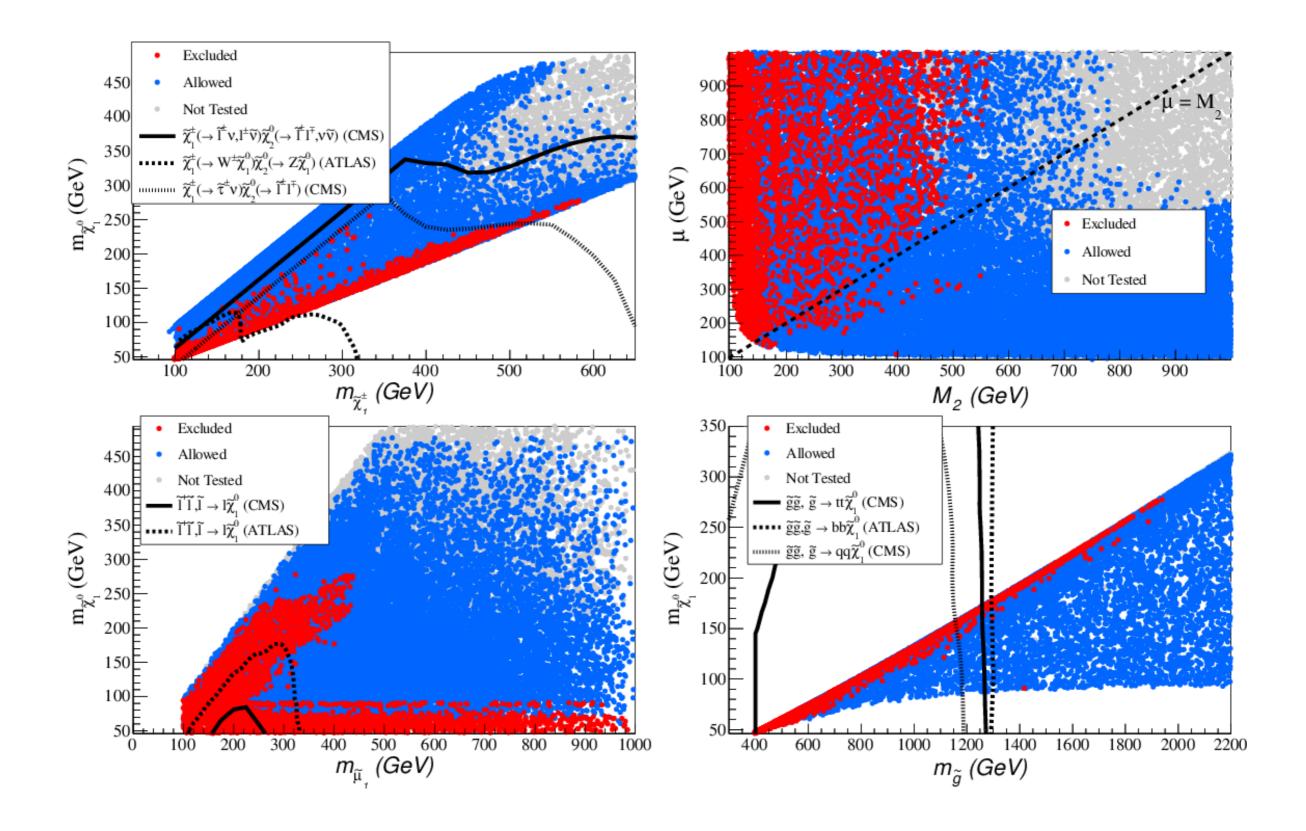
NB: A similar scan was also performed for strong sector particles



