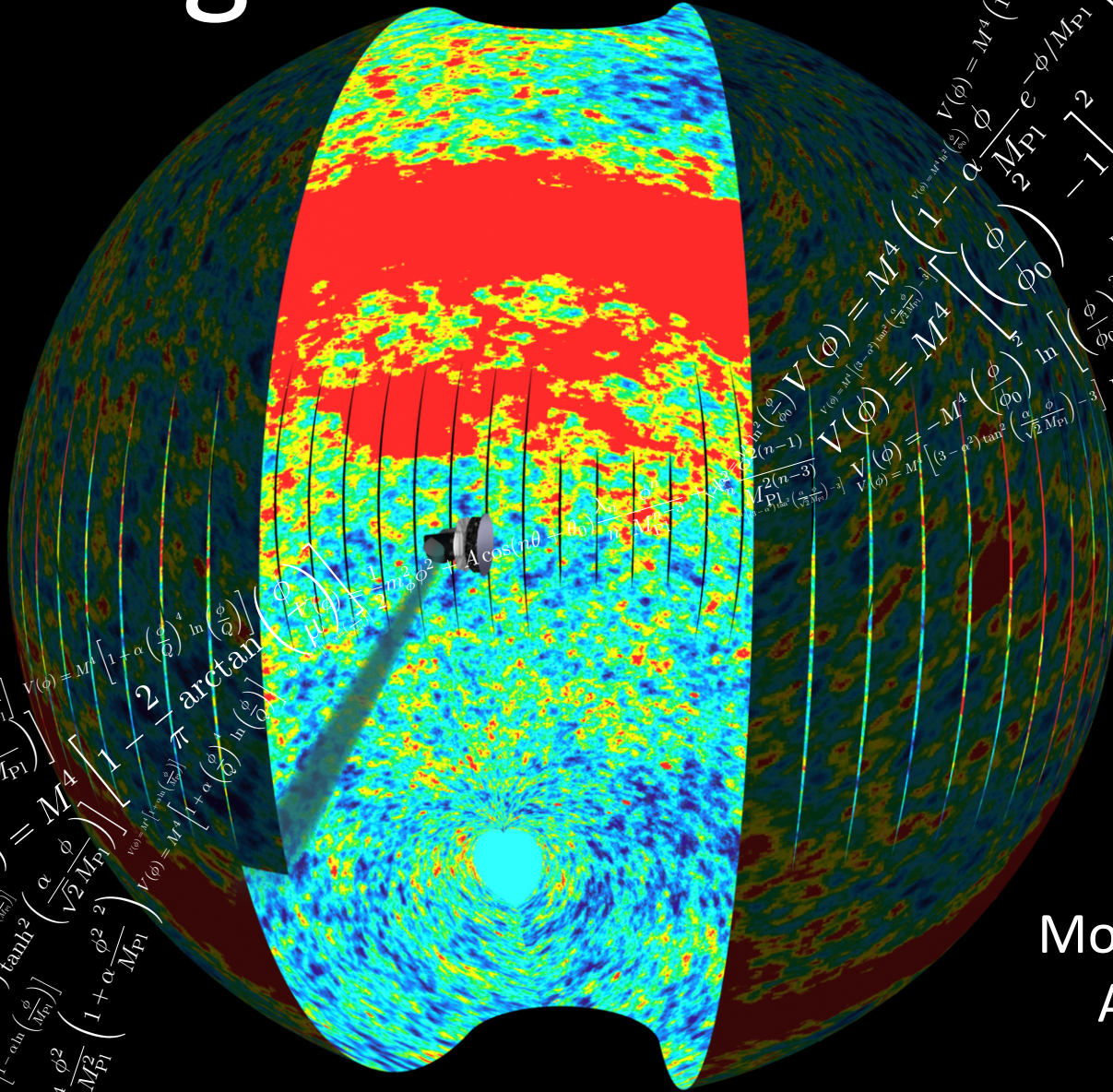


# Cosmological Inflation

to the test



Montpellier L2C  
April 7<sup>th</sup> 2016

Vincent Vennin, ICG Portsmouth

# Outline

- Inflation: Where do we Stand?
- Bayesian Model Comparison for Single-Field Models
- Including (and constraining) Reheating
- Adding a Light Scalar Field

## Collaborators:

Jérôme [Martin](#) (IAP), Christophe [Ringeval](#) (Louvain U. CP3), and Roberto [Trotta](#) (Imp.Coll)  
David [Wands](#) (ICG), Kazuya [Koyama](#) (ICG)

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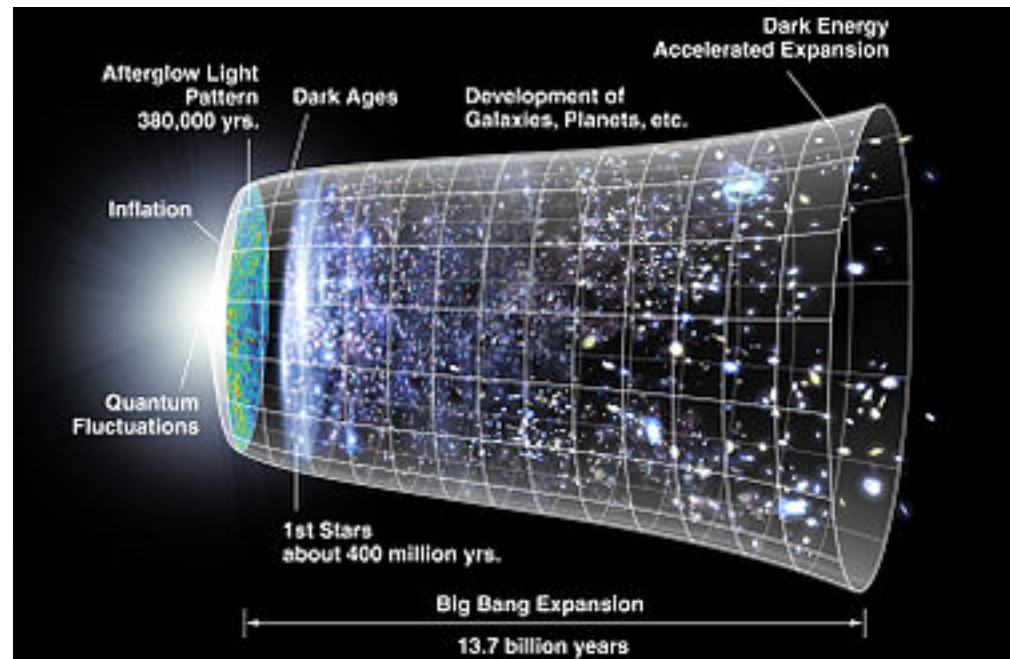
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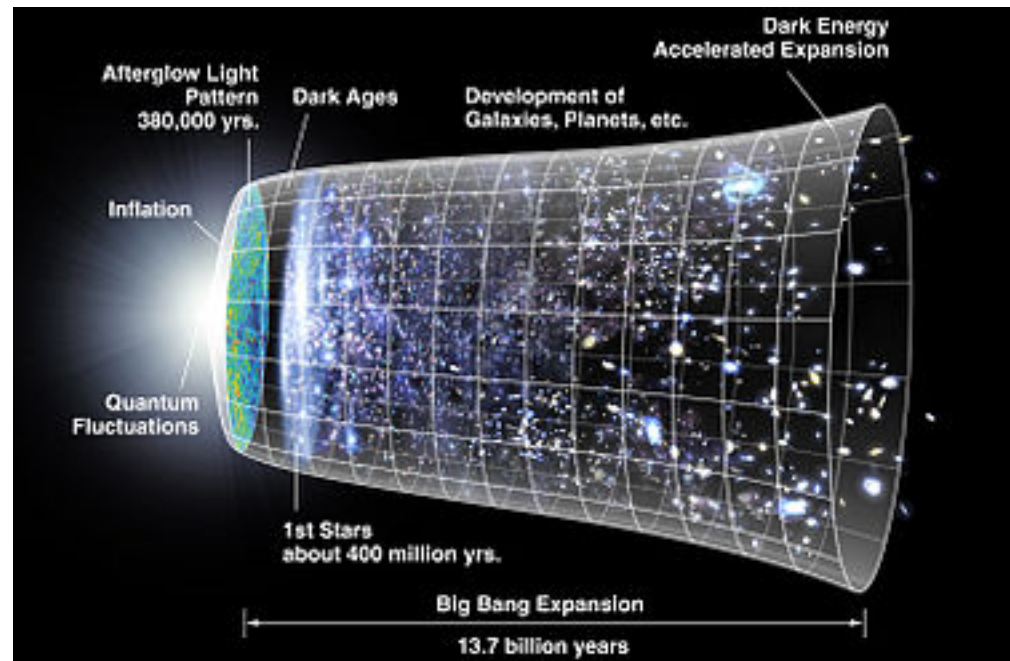


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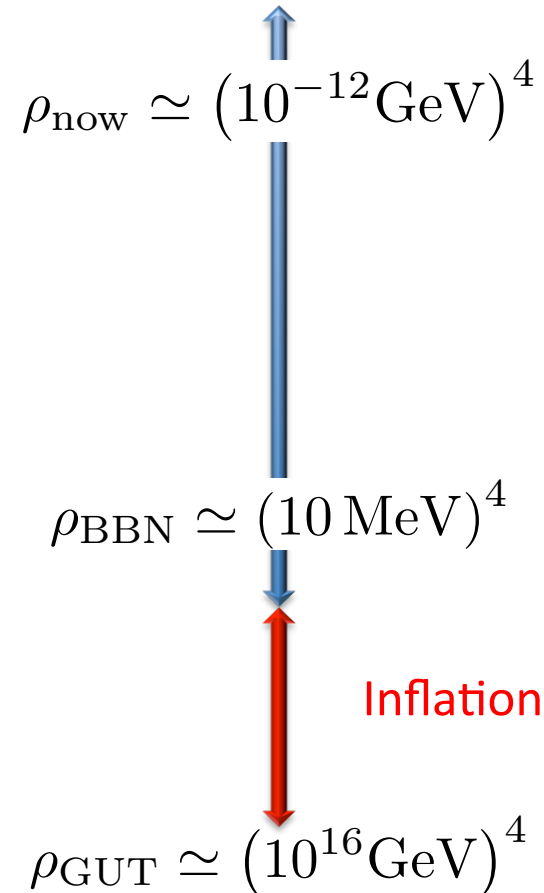
$$ds^2 = -dt^2 + a^2(t) d\vec{x}^2$$



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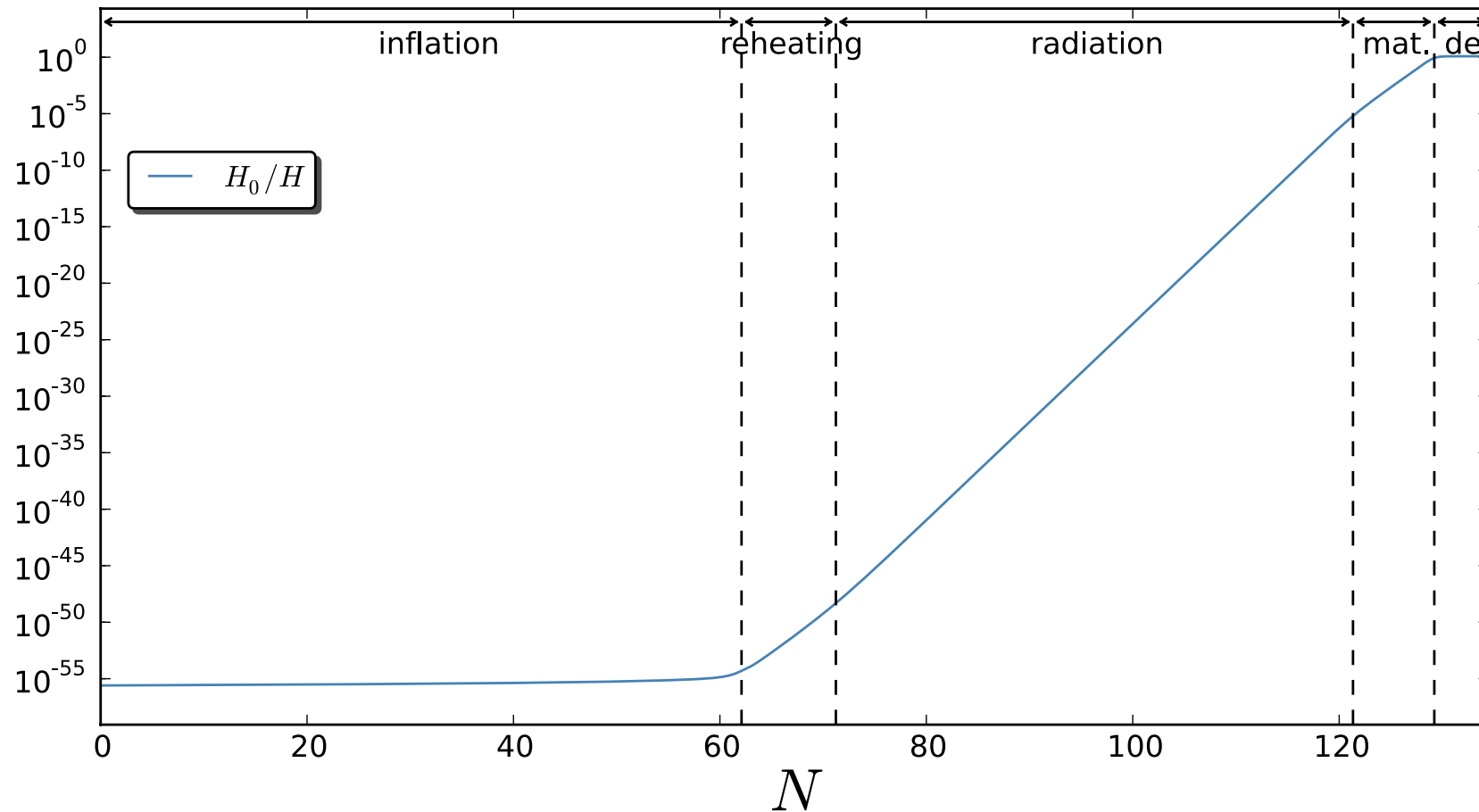
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- Combined with **QM**, accounts for the production of **cosmological perturbations** whose features depend on the underlying inflationary **model**.