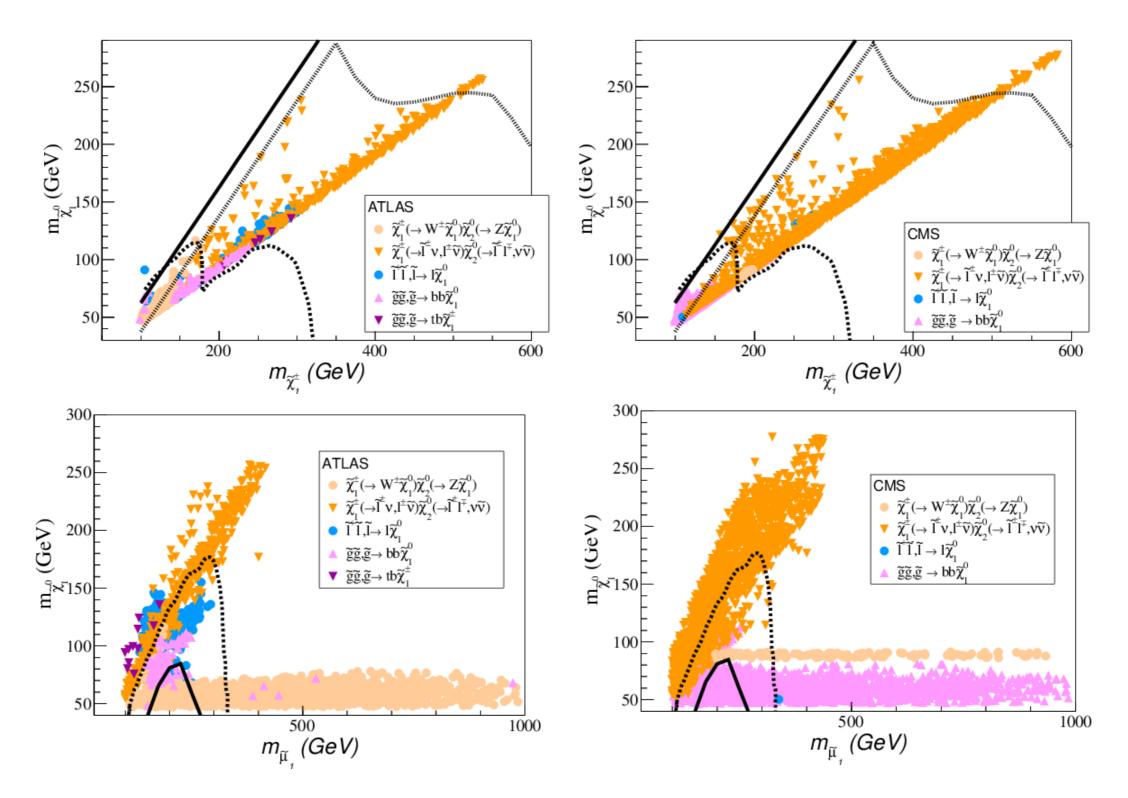




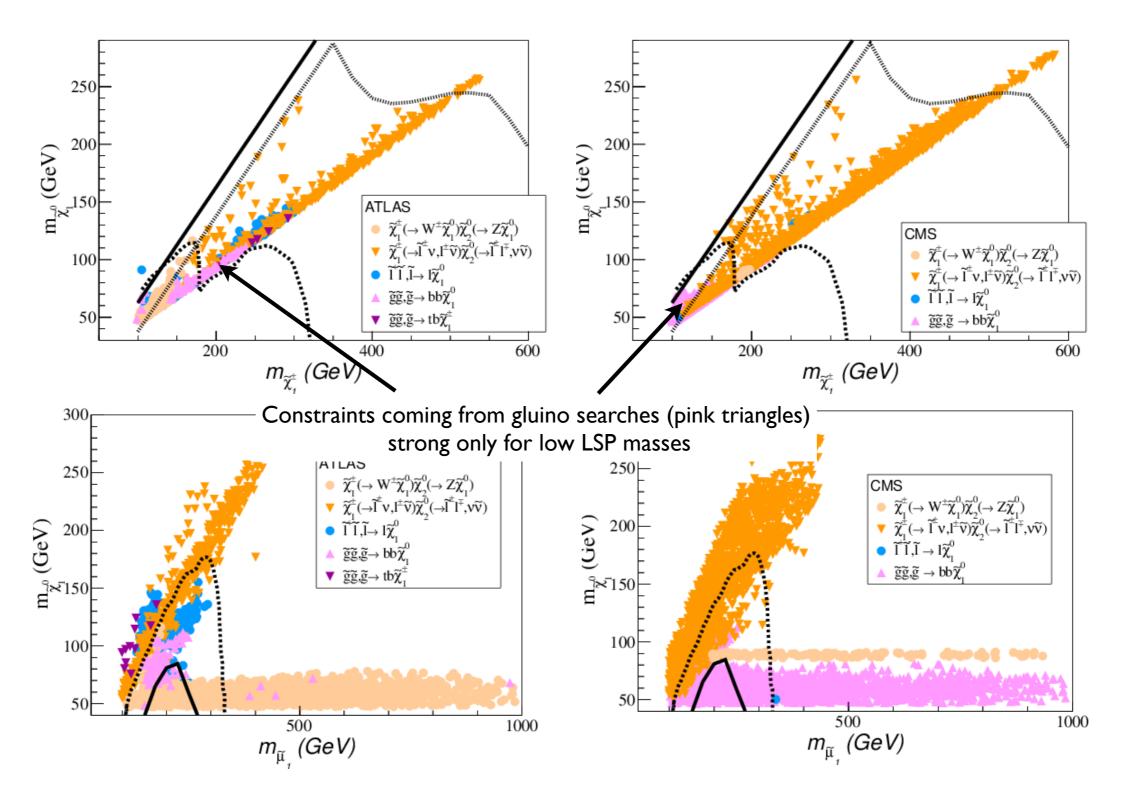




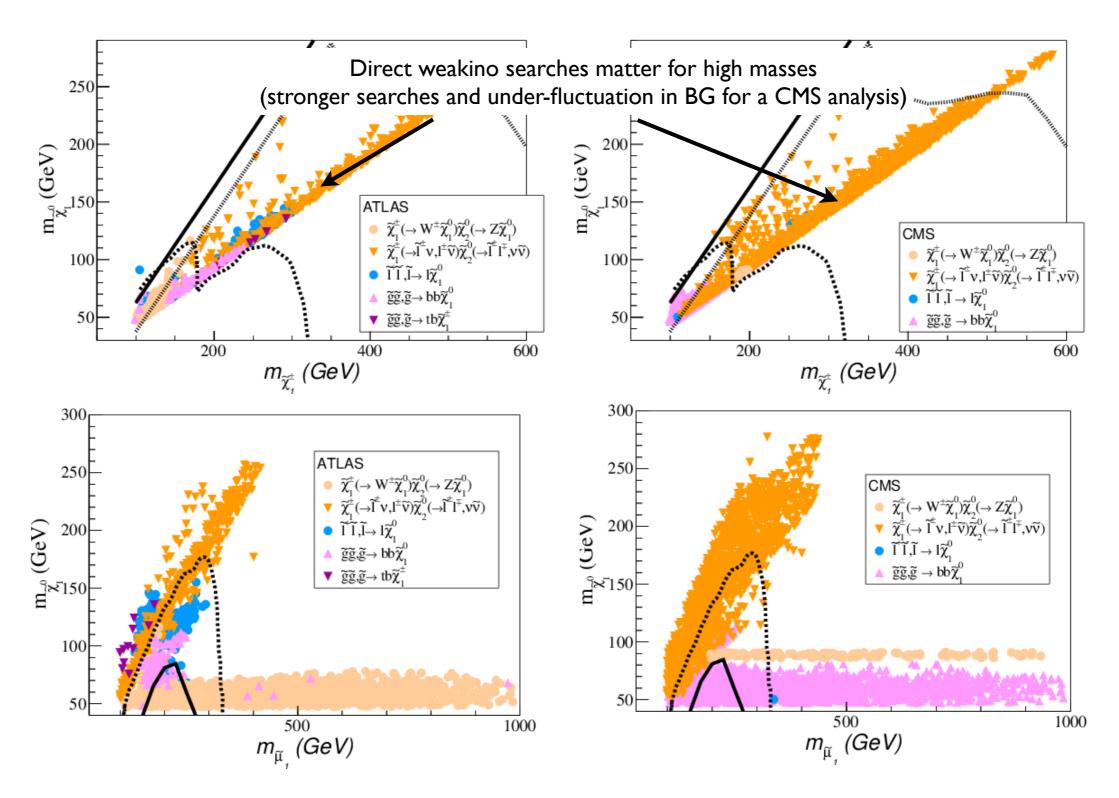
Breakdown of the excluded parameter space by analysis



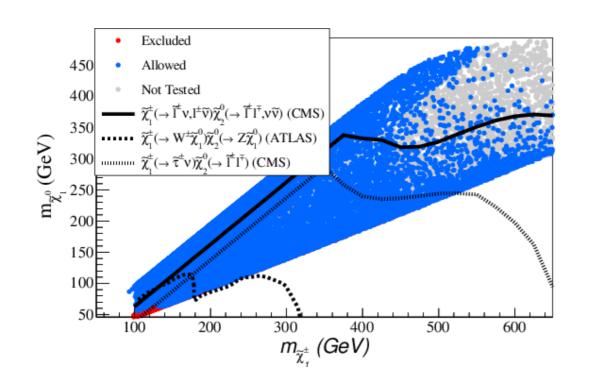
Breakdown of the excluded parameter space by analysis



Breakdown of the excluded parameter space by analysis

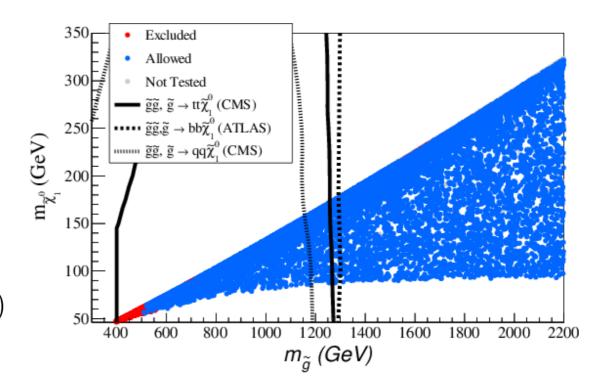


Some allowed points may lie below excluded points



- Gluino decays via on-shell squarks are kinematically forbidden for small masses
- Uncovered gluino decay topologies e.g. $BR(\tilde{g} \rightarrow \tilde{\chi}^{\pm} + tb)$

- Chargino LSP nature, higgsino have smaller production cross-section
- Right handed sleptons have smaller production CS



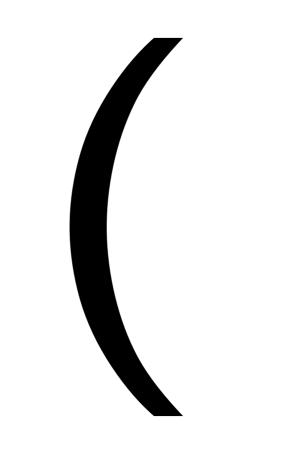
Can it be used to test parameter space for some interesting scenario?

A real life application

Do LHC results on the SUSY particles, Higgs signal strengths and constraints on DM from direct and indirect detection experiments rule out light neutralino DM?

Already many studies exist in literature, I'll not list them here

based on: arxiv:1308.3735 [hep-ph] (published PLB)



 $Ht\bar{t}$ coupling, leading to a strong correlation with the ttH process, this need not be the case in models with suppressed $Ht\bar{t}$ coupling and/or enhanced $Hb\bar{b}$ coupling and most especially in models with BSM loops.