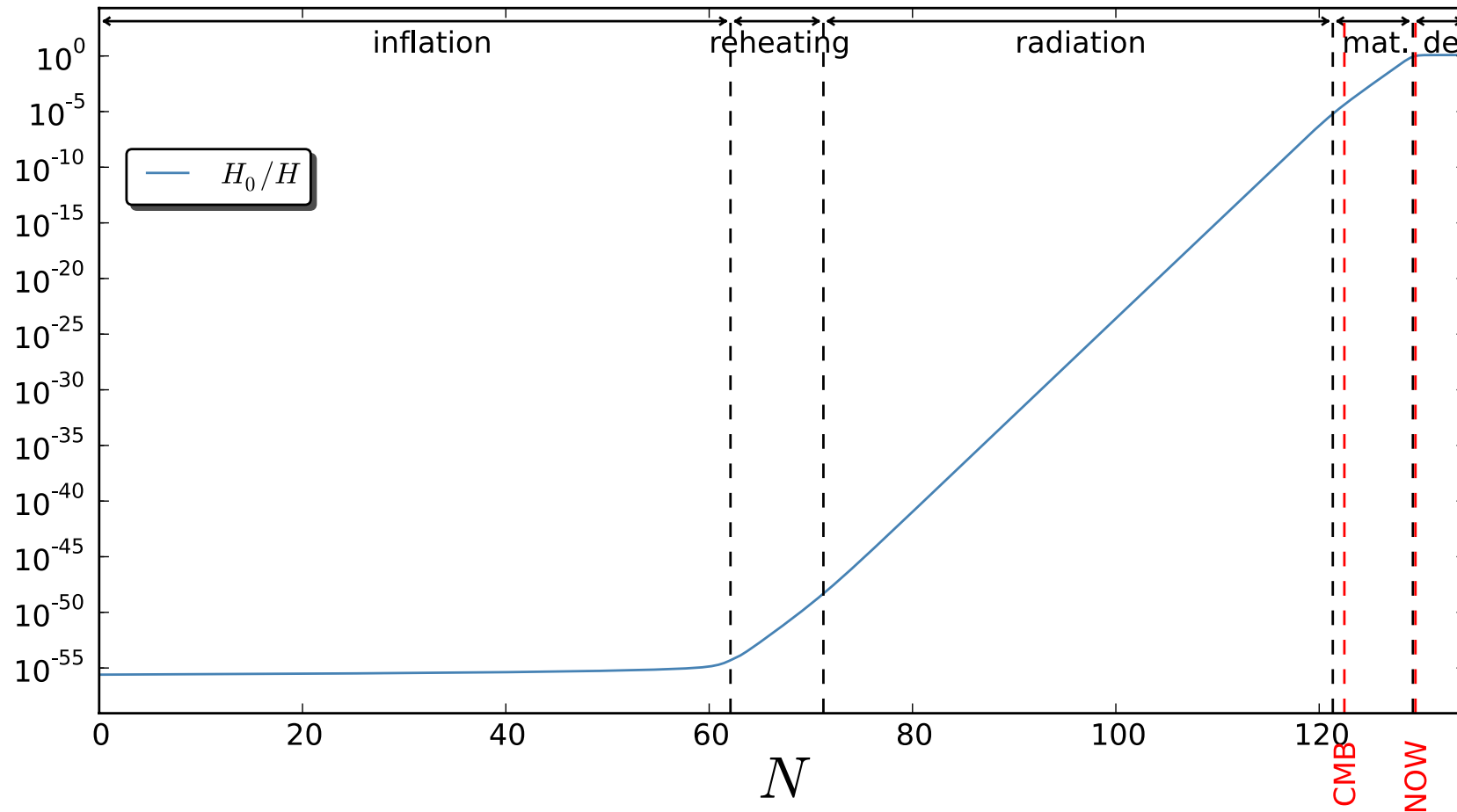
# Cosmological Perturbations

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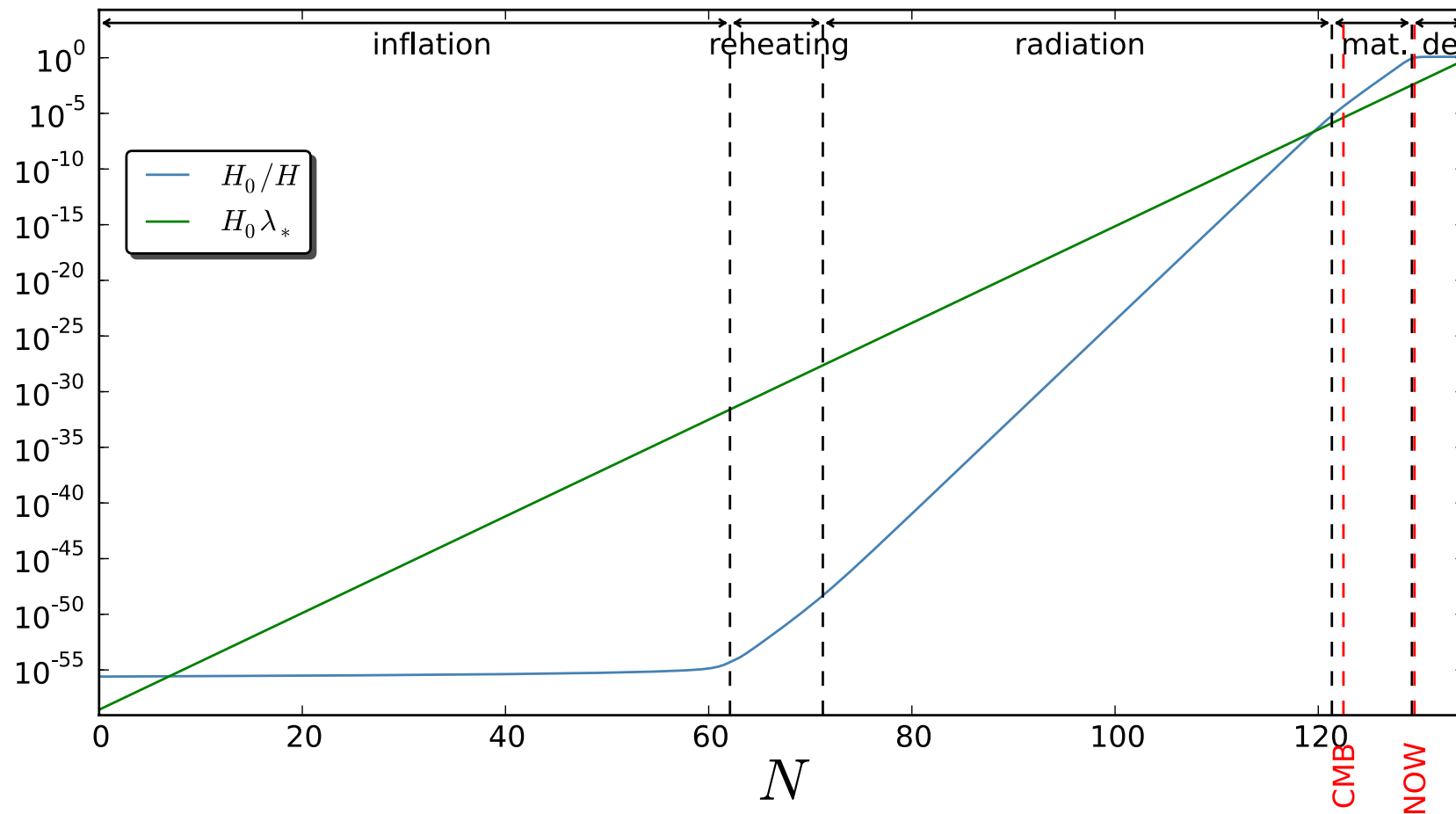
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Coherent, Gaussian, almost scale invariant, adiabatic perturbations

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# Planck Results in brief

Released March 2013, Updated February 2015

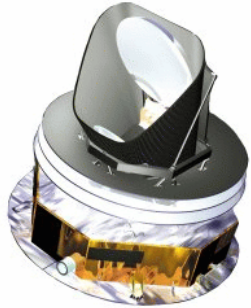
*Planck +...*



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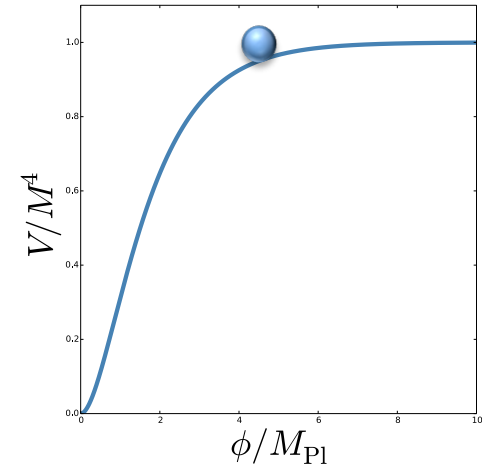
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## •Consequences for Inflationary Models in Particular ...

# Inflationary Observables

of single-field slow-roll models

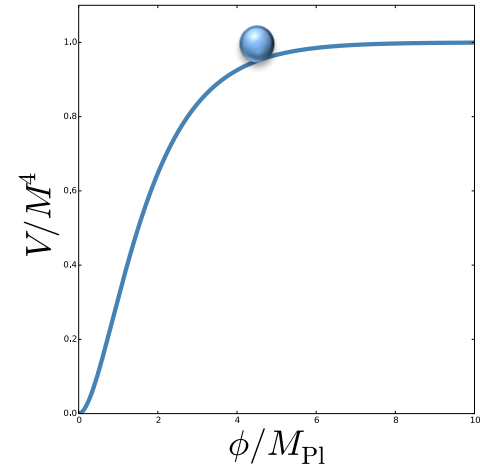


# Inflationary Observables

of single-field slow-roll models

- The slow-roll approximation

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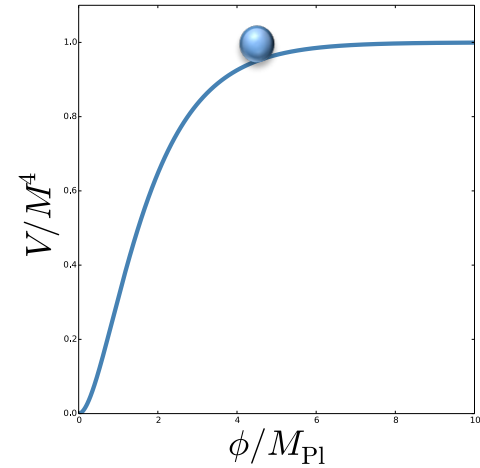
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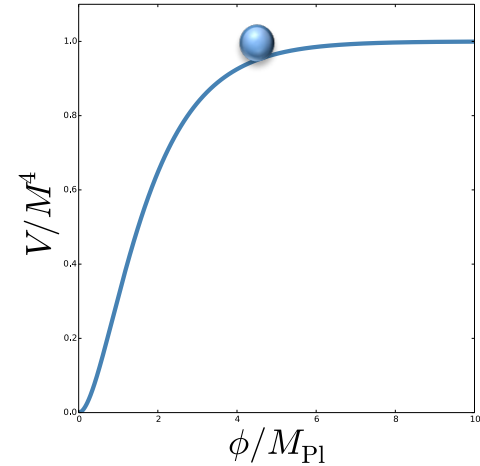
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Slow-Roll hierarchy

$$\epsilon_{n+1} = \frac{1}{\epsilon_n} \frac{d\epsilon_n}{dN}$$



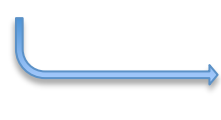
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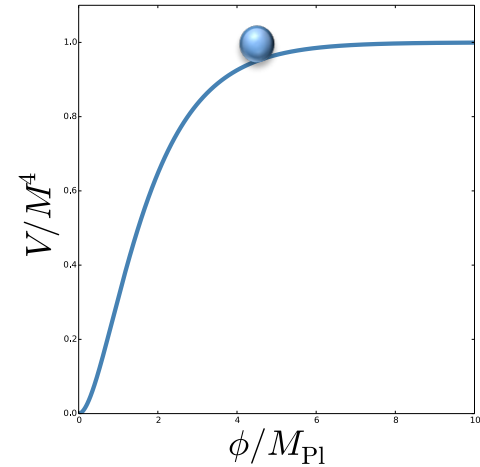
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$$\epsilon_{n+1} = \frac{1}{\epsilon_n} \frac{d\epsilon_n}{dN}$$

$$\epsilon_1 \simeq \frac{1}{2M_{\text{Pl}}^2} \left( \frac{V_\phi}{V} \right)^2$$

$$\epsilon_2 \simeq \frac{2}{M_{\text{Pl}}^2} \left[ \left( \frac{V_\phi}{V} \right)^2 - \frac{V_{\phi\phi}}{V} \right]$$

$\epsilon_3 \simeq \text{etc...}$

# Inflationary Observables

of single-field slow-roll models

Starobinsky (1979)  
Hawking (1982)  
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$$\mathcal{P}_\zeta = \frac{H_*^2}{8\pi^2 M_{\text{Pl}}^2 \epsilon_{1*}} \left[ 1 - 2(C+1)\epsilon_{1*} - C\epsilon_{2*} - (2\epsilon_{1*} + \epsilon_{2*}) \ln \left( \frac{k}{k_*} \right) \right]$$

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$$f_{\text{NL}}^{\text{loc}} = \frac{5}{12} (2\epsilon_{1*} + \epsilon_{2*}) \quad \text{Maldacena (2002)}$$

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$$n_S \simeq 0.96$$

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$r$   $\longrightarrow$  energy scale of inflation

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$n_T$   $\longrightarrow$  consistency relation