The final states in which the Higgs is observed include $\gamma\gamma$, $ZZ^{(*)}_{Q}WW^{(*)}$, $b\bar{b}$ and $\tau\tau$. However, they do not all scale independently. In particular, custodial symmetry implies that the branching fractions into $ZZ^{(*)}$ and $WW^{(*)}$ are rescaled by the same factor with respect to the SM. We are then left with two independent production modes (VBF+VH) and (cgF+ttH), and for independent production modes (VBF+VH) and (cgF+ttH), and for independent production modes (VBF+VH) and (cgF+ttH), and coupling to Four independent production modes (VBF+VH) be and (rgF+ttH), and coupling to Four independent production modes (VBF+VH) be and (rgF+ttH), and coupling to Four independent production modes (VBF+VH) be and $\tau\tau$ rescale by a

• common factor, leading to identical μ values for the $b\bar{b}$ and $\tau\tau$ final states. Five first purpose of the present paper and contained the production provided by ATLAS, Clinclependential UZE and BW Wei electrony course to dials symmetry eluction ding the error completies among the (VBF+VH) and (ggF+ttH) production modes. Using a Gaussian approximation of privation electron and the approximately state of the present paper is the error completies among the (VBF+VH) and (ggF+ttH) production modes. Using a Gaussian approximation of private electron of the present paper is the error completies among the (VBF+VH) and (ggF+ttH) production modes. Using a Gaussian approximation of private electron of the present paper is the error completies among the error completies among the (VBF+VH) and (ggF+ttH) production modes. Using a Gaussian approximation of private electron of the present paper is the error completies among the error completies among the error of the present of the present paper is the error production modes. Using a Gaussian approximation of the present paper is the error of th

• where index i stands for $\gamma\gamma$, $VV^{(*)}$, $b\bar{b}$ and $\tau\tau$ (or $b\bar{b} = \tau\tau$), and $\hat{\mu}_i^{\text{ggF}}$ and $\hat{\mu}_i^{\text{VBF}}$ denote the bestfit points obtained from the μ neas (remeBR)_i (or thuBR))^M in "combined likelihood ellipses", which can be used in a simple, generic way to constrain non-standard Higgs sectors and new contributions to the loop-induced processes, provided they have the same Lagrangian structure as for Sach final state

In particular, these likelihoods can be used to derive constraints on a model-dependent choice of generalized Higgs couplings, the implications of which we study subsequently for

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The final states in which the Higgs is observed include $\gamma\gamma$, $ZZ^{(*)}_{Q}WW^{(*)}$, $b\bar{b}$ and $\tau\tau$. However, they do not all scale independently. In particular, custodial symmetry implies that the branching fractions into $ZZ^{(*)}$ and $WW^{(*)}$ are rescaled by the same factor with respect to the SM. We are then left with two independent production modes (VBF+VH) and (reF+ttH), and four independent production modes (VBF+VH, geF+ttH), and four independent final states $\gamma\gamma$, $VV^{(*)}$, bb, $\tau\tau$. In addition, in many models there is a common • coupling to Four tindependent final states ranching fiberious into $b\bar{b}$ and $\tau\tau$ rescale by a complete for the first purpose of the present paper and complete the first purpose of the present paper and complete the mean provided by $\Delta TLAS$,

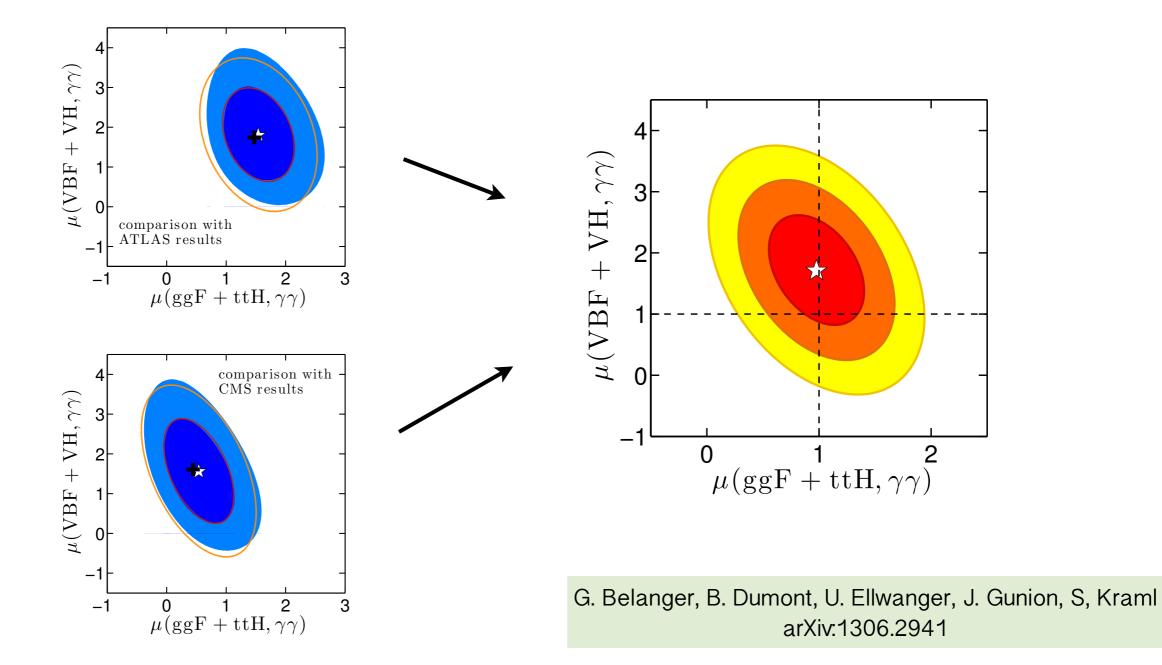
$$\mathcal{L} = g \left[C_V \left(M_W W_\mu W^\mu + \frac{M_Z}{\cos \theta_W} Z_\mu Z^\mu \right) - C_U \frac{m_t}{2M_W} \bar{t}t - C_D \frac{m_b}{2M_W} \bar{b}b - C_D \frac{m_\tau}{2M_W} \bar{\tau}\tau \right] H.$$

C's scale couplings relative to SM ones; $C_U=C_D=C_V=1$ is SM.

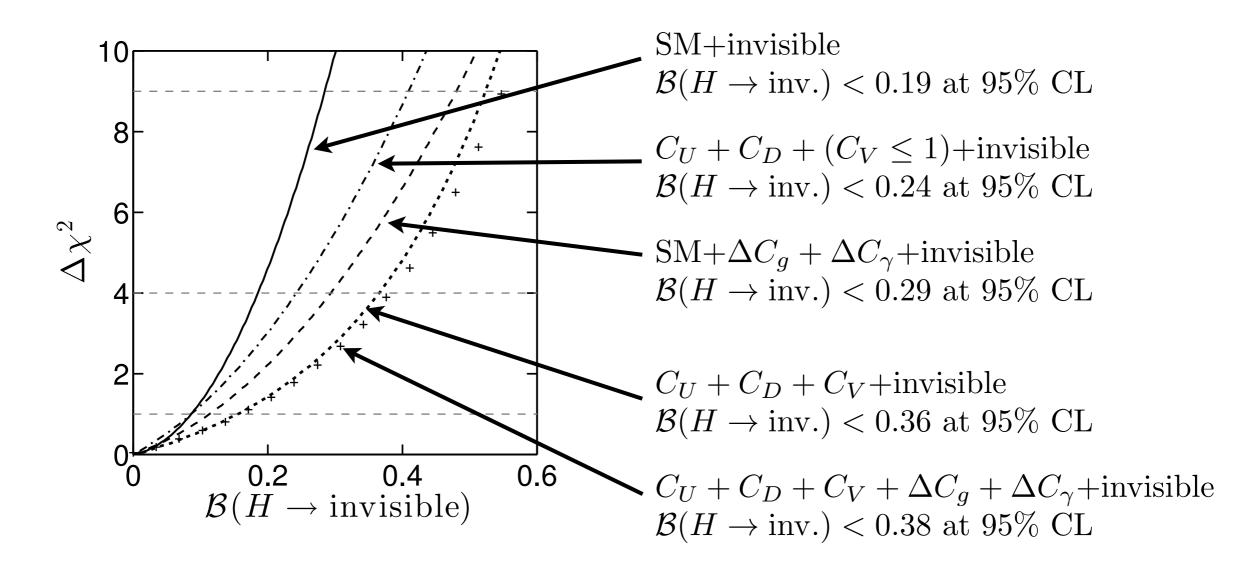
 $\chi_i^2 = q_i (\mu_i^{\text{ggF}} - \hat{\mu}_i^{\text{ggF}})^2 + 2b_i (\mu_i^{\text{ggF}} - \hat{\mu}_i^{\text{ggF}}) (\mu_i^{\text{VBF}} - \hat{\mu}_i^{\text{VBF}}) + c_i (\mu_i^{\text{VBF}} - \hat{\mu}_i^{\text{VBF}})^2,$ Additional loop contribution modify the couplings to gluons and (1)where the upper indices ggF and VBF stand for (ggF+ttH) and (VBF+VH), respectively, the lower index i stands for $\gamma\gamma$, $VV^{(*)}$, $b\bar{b}$ and $\tau\tau$ (or $b\bar{b} = \tau\tau$), and $\hat{\mu}_i^{\text{ggF}}$ and $\hat{\mu}_i^{\text{VBF}}$ denote the bestfit points obtained from the measurements. We thus obtain "combined likelihood ellipses", which can be used in a simple, generic way to constrain non-standard Higgs sectors and new vendred 27 septembre 2013 contributions to the loop-induced processes, provided they have the same Lagrangian structure as the SM.

In particular, these likelihoods can be used to derive constraints on a model-dependent choice of generalized Higgs couplings, the implications of which we study subsequently for

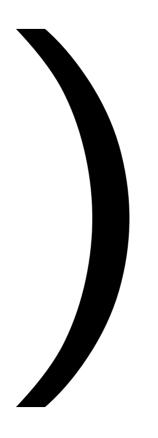
 A combined likelihood in (ggF+ttH) and (VBF+VH) planes was derived using ATLAS, CMS and Tevatron results



How much invisible Higgs decay is allowed?



G. Belanger, B. Dumont, U. Ellwanger, J. Gunion, S, Kraml arXiv:1306.2941



A real life application

Do LHC results on the SUSY particles, Higgs signal strengths and constraints on DM from direct and indirect detection experiments rule out light neutralino DM?

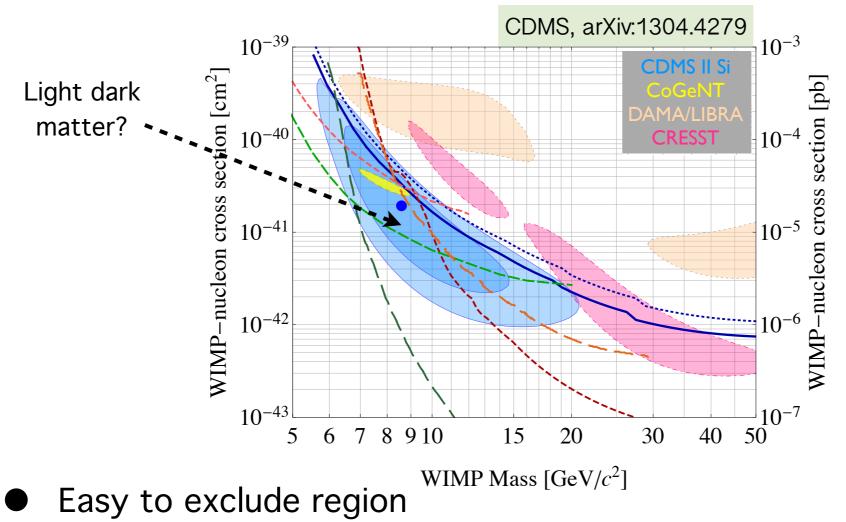
Already many studies exist in literature, I'll not list them here

based on: arXiv:1308.3735 [hep-ph] (published PLB)

Why light neutralino?

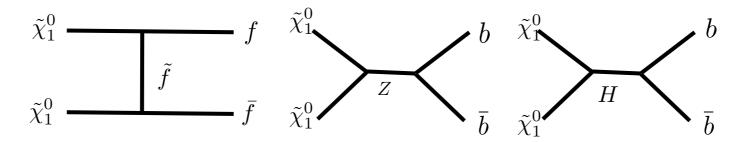
 $\stackrel{(\widetilde{B}, \widetilde{W}^0, \widetilde{H}^0_d, \widetilde{H}^0_d, \widetilde{H}^0)}{\bullet} \stackrel{\mathrm{EWSB}}{\to} \stackrel{(\widetilde{\chi}^0_1, \widetilde{\chi}^0_2, \widetilde{\chi}^0_3, \widetilde{\chi}^0_4)}{\bullet}$

Hints from direct (and may be indirect detection) ~ 10 GeV

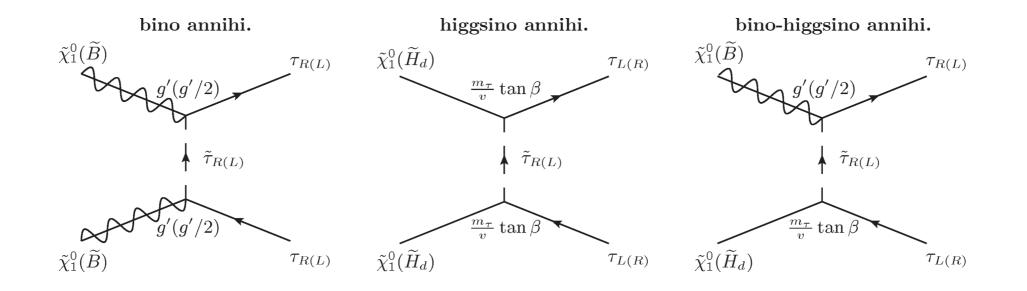


- No resonance below 45 GeV (Mz/2)
- No co-annihilation under 100 GeV (LEP limits) (counter example light sbottoms) arXiv:1308.2153

- Relaxing gaugino universality: few collider constraints
 - Z width, LEP bounds, invisible Higgs decays
- Most important annihilation channels:



- Region of interest: $m_{\tilde{\chi}_1^0} < m_h/2$ Z and H exchange not effective
- Light slepton exchange of interest to us here

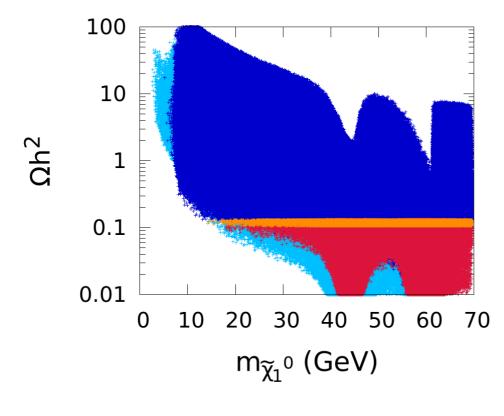


RH stau annihilation is more efficient, also get enhancements for high tan(beta) and higgsino LSP

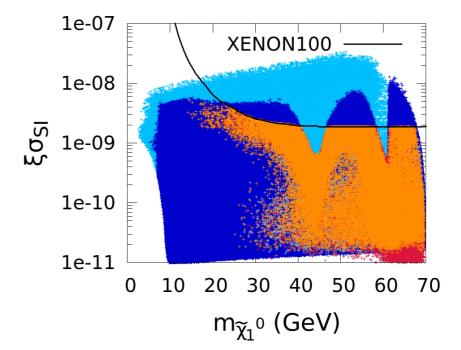
Light chargino	LEP and LHC	\checkmark
Invisible Z, Higgs decays	LEP and LHC	\checkmark
Light neutralino 2	$\text{LEP}\sigma(e^+ e^- \to \tilde{\chi}_1^0 \tilde{\chi}_2^0)$	\checkmark
Slepton and stau	LEP and LHC	\checkmark

- pMSSM scan over relevant parameters pMSSM scan over 11 parameters – $M_1, M_2, \mu, \tan\beta, M_A, A_t, M_{IL}, M_{IR}, M_{3L}, M_{3R}, A_t$
 - $M_{L} \neq P^{1}, Hint \in S^{3}, Hivisible Z^{M}, B = p = Mysics, R Higgs mass$
 - <u>teouplings, heavy Higgs@LHC, Xenon100</u>