$$n_T \simeq -r/8$$

# Model Predictions

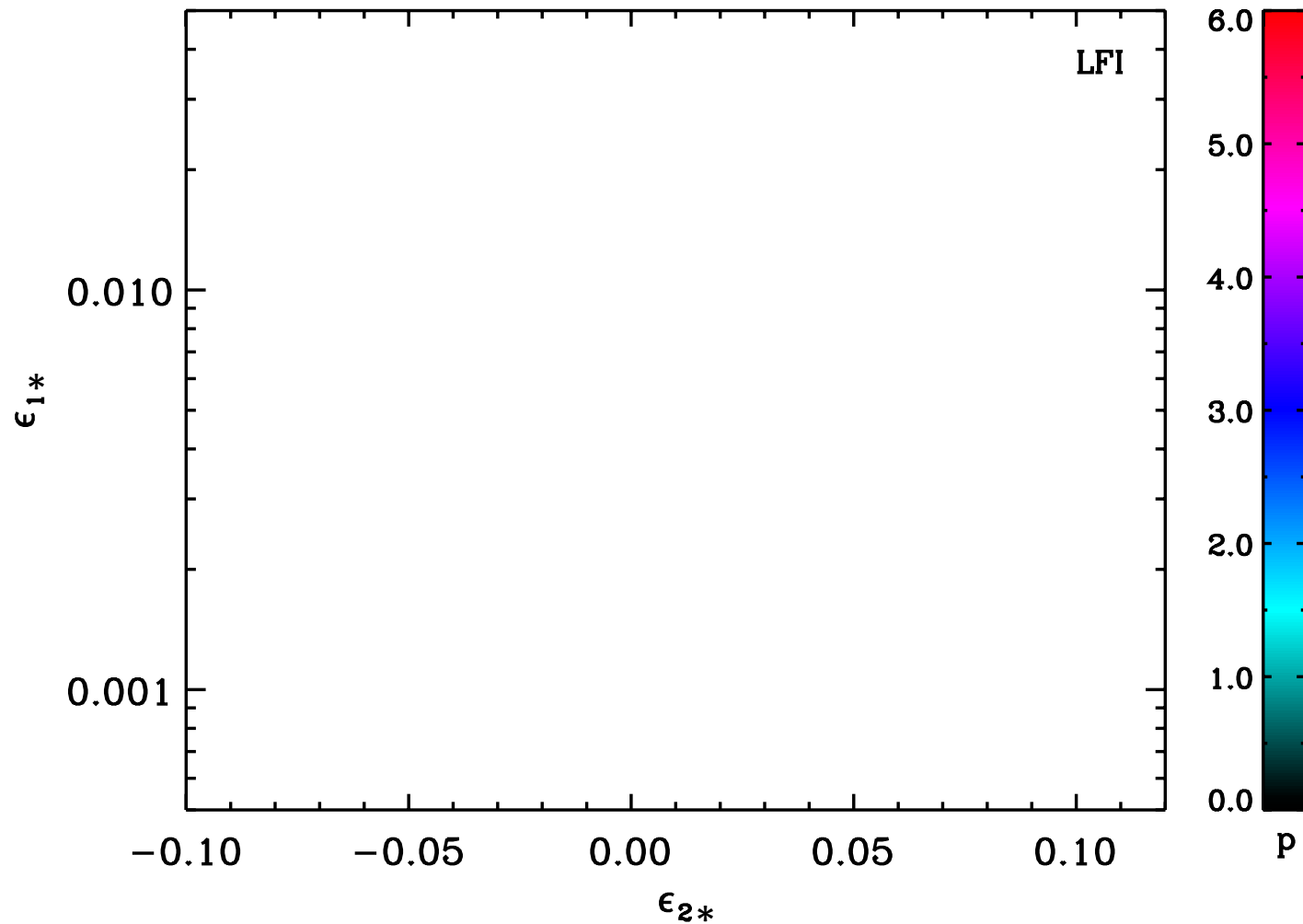
Described in terms of  $\epsilon_{i*}$



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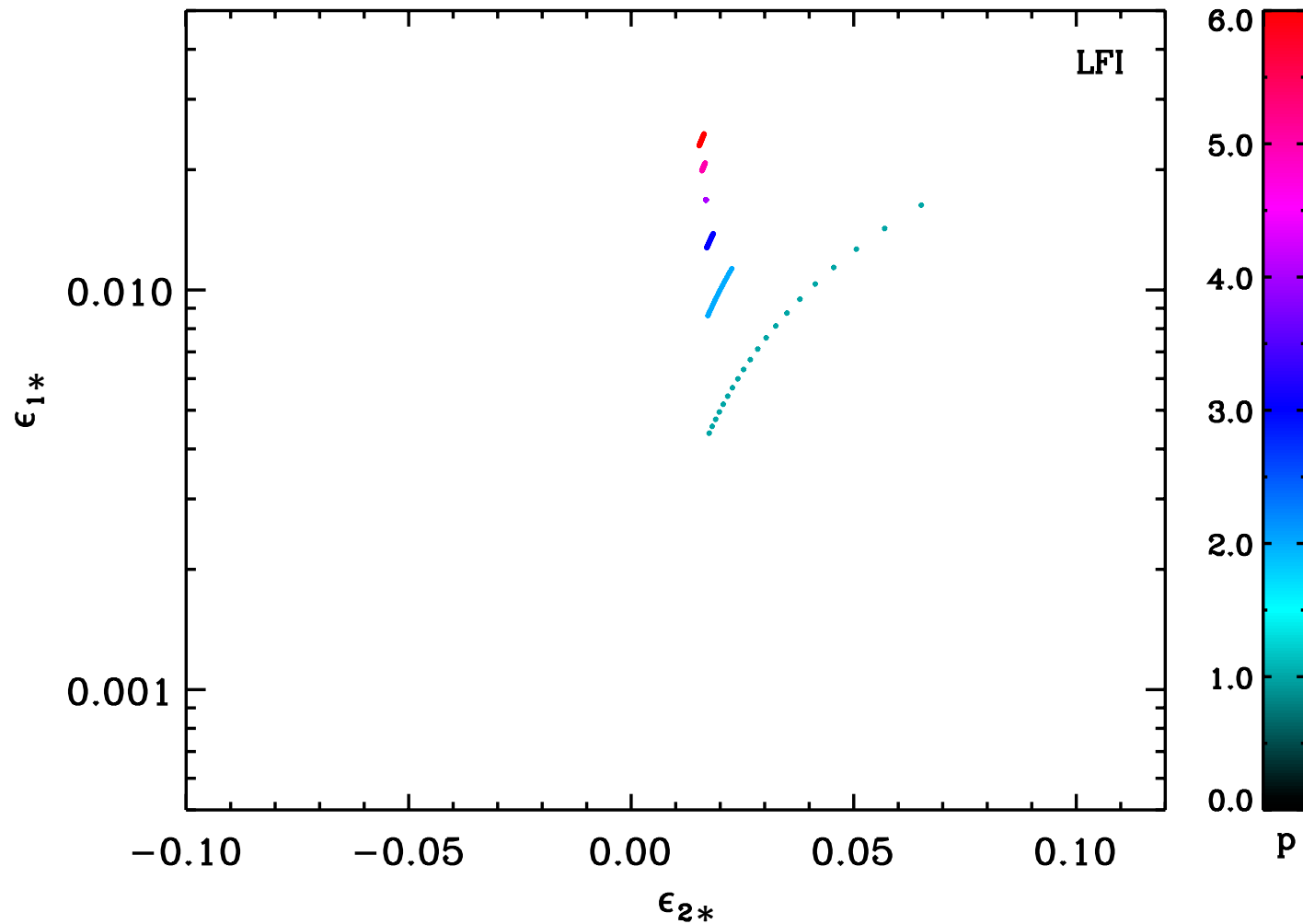
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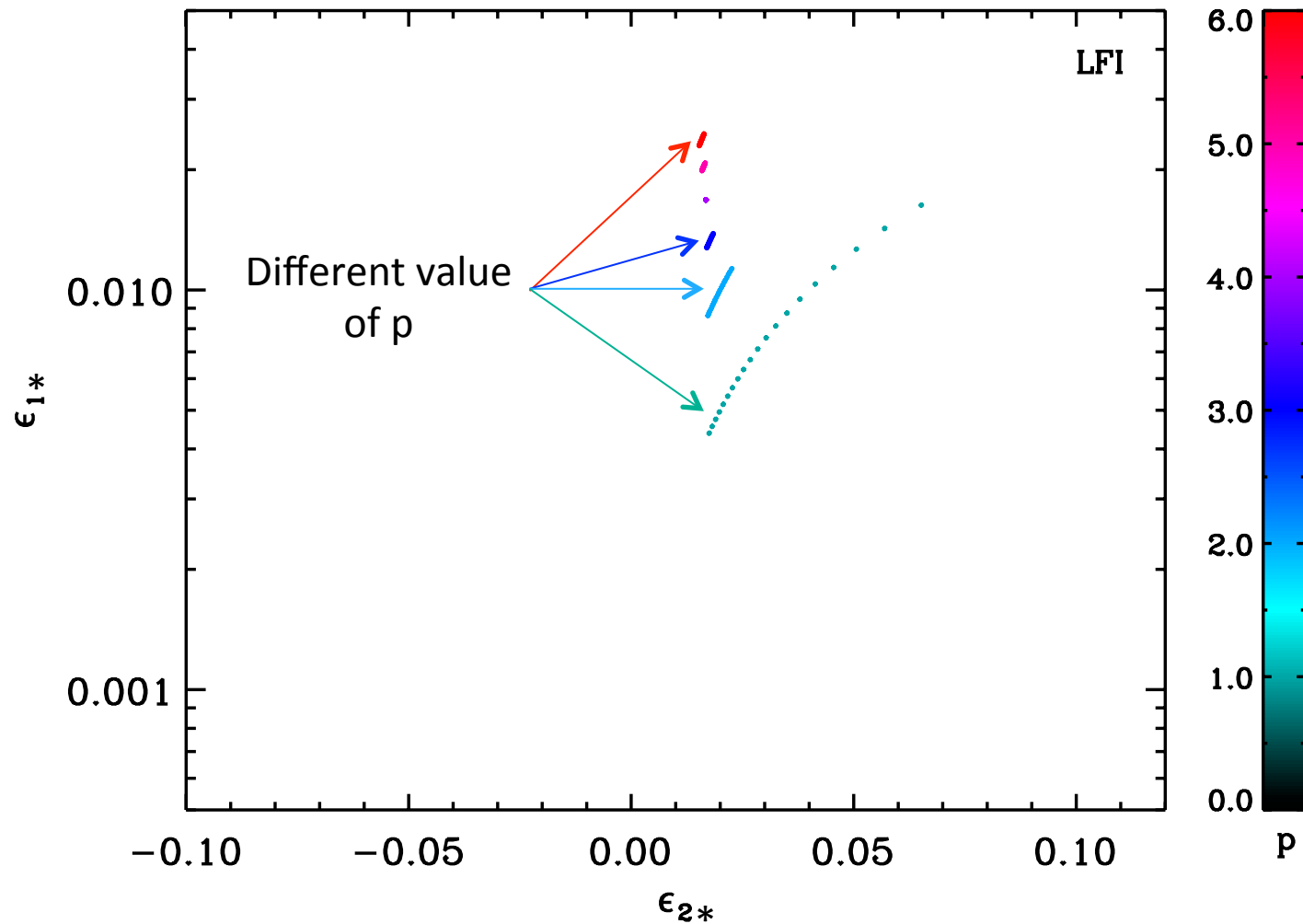
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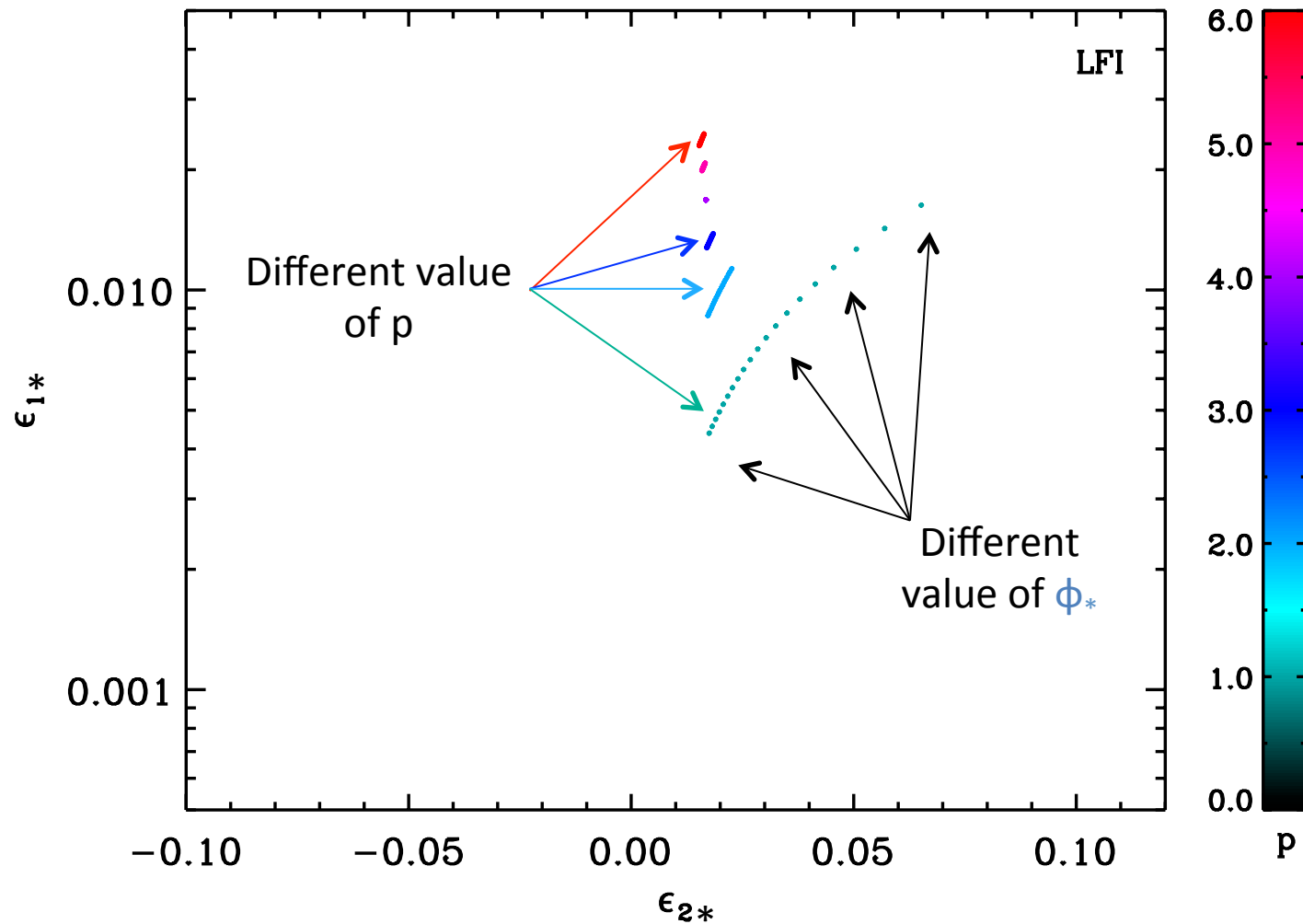
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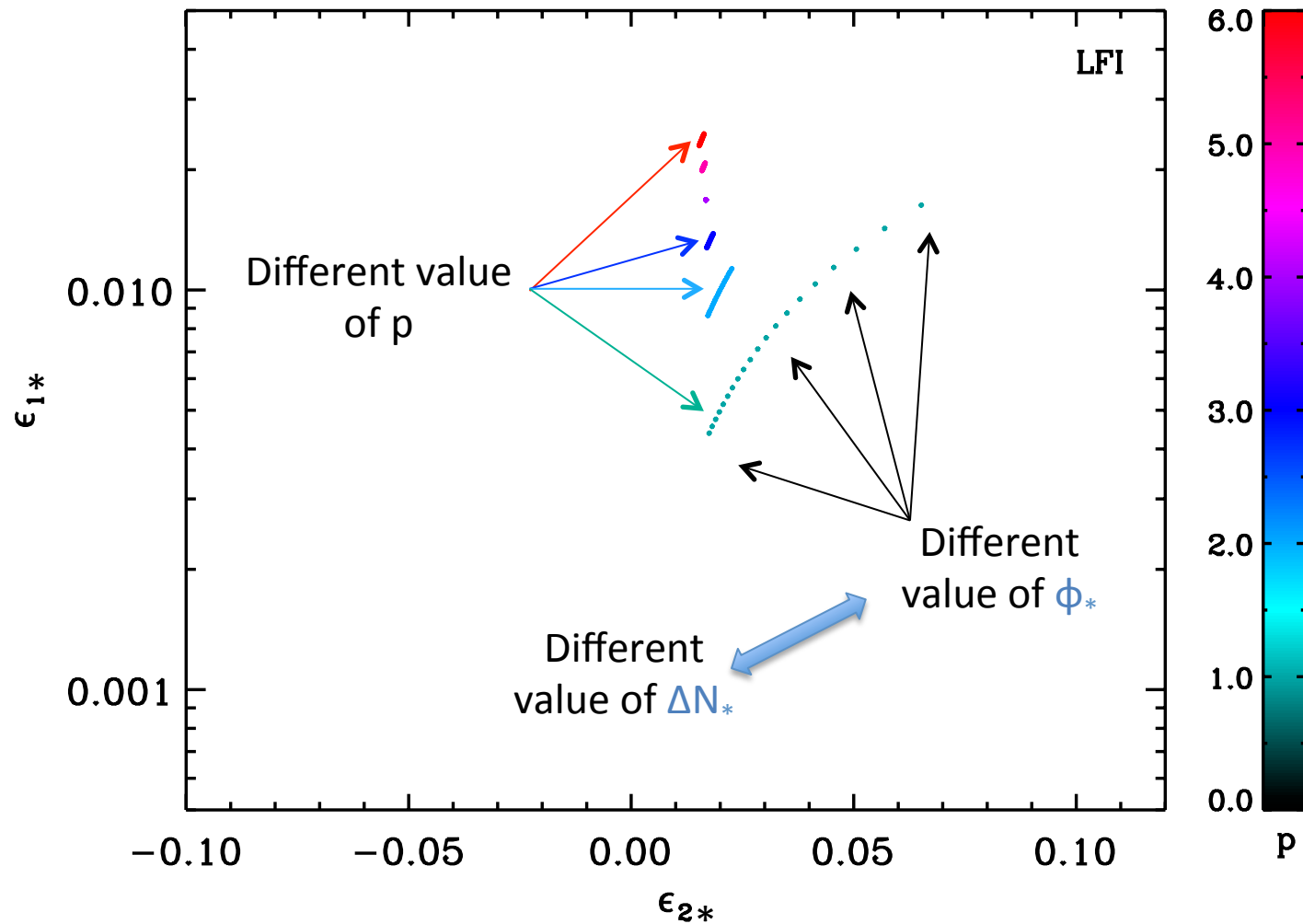
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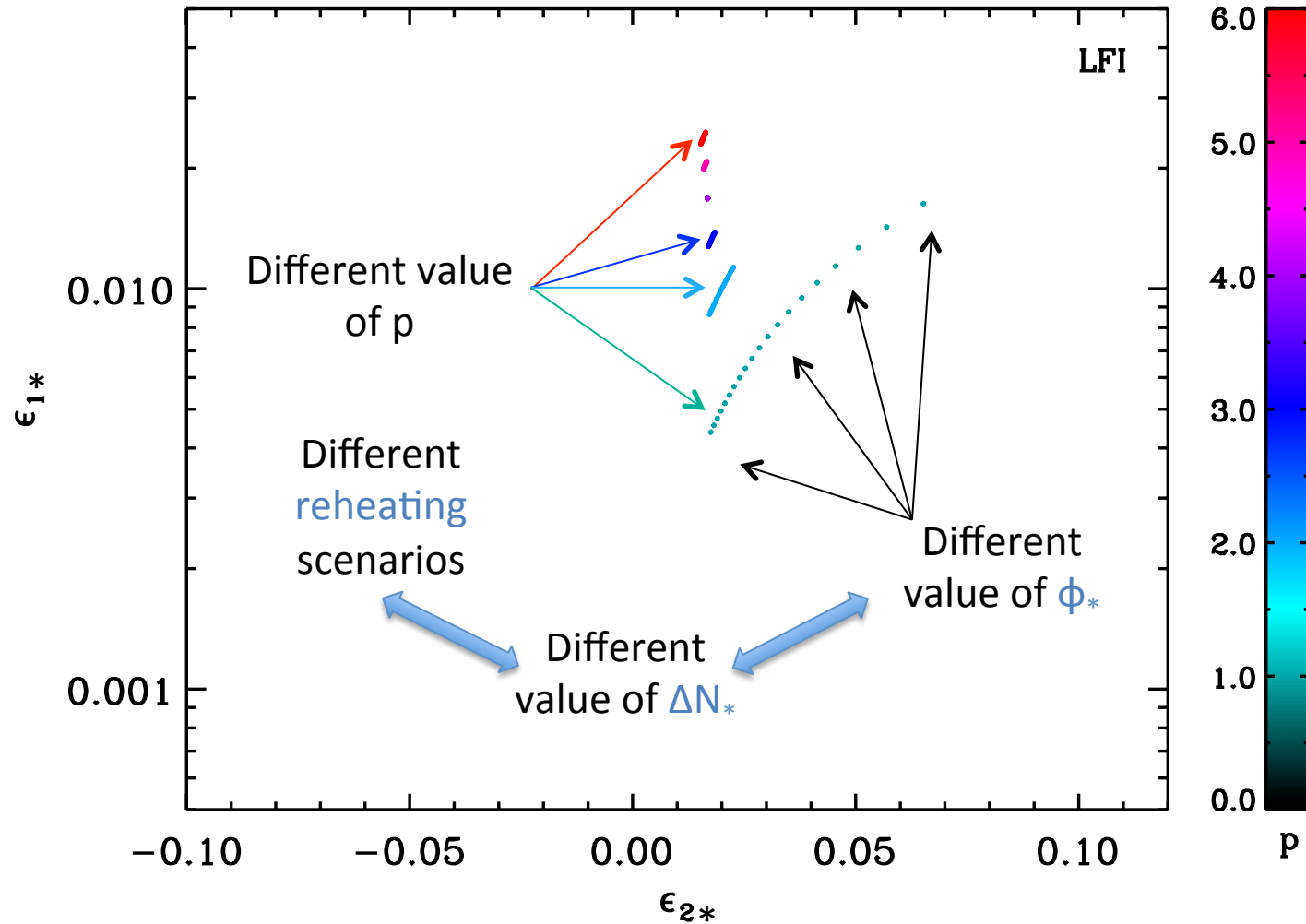
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# Model Predictions