@LHC, Higgs mass, Higgs couplings, Xenon100

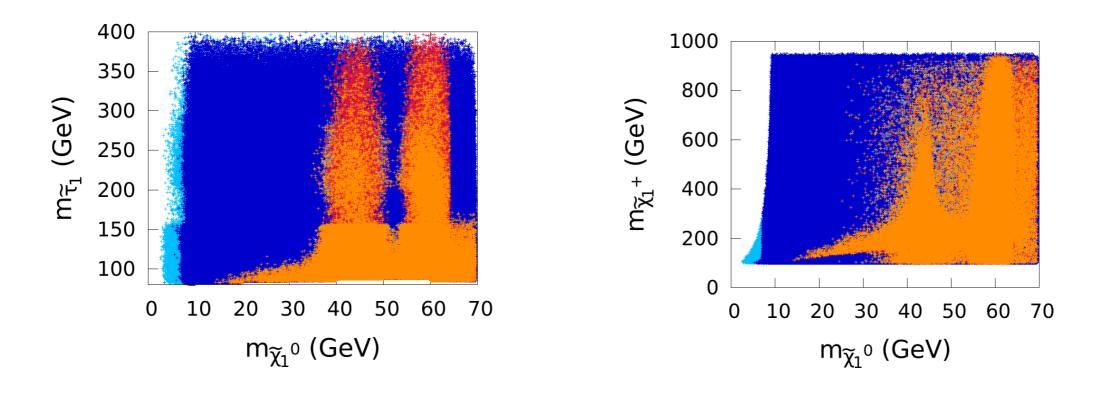


Basic constraints Higgs couplings fits



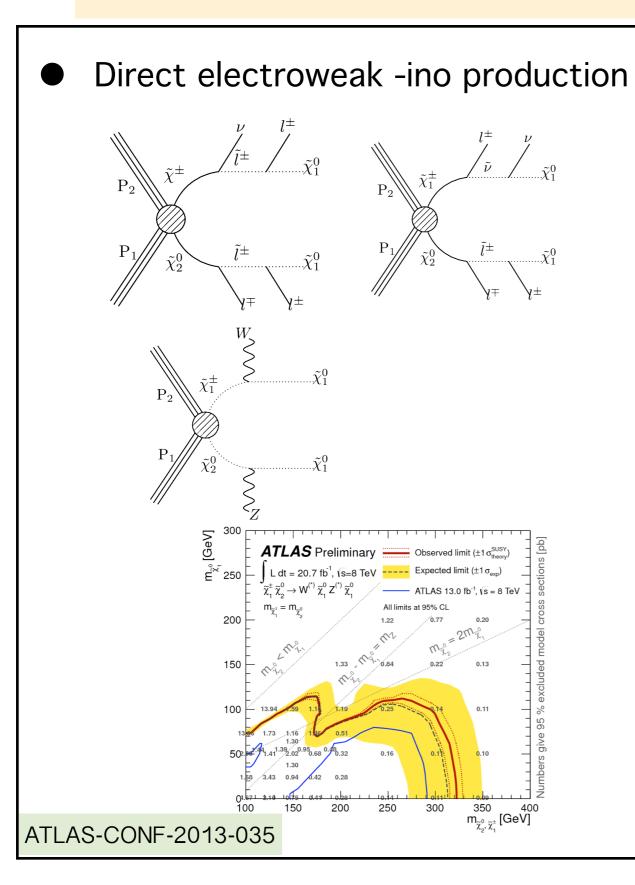
LHC results + upper limit of relic LHC results + exact relic

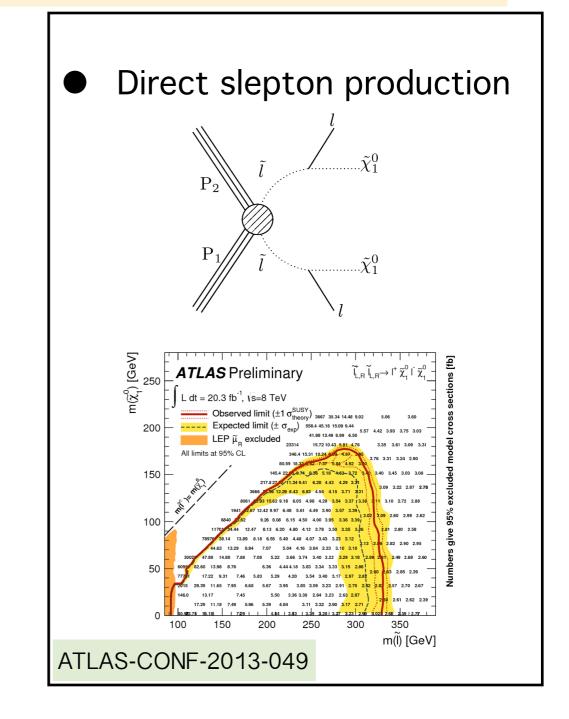
DM < 35 GeV associated with light stau + light chargino



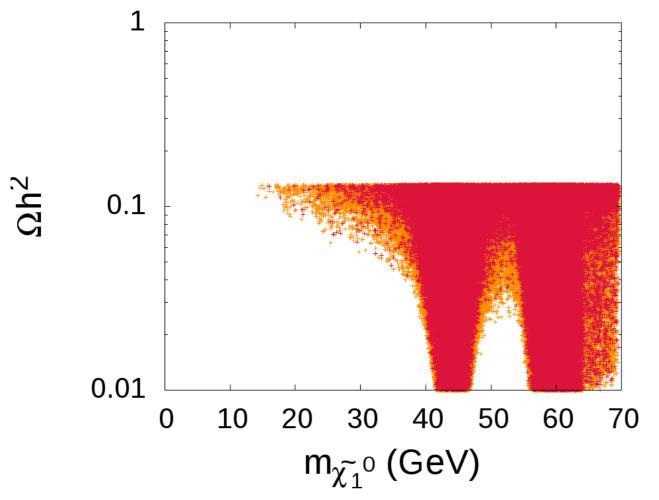
- LHC searches put constraints on light electroweak -ino and DMIEpt35 product associated with light sparticles : light stau + light chargino
- ATLAS and CNS have started to proble electroweakino and sleptons

LHC searches





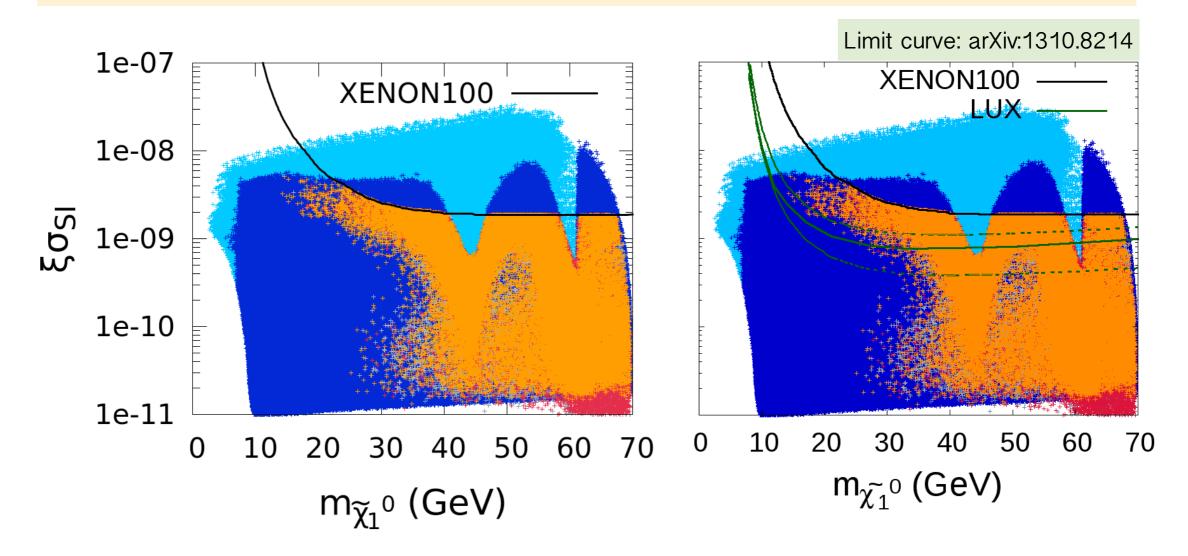
Applying SModelS



All points passing relic density upper limits Points excluded by the LHC limits

- SMS results used from ATLAS-CONF-2013-049, CMS-PAS-SUS-12-022, ATLAS-CONF-2013-035
- Density of points reduced LHC SMS results do rule out some scenarios
- In general light neutralino still possible

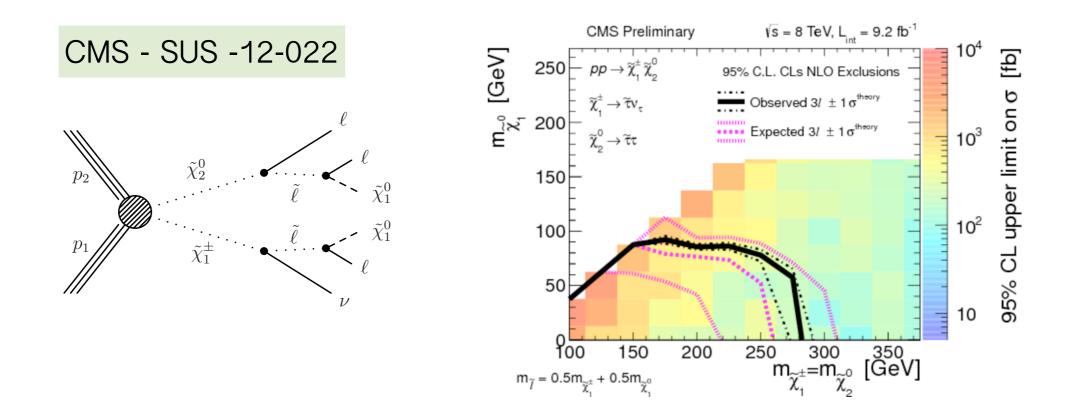
LUX limits



- Neutralino DD CS is driven by higgsino component, suppressed when LSP has small higgsino component
- LUX disfavors the light neutralino DM region we had identified to be viable

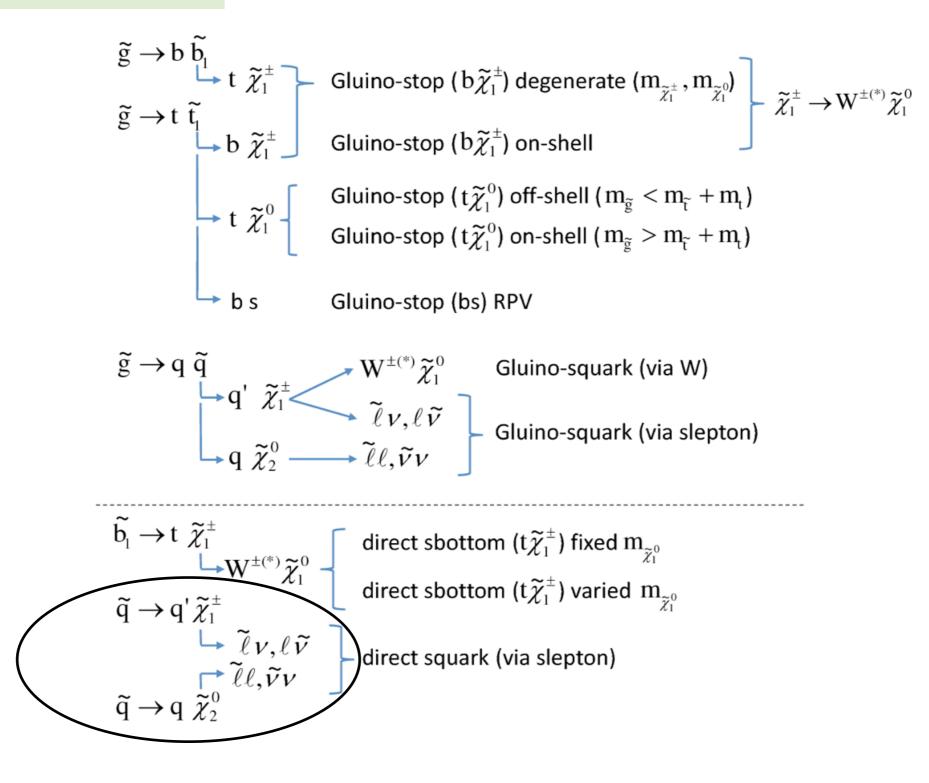
Basic constraints Higgs couplings fits LHC results + upper limit of relic LHC results + exact relic

- SMS approach is not perfect yet
- Not all SMS topologies are present
- It is not always possible to use experimental SMS results, sometimes the results have a coarse grid or in case of a one step decay, only one mass slice is given

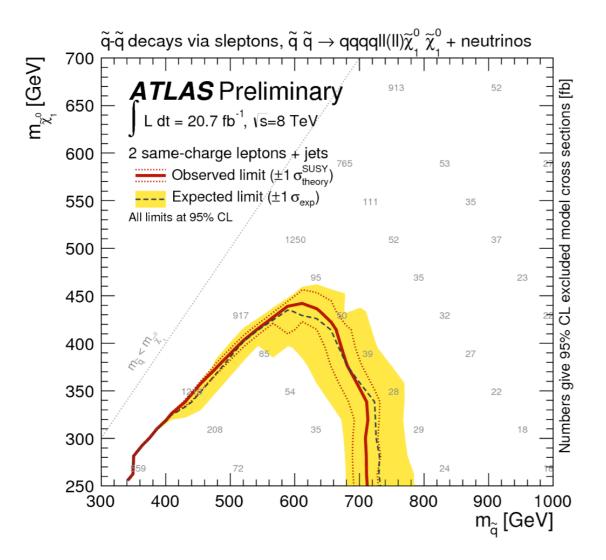


- Results presented are not always usable
- My nightmare SMS analysis: ATLAS-CONF-2013-007
 - Involves topologies with more than four SUSY particles
 - Plots often include strong assumptions on the masses involved
 - Binning is not uniform

ATLAS-CONF-2013-007



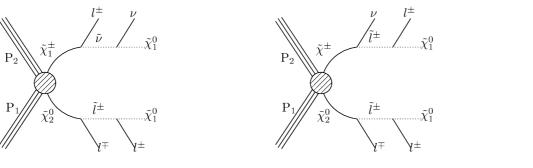
ATLAS-CONF-2013-007

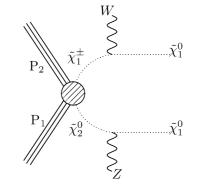


- Two of the four masses fixed
- Only democratic slepton decays
- Irregular binning (less severe)

 In principle several topologies can contribute to the same final state with different efficiencies

Tri-lepton final state: ATLAS-CONF-2013-035



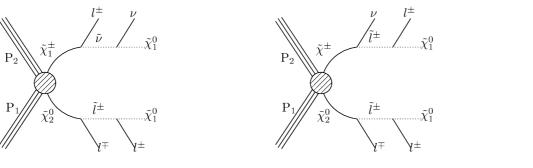


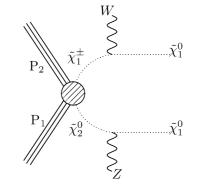
 In principle several topologies can contribute to the same final state with different efficiencies

Tri-lepton final state: ATLAS-CONF-2013-035 Real life $\epsilon_1 \xrightarrow[P_2]{\tilde{\chi}_1^{\pm}} \xrightarrow{\tilde{\chi}_1^{\pm}} + \epsilon_2 \xrightarrow[P_1]{\tilde{\chi}_2^{\pm}} \xrightarrow{\tilde{\chi}_1^{\pm}} + \epsilon_3 \xrightarrow[P_1]{\tilde{\chi}_2^{\pm}} \xrightarrow{\tilde{\chi}_1^{\pm}} + \epsilon_3$