Described in terms of  $\epsilon_{i*}$

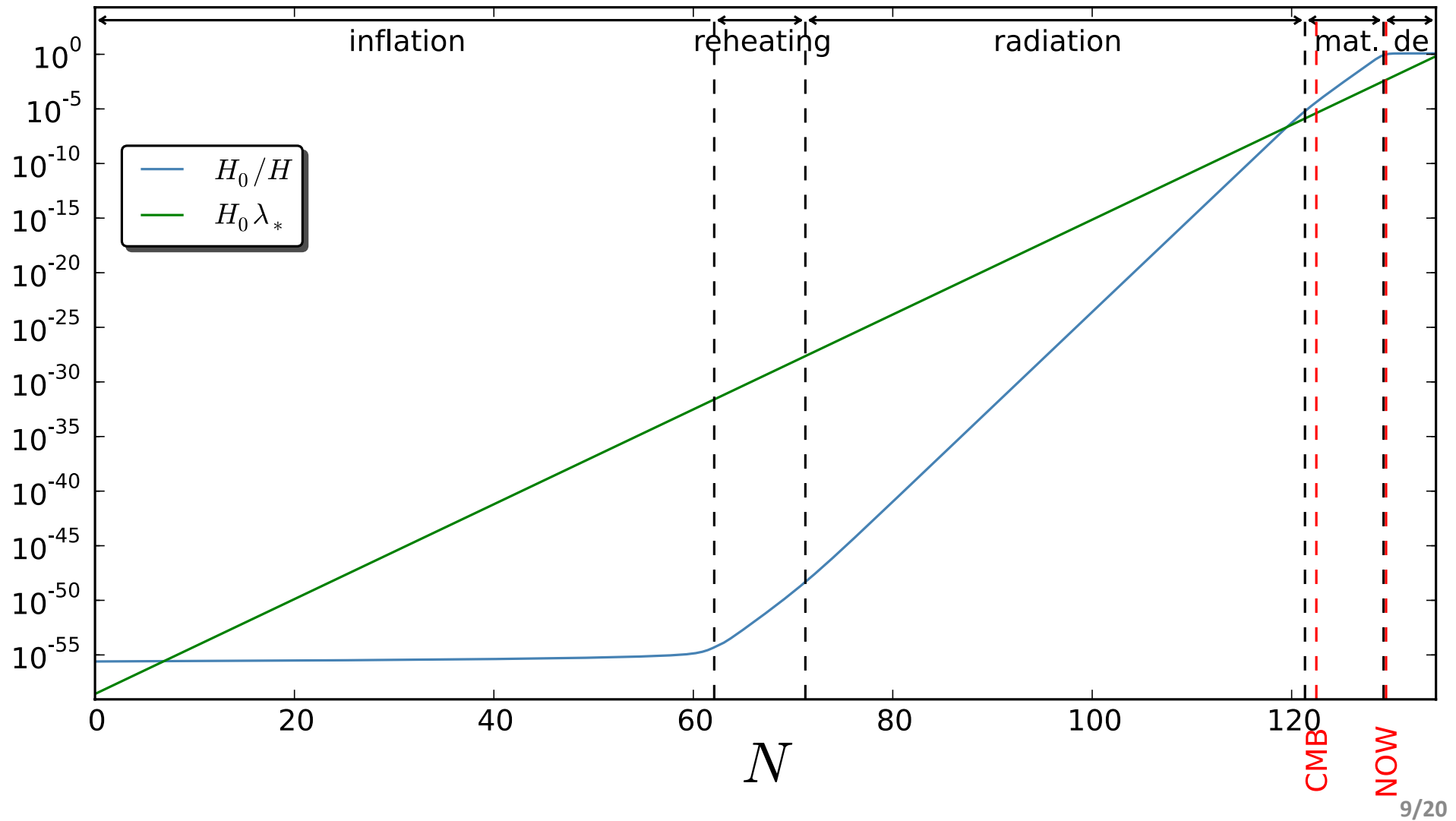
An example: « large field inflation »  $V(\phi) = M^4 \left( \frac{\phi}{M_{\text{Pl}}} \right)^p$



# Role of Reheating

Martin & Ringeval (2010)  
Easter & Peiris (2011)

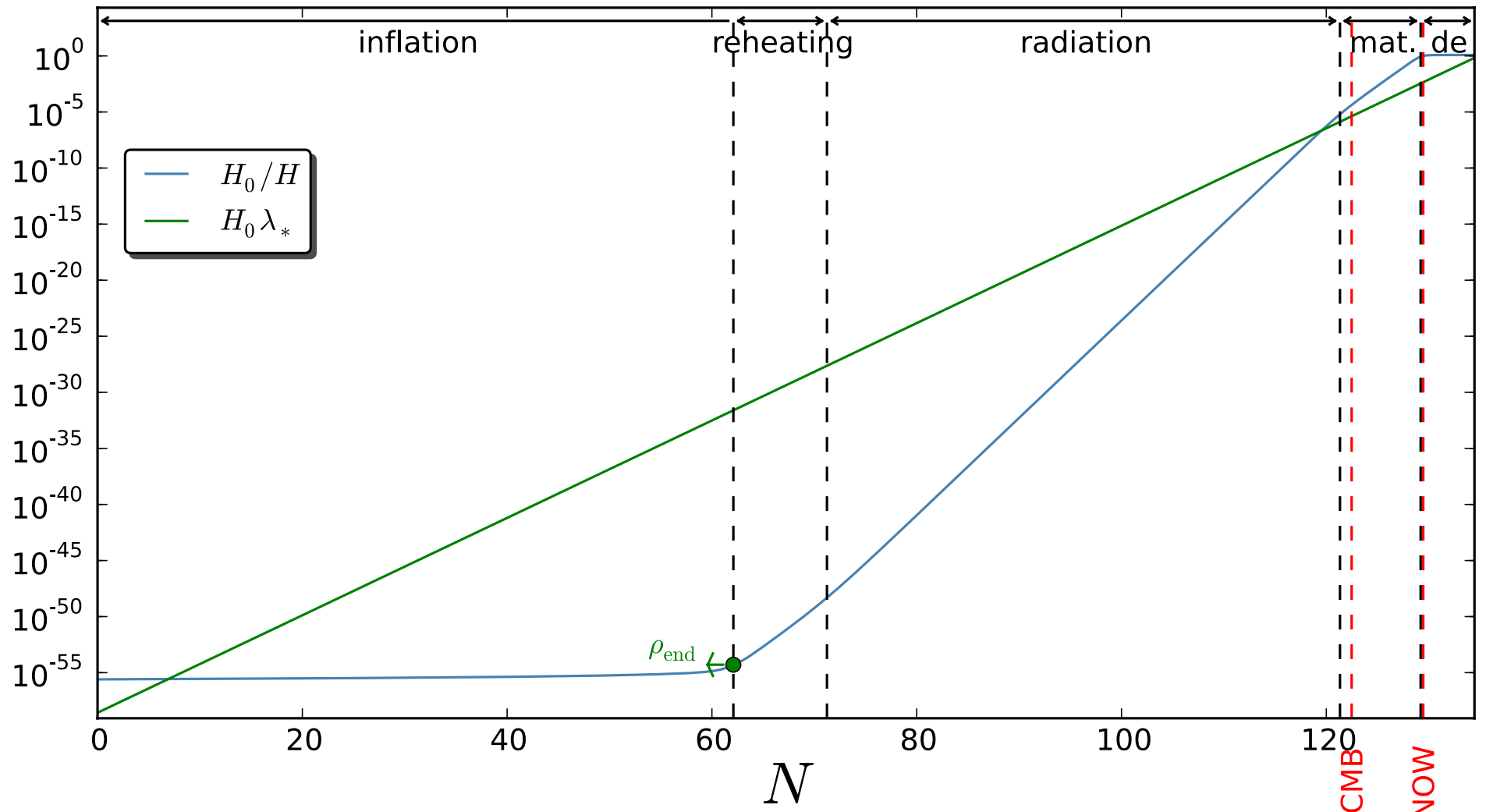
A technical aspect



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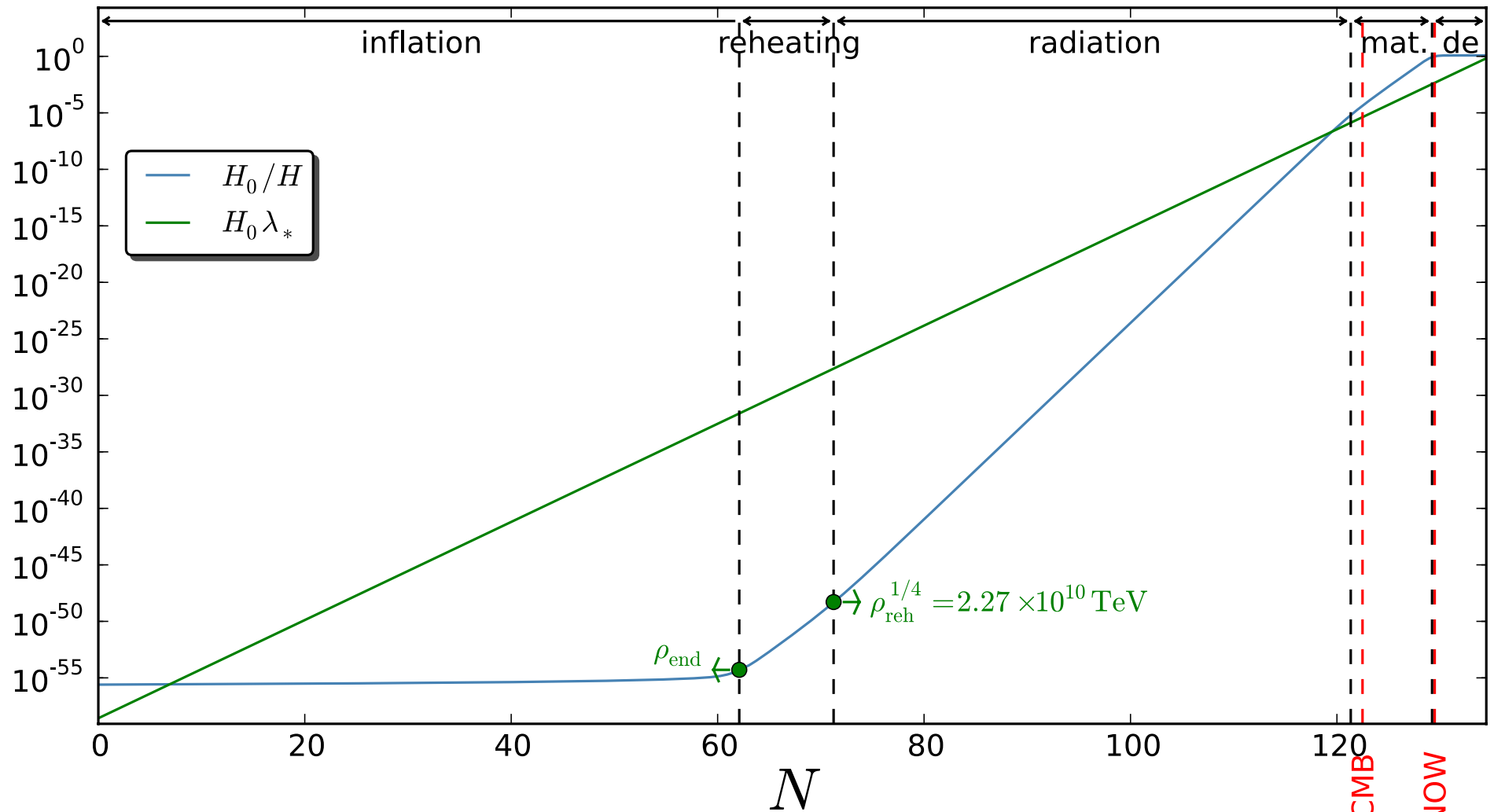