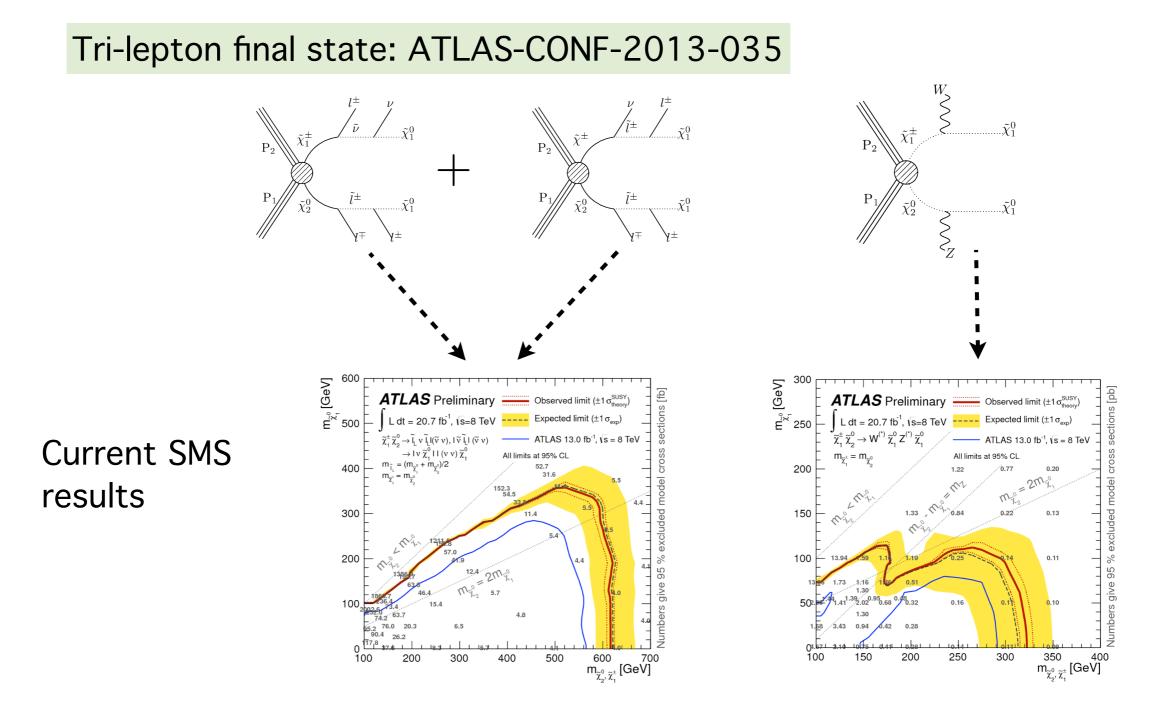
 In principle several topologies can contribute to the same final state with different efficiencies

Tri-lepton final state: ATLAS-CONF-2013-035





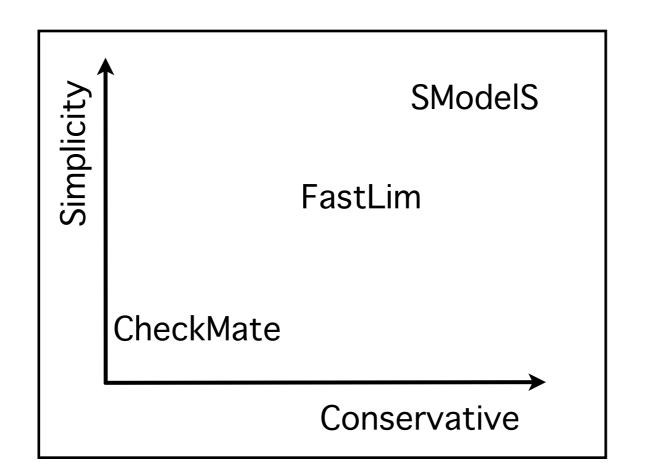
 In principle several topologies can contribute to the same final state with different efficiencies



- To utilize this approach one needs to develop efficiency maps for each analysis and for each topology that can potentially contribute to the final state
- Need to re-implement the analysis requires manpower and availability of information from experimental collaborations in a systematic manner
- FastLim is developing the efficiency maps, aim is to reconstruct the number of events for each signal region
- SModelS is also capable of supporting efficiency maps approach, and in future we might consider exploiting this feature

- Generally the development will be slow a community effort to contribute to the efficiency maps is underway, this effort should also make the re-implemented analysis publicly available
- Also need to develop a reasonable likelihood in order to be able to combine several signal regions - e.g. CMS-SUS-PAS-13-011 has 16 signal regions

- It will be difficult to tackle long cascade decays with SMS approach
- A completely different approach is being taken by checkMATE
- The tool identifies the most sensitive topology and then tests it via Monte Carlo simulation



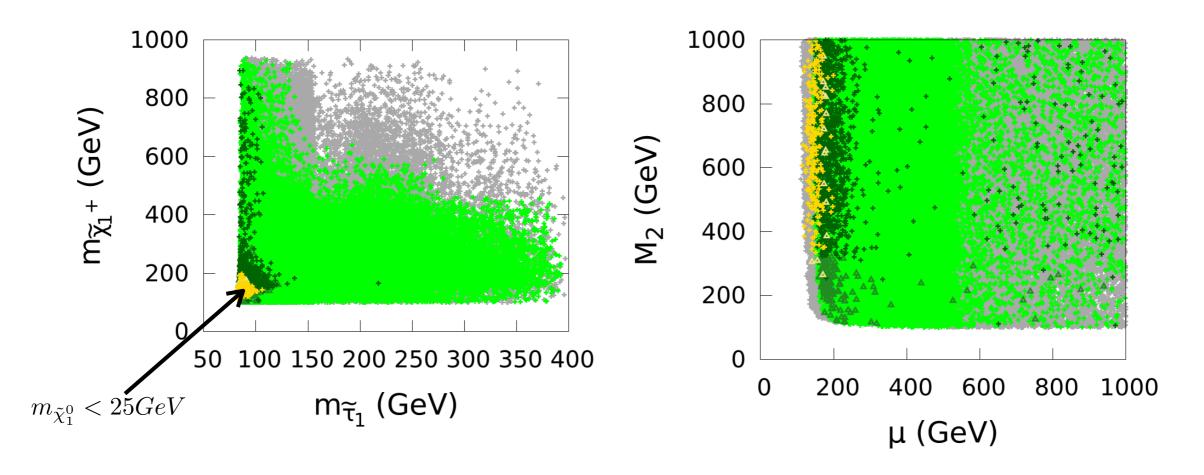
CheckMATEarXiv:1312.2591

Conclusions

- SMS results are a good way to test BSM theories and can have a good constraining power
- SModelS is designed to utilize this power and constrain BSM scenarios
- The formalism of the code is generic and can be applied to any BSM spectra for which SMS results are applicable
- Currently, the code can handle scenarios with Z2 parity
- It contains the most comprehensive database of SMS SUSY results
- It can be used in order to understand the features of parameter space under consideration, it can also be used to study viability of an interesting BSM scenario
- Stay tuned applying LHC searches to your favorite BSM model is being made easy!