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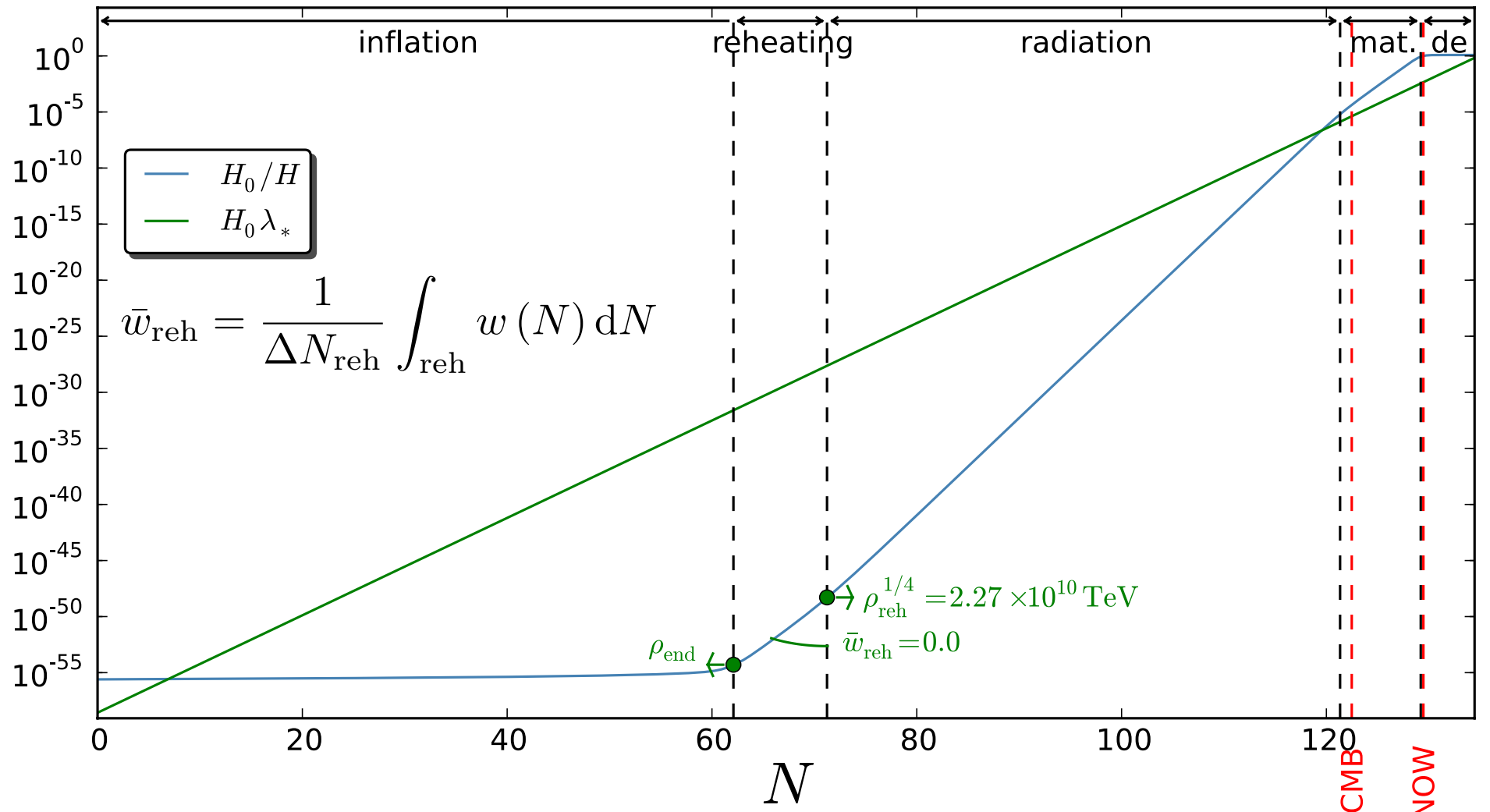
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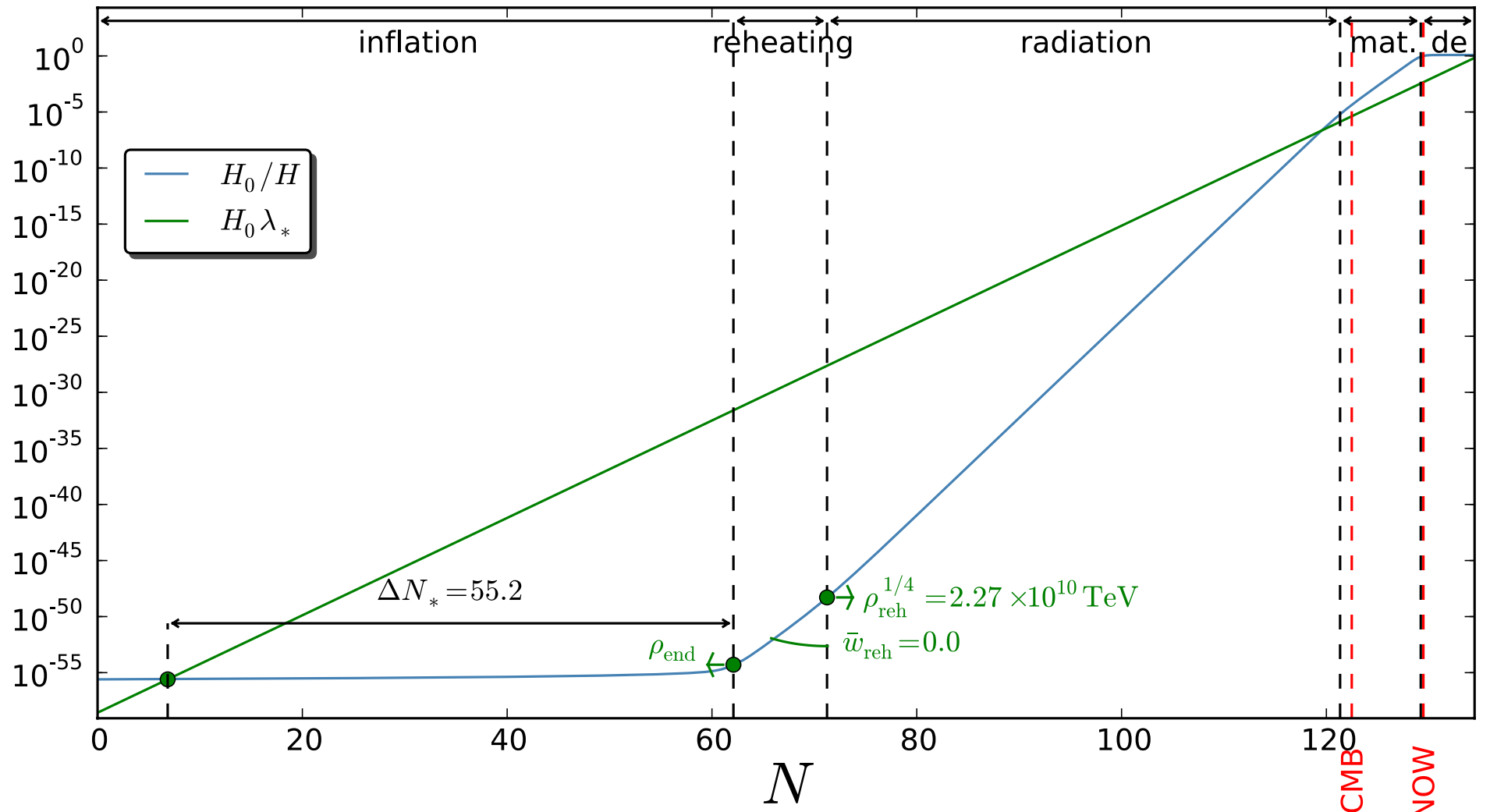
$$\bar{w}_{\text{reh}} = \frac{1}{\Delta N_{\text{reh}}} \int_{\text{reh}} w(N) dN$$

CMB  
NOW

# Role of Reheating

Martin & Ringeval (2010)  
Easter & Peiris (2011)

A technical aspect

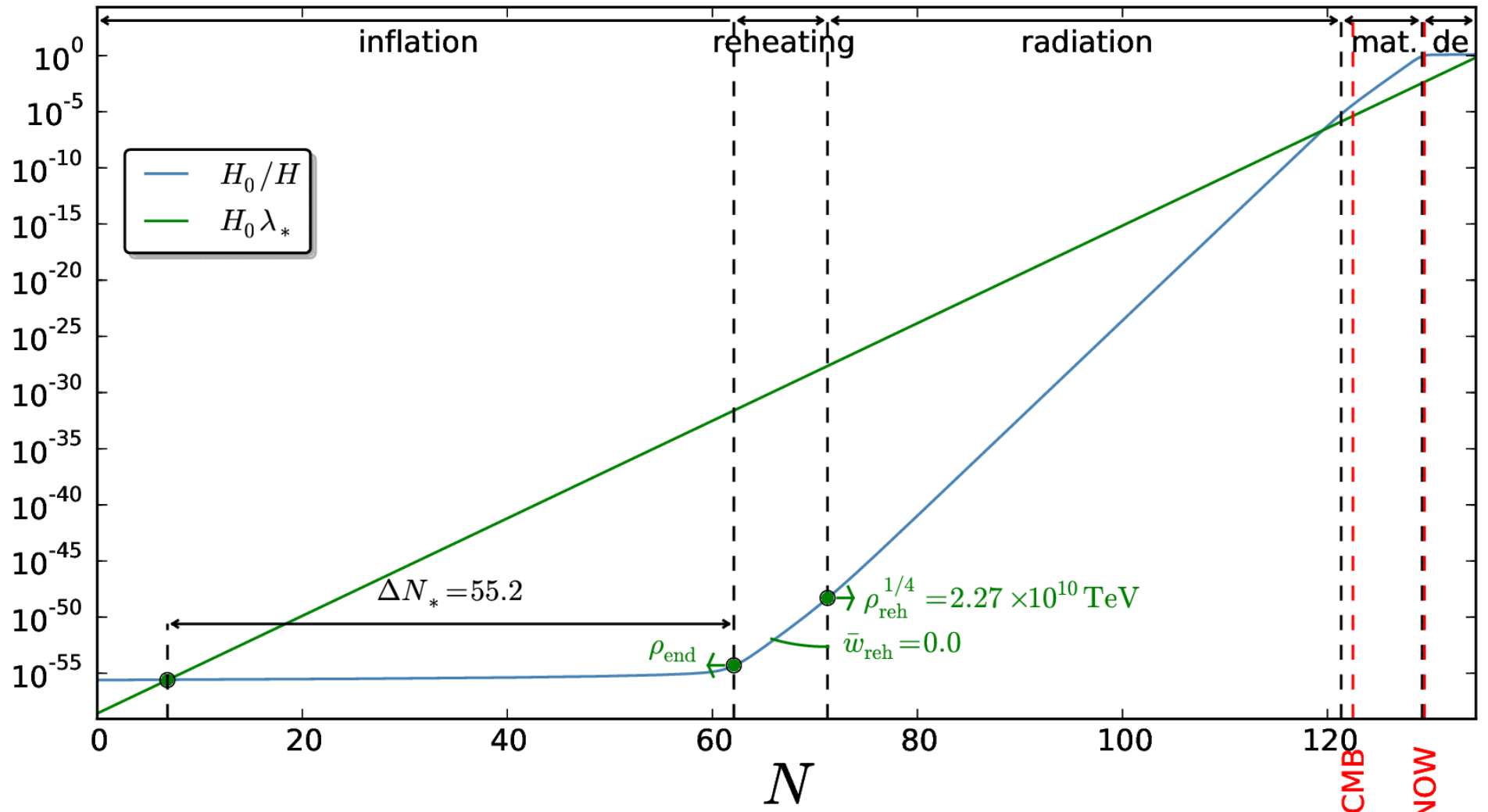


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Easter & Peiris (2011)

A technical aspect

$$\rho_{\text{BBN}} < \rho_{\text{reh}} < \rho_{\text{end}}$$

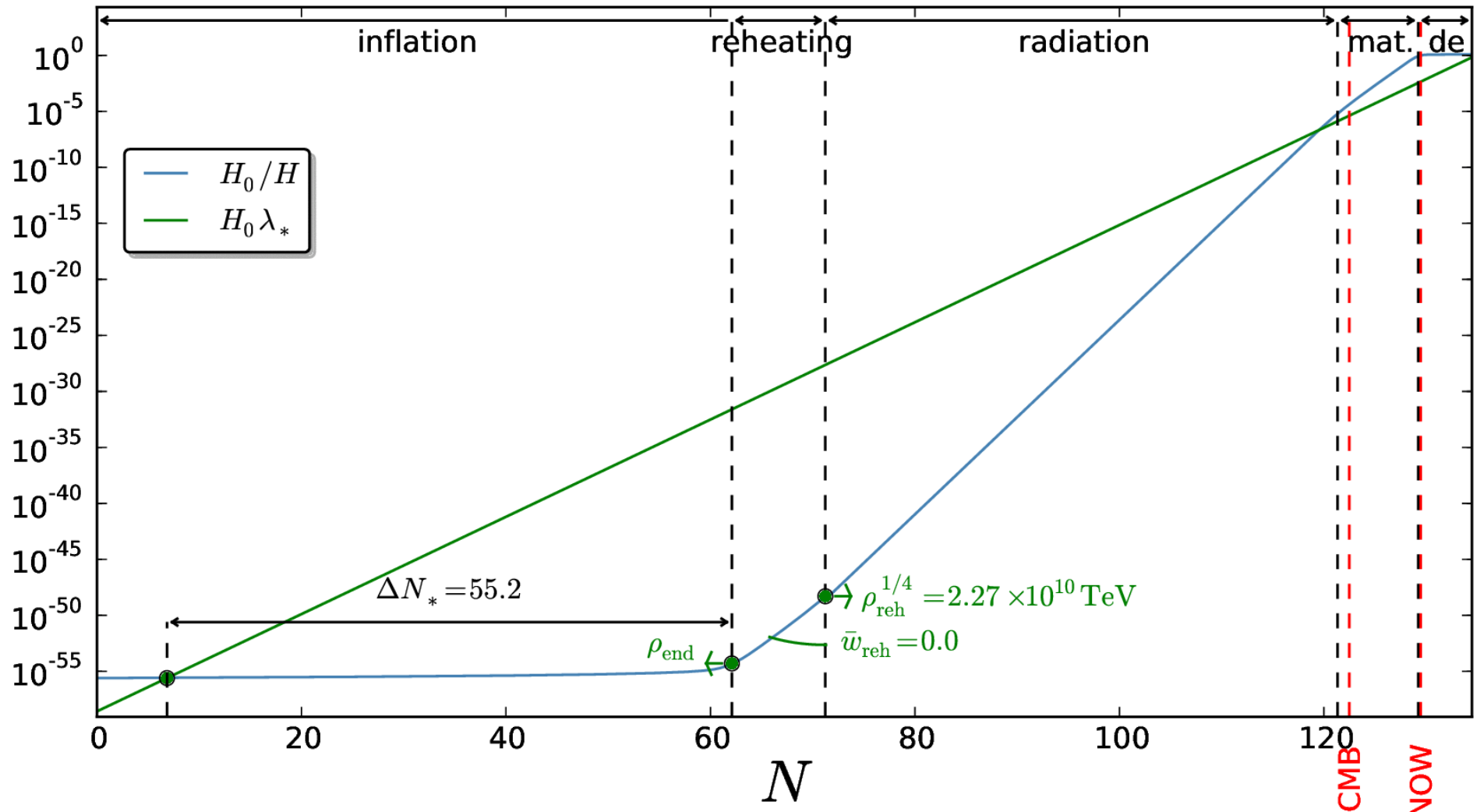


# Role of Reheating

Martin & Ringeval (2010)  
Easter & Peiris (2011)

A technical aspect

$$-1/3 < \bar{w}_{\text{reh}} < 1$$



# Role of Reheating

Martin & Ringeval (2010)  
Easter & Peiris (2011)

$$\begin{aligned}\Delta N_* = & \frac{1 - 3\bar{w}_{\text{reh}}}{12(1 + \bar{w}_{\text{reh}})} \ln \left( \frac{\rho_{\text{reh}}}{\rho_{\text{end}}} \right) \\ & + \frac{1}{4} \ln \left( \frac{\rho_*}{3M_{\text{Pl}}^4} \frac{\rho_*}{\rho_{\text{end}}} \right) \\ & - \ln \left( \frac{k_*/a_{\text{now}}}{\rho_{\gamma, \text{now}}^{1/4}} \right)\end{aligned}$$