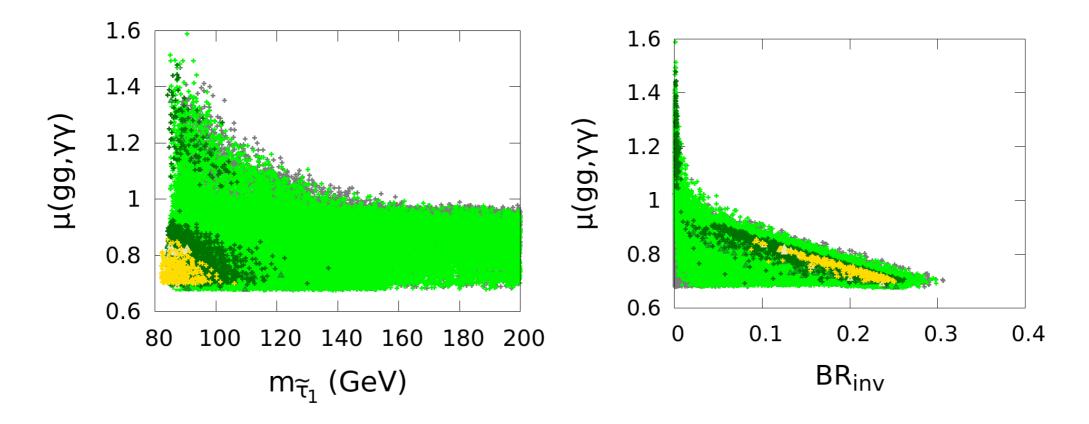


Applying SModelS



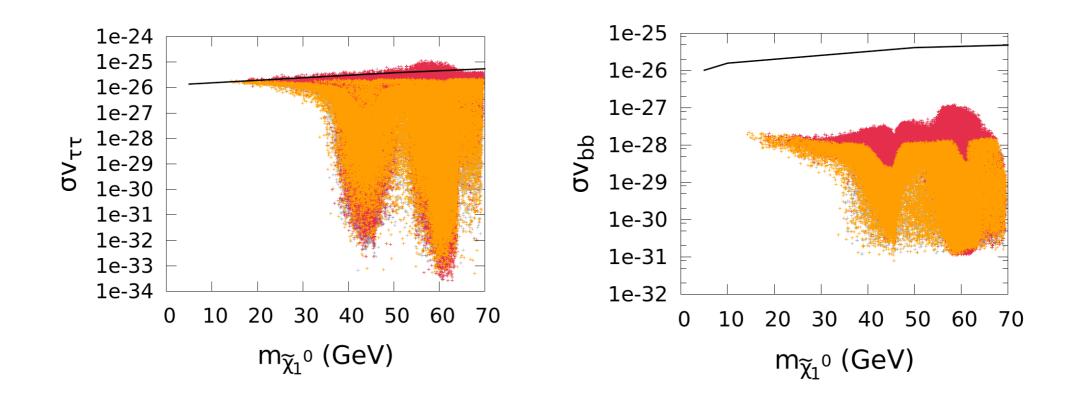
- Mass slices 10-15; 25-35; 35-50; 50-60 GeV
- Density of points reduced LHC does rule out some scenarios
- In general light neutralino still possible

Higgs signal strengths



- Mass slices 10-15; 25-35; 35-50; 50-60 GeV
- Light neutralino and light stau lead to modifications
- Heavily mixed stau increases diphoton rate
- Lightest neutralino associated with some invisible Higgs decays
- At 14 TeV with ZH -> invisible, better sensitivity expected

Indirect detection limits



- We test for FERMI-LAT limits photons produced from DM annihilation in dwarf spheroidal galaxies in bbar or tautau channel
- Large fraction of LSP < 30 GeV points are several orders of magnitude below the limit

Why not use monojet channel?

- Direct LSP production probed via monojet signature at the LHC
- Limits given on the spin-independent interactions of DM
- Limits applicable for models involving heavy mediators
- Not applicable for MSSM since the mediators are not heavy enough

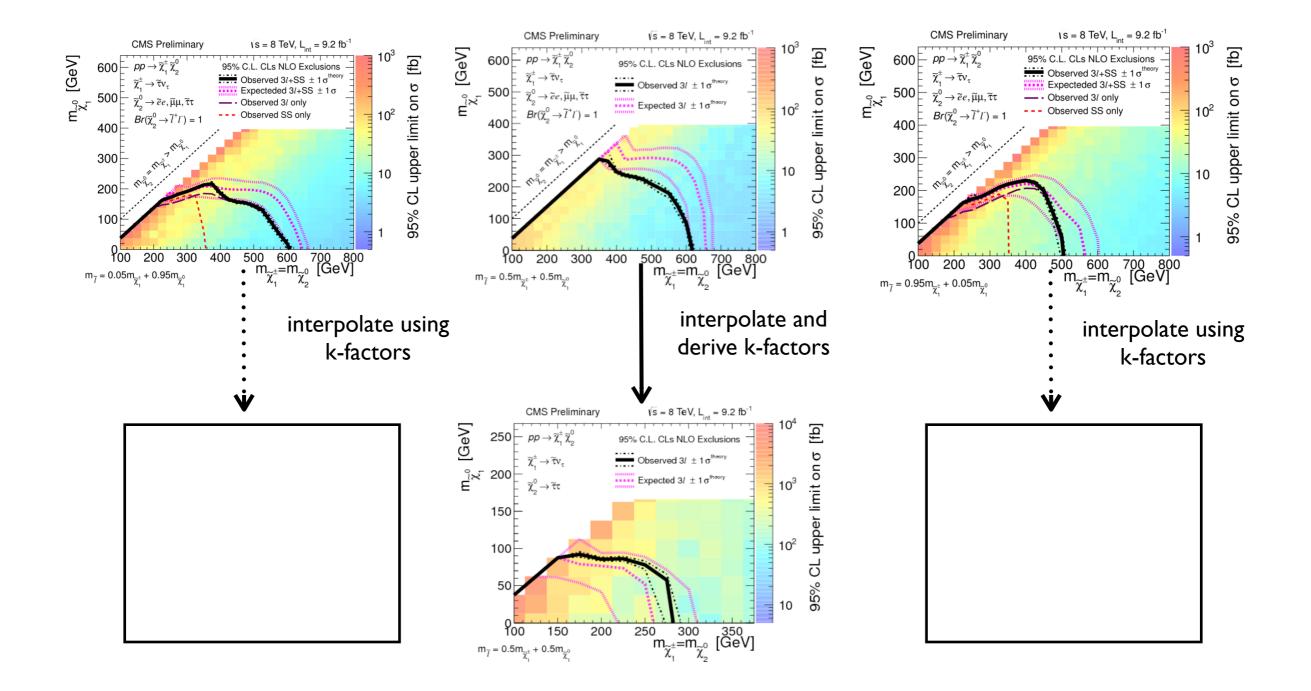
Light neutralino DM - scan details

$ aneta\ M_{A^0}$	[5, 50] [100, 1000]	$ \begin{array}{ c } M_{L_3} \\ M_{R_3} \end{array} $	[70, 500] [70, 500]
M_1	[10, 70]	A_{τ}	[-1000, 1000]
M_2	[100, 1000]	M_{L_1}	[100, 500]
μ	[100, 1000]	M_{R_1}	[100, 500]

	> 100 C - V
LEP limits	$m_{\tilde{\chi}_1^{\pm}} > 100 \text{ GeV}$
	$m_{\tilde{\tau}_1} > 84 - 88 \text{ GeV} (\text{depending on } m_{\tilde{\chi}_1^0})$
	$\sigma(e^+e^- \to \tilde{\chi}^0_{2,3} \tilde{\chi}^0_1 \to Z^{(*)}(\to q\bar{q}) \tilde{\chi}^0_1) \lesssim 0.05 \text{ pb}$
invisible Z decay	$\Gamma_{Z \to \tilde{\chi}_1^0 \tilde{\chi}_1^0} < 3 \text{ MeV}$
μ magnetic moment	$\Delta a_{\mu} < 4.5 \times 10^{-9}$
flavor constraints	$\mathrm{BR}(b \to s\gamma) \in [3.03, 4.07] \times 10^{-4}$
	$BR(B_s \to \mu^+ \mu^-) \in [1.5, 4.3] \times 10^{-9}$
Higgs mass	$m_{h^0} \in [122.5, 128.5] \text{ GeV}$
$A^0, H^0 \to \tau^+ \tau^-$	CMS results for $\mathcal{L} = 17 \text{ fb}^{-1}$, m_h^{max} scenario
Higgs couplings	ATLAS, CMS and Tevatron global fit, see text
relic density	$\Omega h^2 < 0.131 \text{ or } \Omega h^2 \in [0.107, 0.131]$
direct detection	XENON100 upper limit
indirect detection	Fermi-LAT bound on gamma rays from dSphs
$pp \to \tilde{\chi}_2^0 \tilde{\chi}_1^\pm$	Simplified Models Spectra approach, see text
$pp \to \tilde{\ell}^+ \tilde{\ell}^-$	
	1

Tau dominated scenario

 For topologies involving an intermediate particles, three mass slices are given. We interpolate over these slices



- pMSSM scan over 6 parameters
- Gaugino masses obey GUT relation
- Flavor constraints, invisible Z width, Higgs mass, LEP limits imposed
- Scan ranges

M_2	μ	$\tan\beta$	$M_{\tilde{q}}$	$M_{\tilde{Q}_3}$	$M_{\tilde{D}_3}$	$M_{\tilde{U}_3}$	A_t	A_b
0.1-1	0.1–1	3-60	0.1 - 5	0-2	0-2	0-2	$[1,3]\max(M_{\tilde{Q}_3}, M_{\tilde{Q}_3})$	± 1