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- depends on reheating parameters
- depends on V parameters (model dependent)



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- depends on reheating parameters
- depends on V parameters
- accurately measured
- implicit equation  
(requires numerical solving)

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In practice:

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In practice: •  $\rho_{\text{BBN}} < \rho_{\text{reh}} < \rho_{\text{end}}$

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not inflation

dominant energy condition

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$$\bullet \quad \left( -\frac{1}{3} \right) < \bar{w}_{\text{reh}} < 1$$

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- $\rho_{\gamma}$  set to measured value,  $k_*/a_{\text{now}}$  chosen to pivot scale

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Yields a model dependent range of admitted values for  $\Delta N_*$

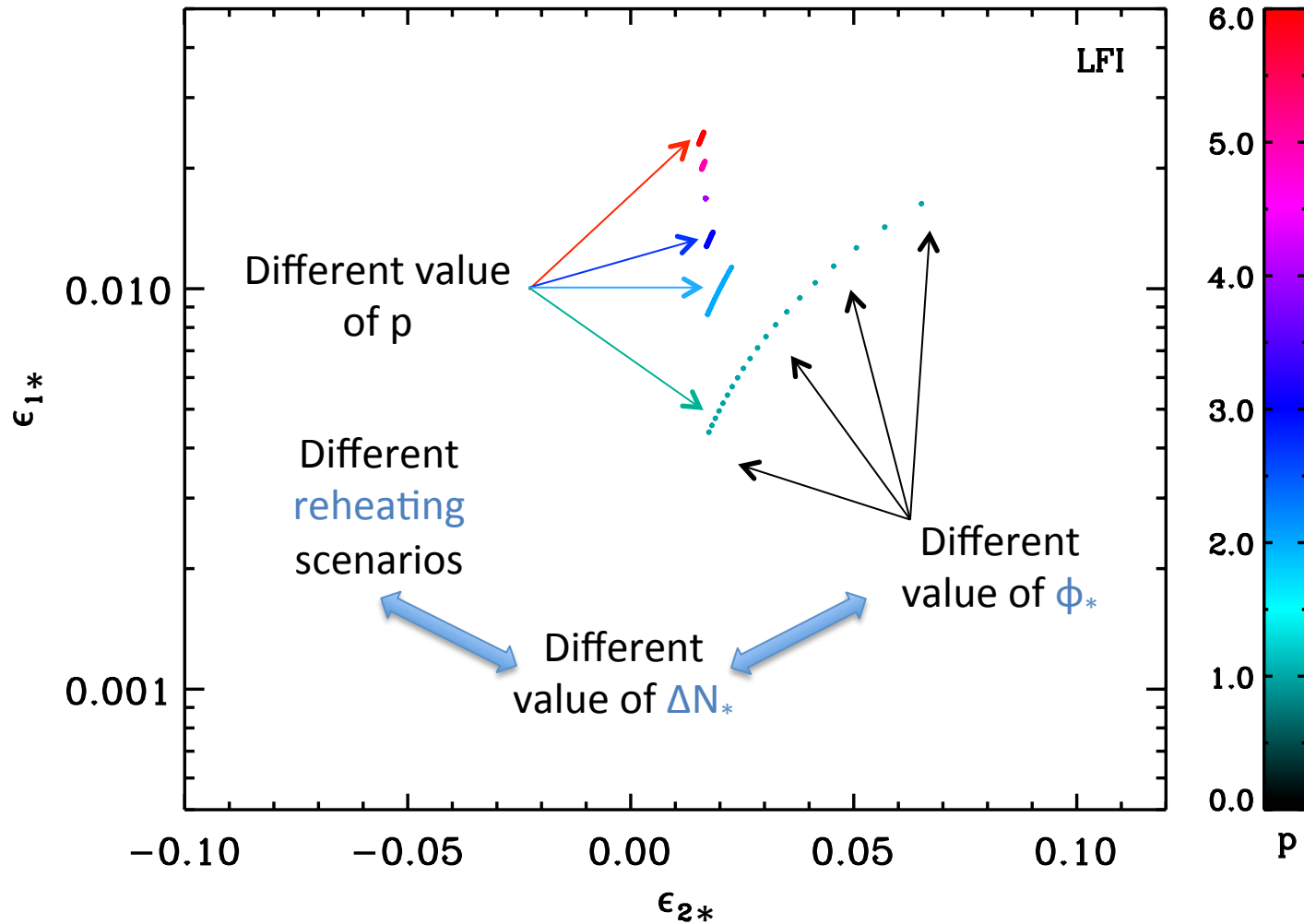
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# Model Predictions

Described in terms of  $\epsilon_{i*}$

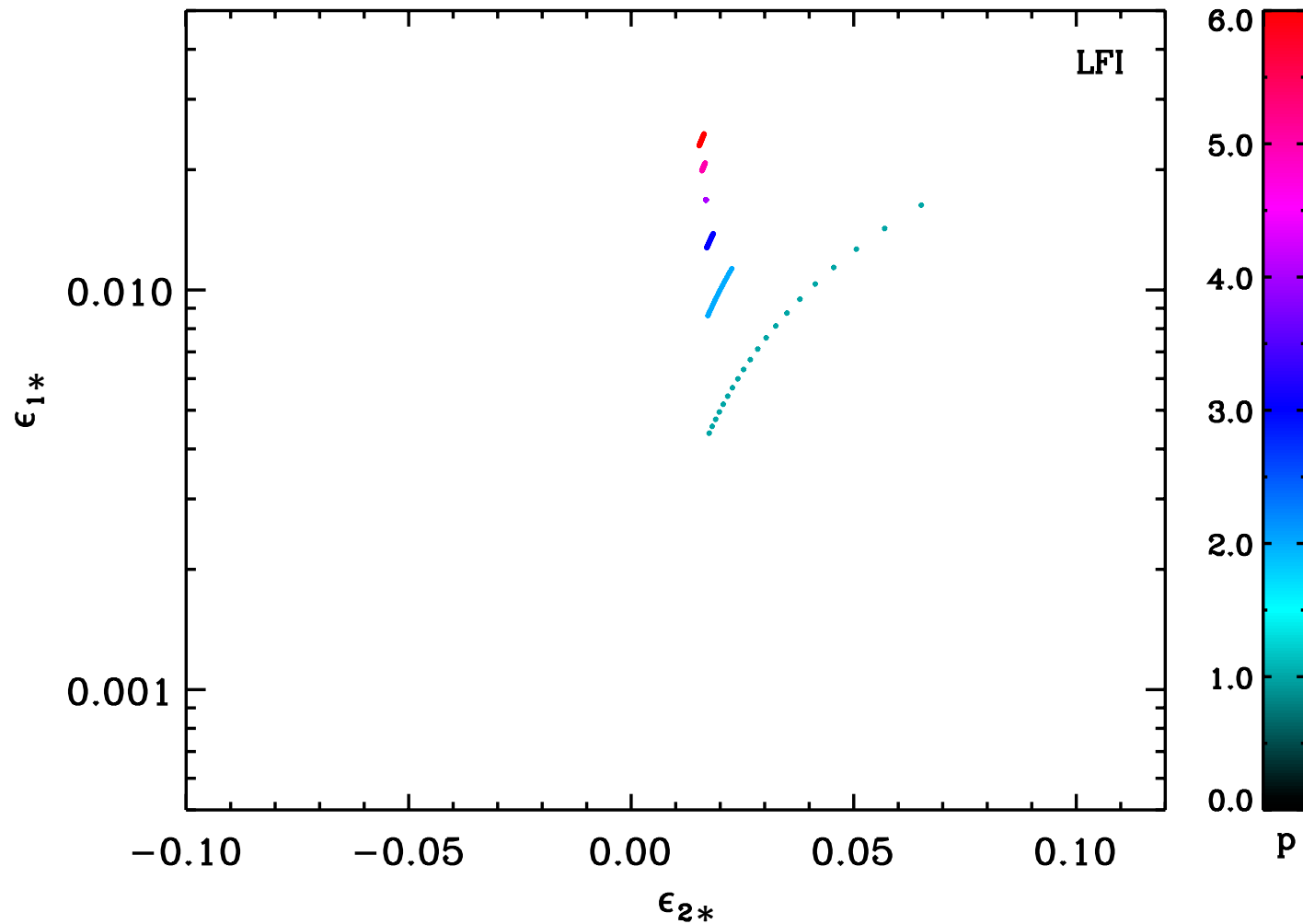
An example: « large field inflation »  $V(\phi) = M^4 \left( \frac{\phi}{M_{\text{Pl}}} \right)^p$



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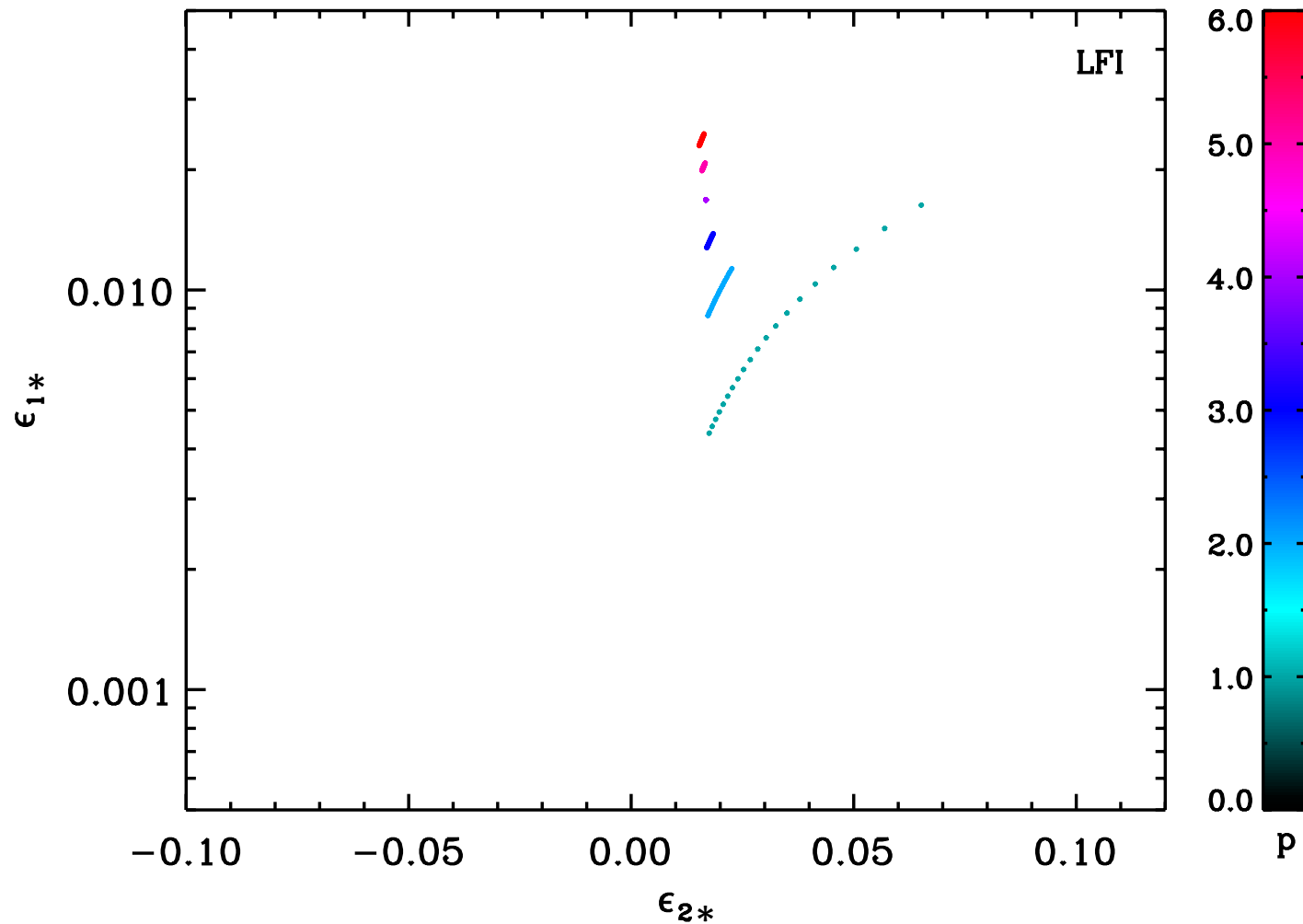


# Model Predictions and Observations

Martin, Ringeval, V.V (2013)

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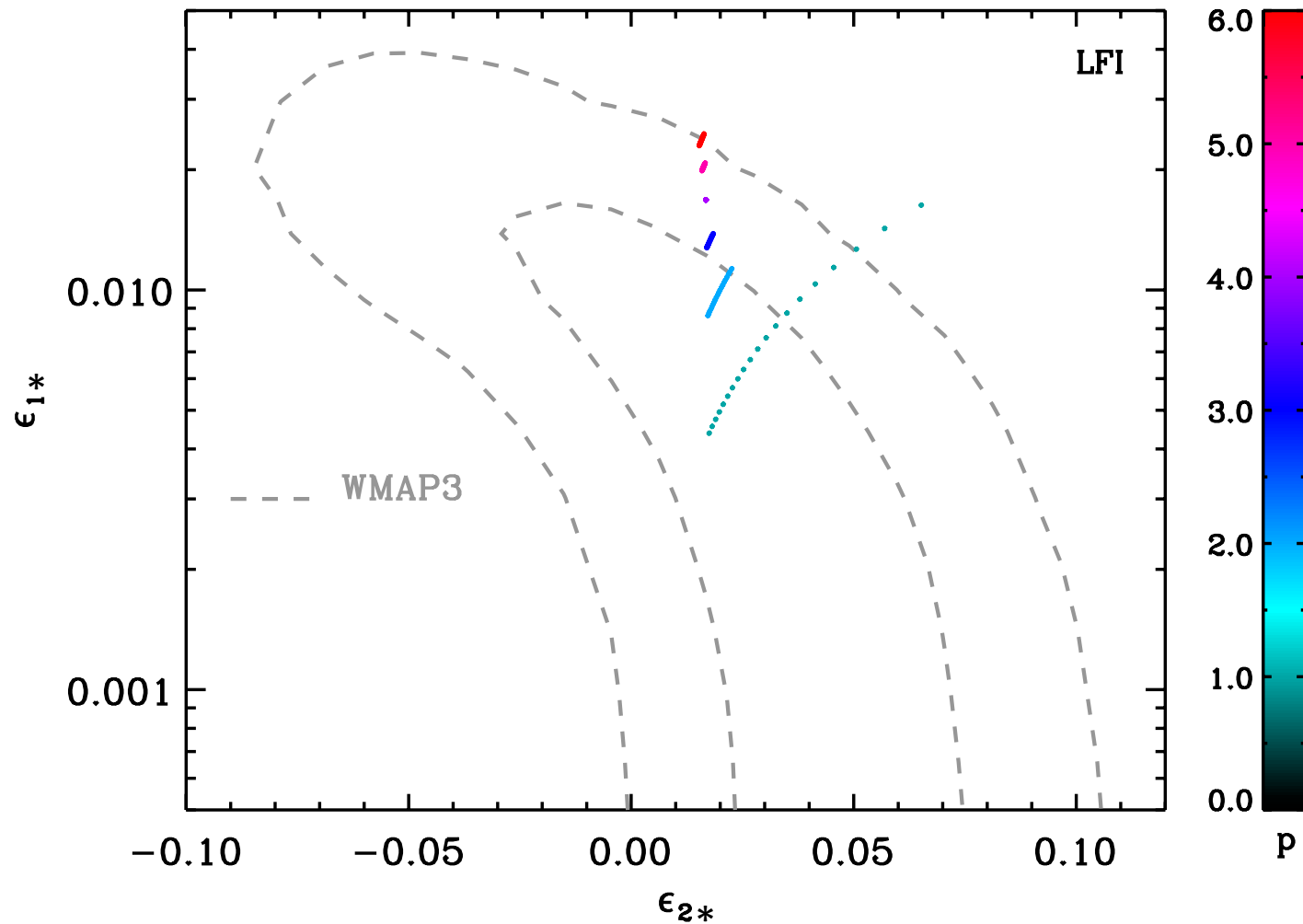


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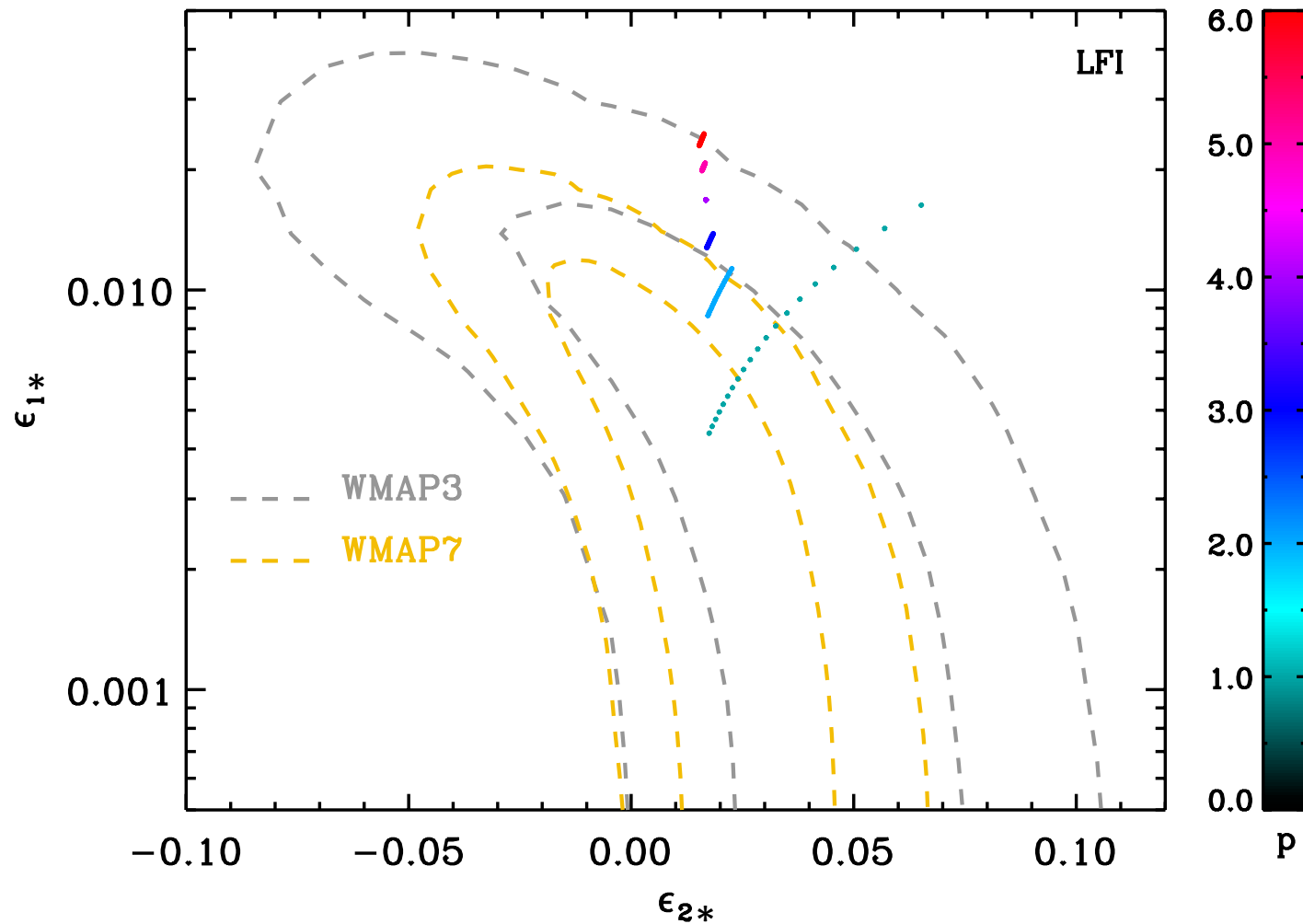


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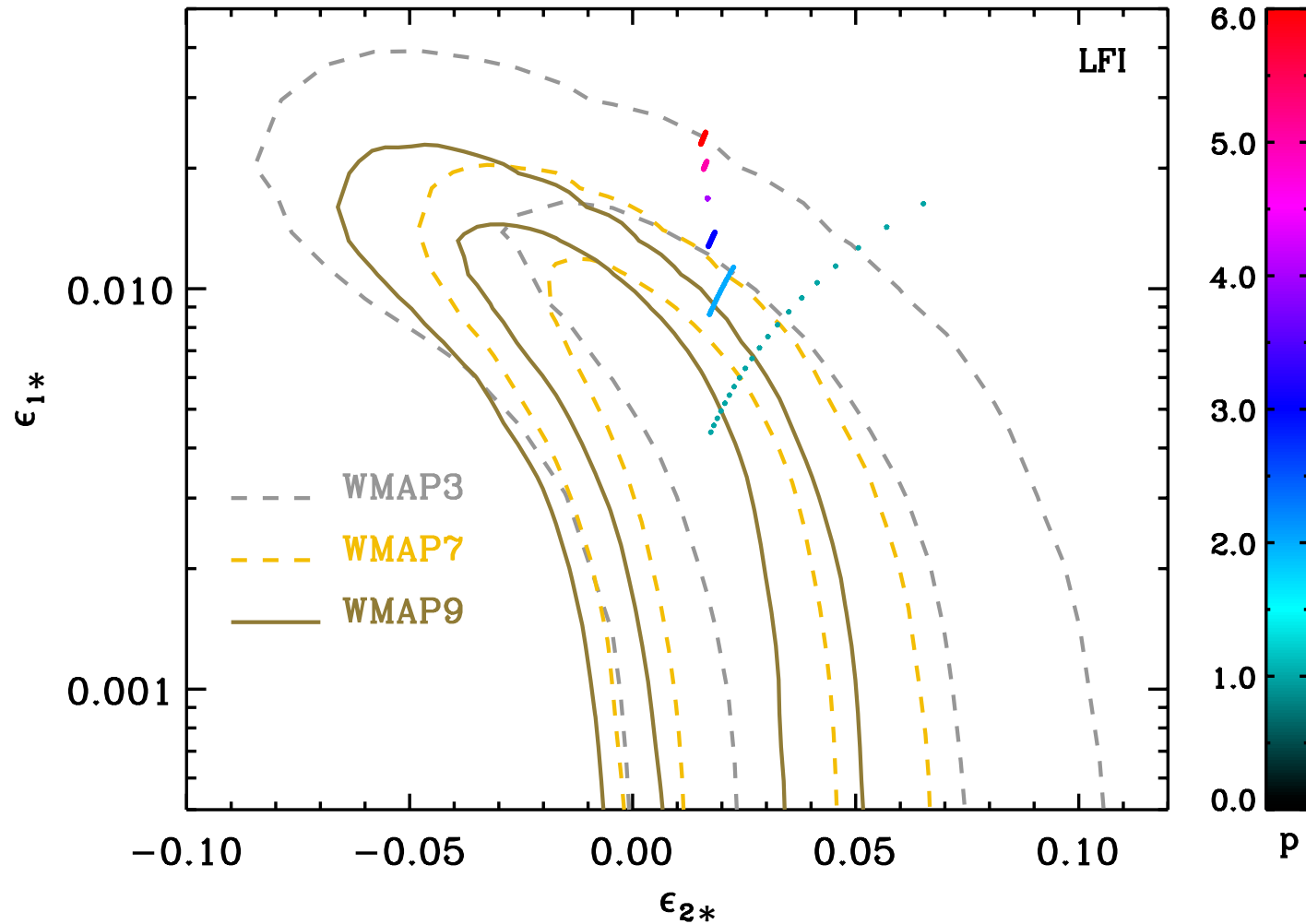


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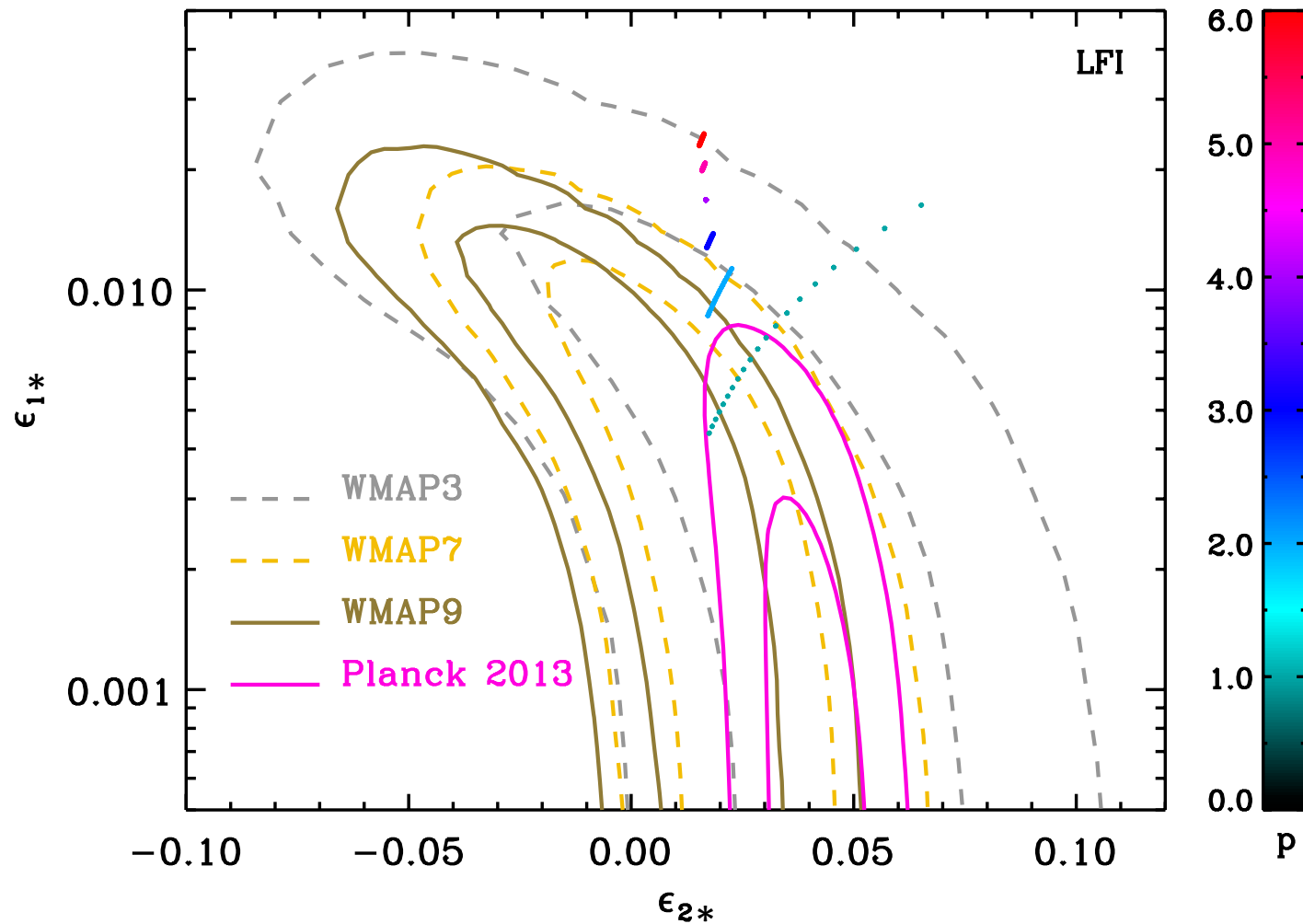


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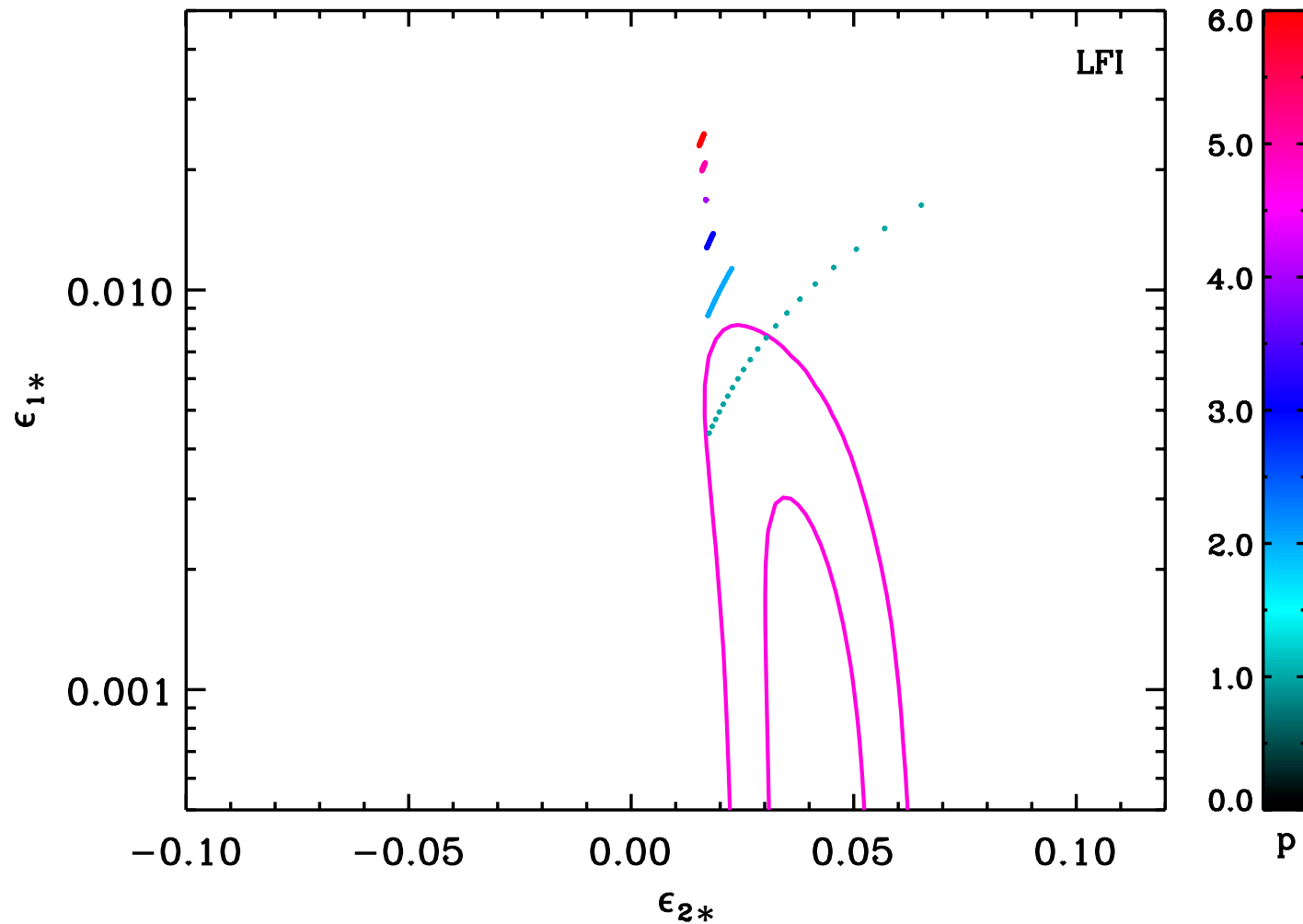


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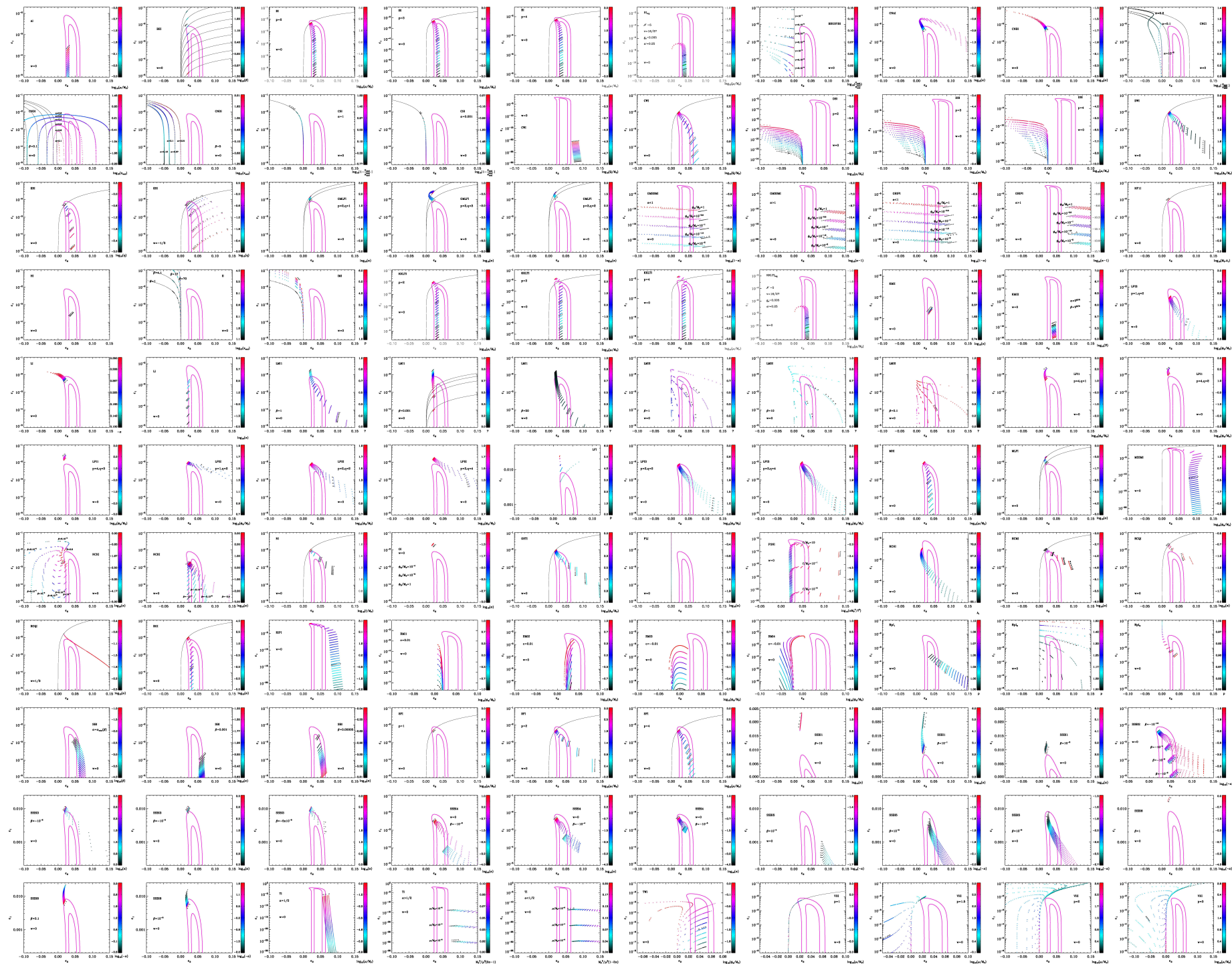
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# Bayesian Approach

to model comparison

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*Bayesian evidence: Integral of the likelihood over parameter prior*

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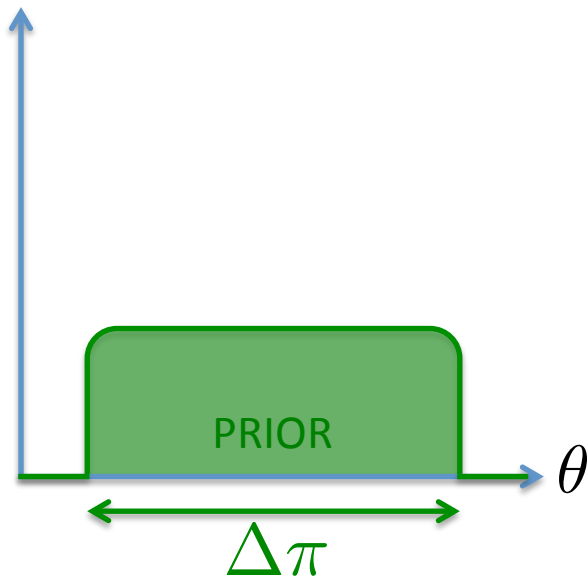
*Bayesian evidence: Integral of the likelihood over parameter prior*



# Bayesian Approach

to model comparison

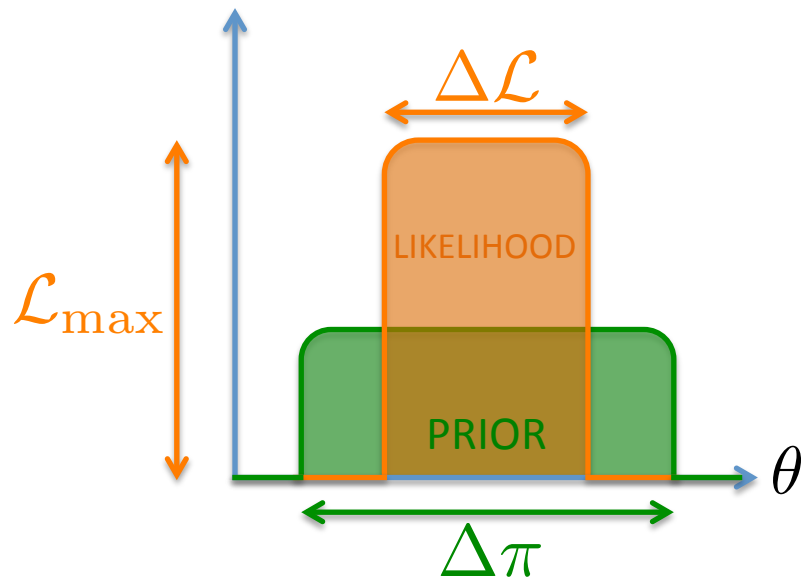
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# Bayesian Approach

to model comparison

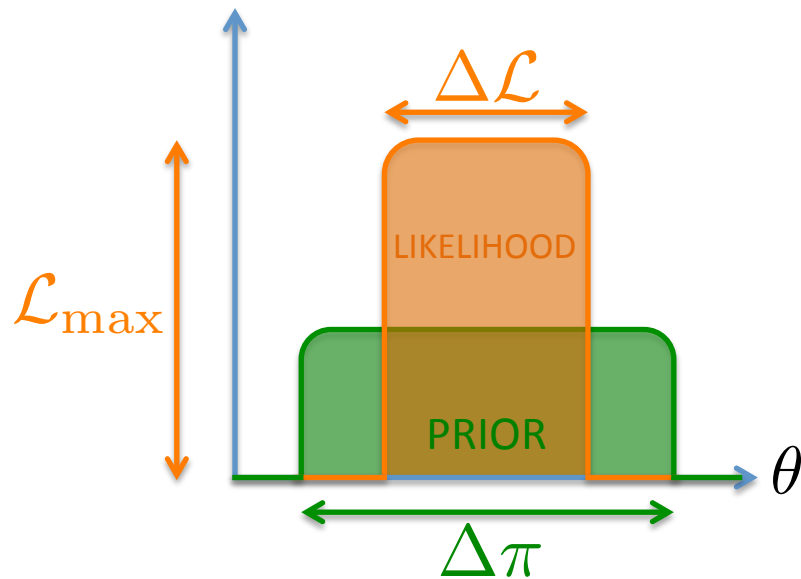
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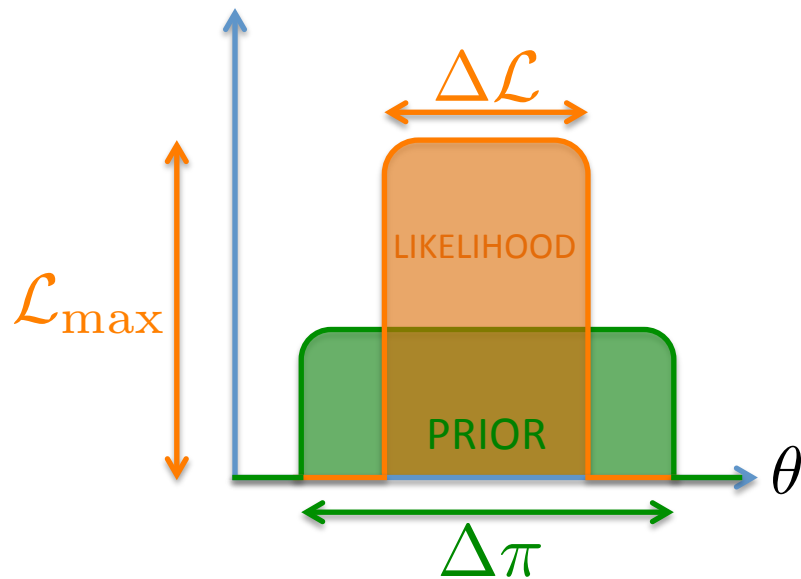


$$\mathcal{E}(\mathcal{M}) = \mathcal{L}_{\max} \frac{\Delta \mathcal{L}}{\Delta \pi}$$

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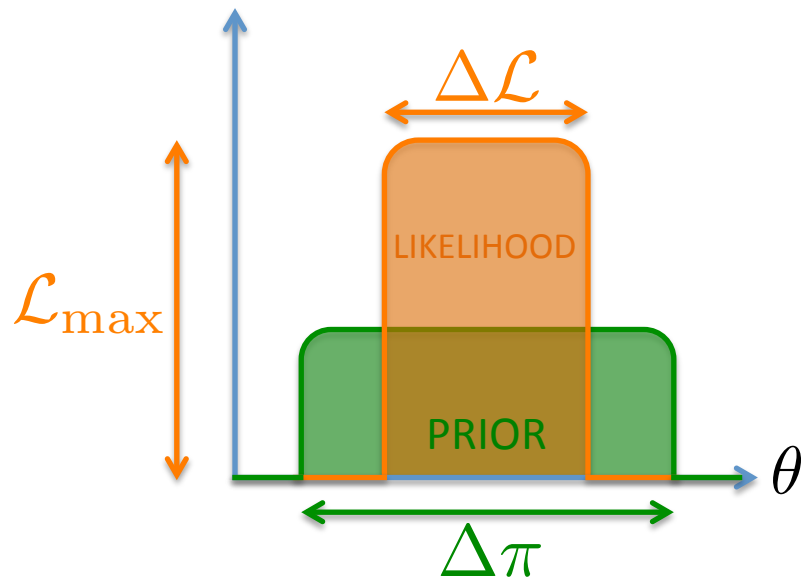
Compromise between **quality of fit** and **simplicity**



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Compromise between **quality of fit** and **simplicity**

For a set of models  $\{\mathcal{M}_j\}$ ,

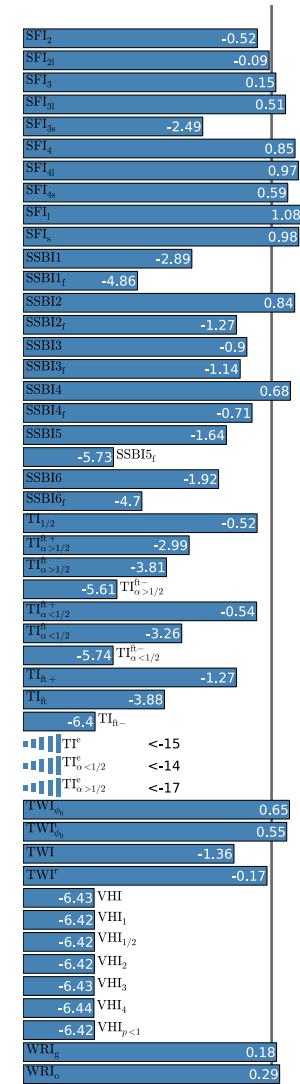
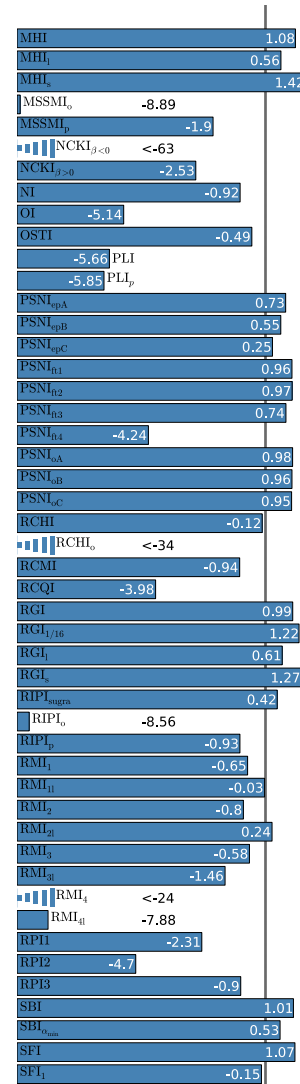
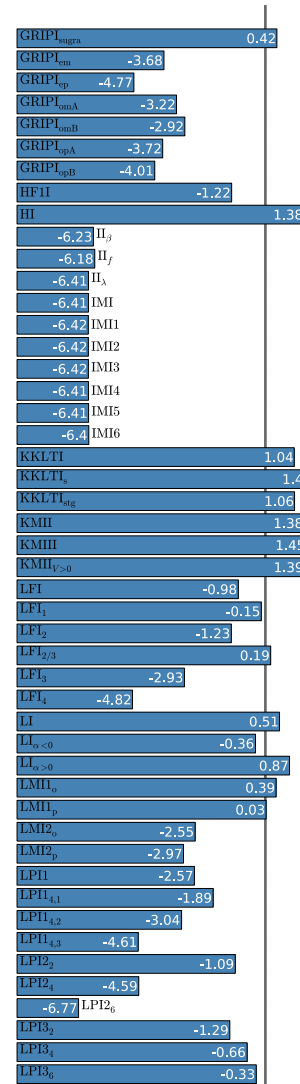
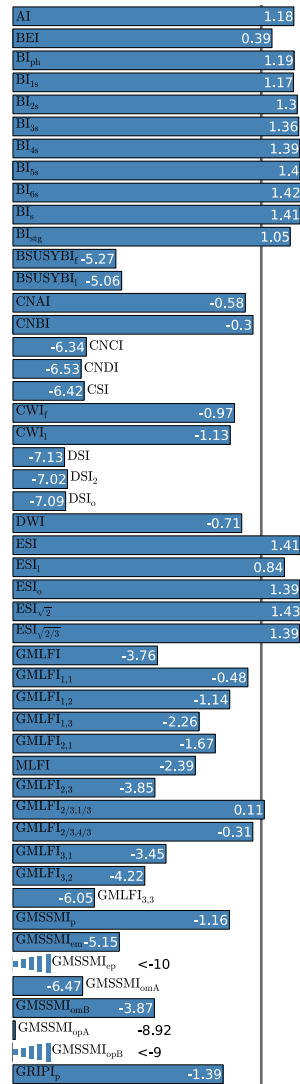
the posterior-to-prior ratio  $N_{\text{models}} \mathcal{E}(\mathcal{M}_i) / \sum \mathcal{E}(\mathcal{M}_j)$

represents the degree by which the data have modified our a priori relative belief

in a given model  $\mathcal{M}_i$ .

# Posterior-to-Prior Ratio computed with Planck

Martin, Ringeval, V.V (2014)



# Bayesian evidences computed with Planck

*Summary of the results*

Martin, Ringeval, V.V (2014)

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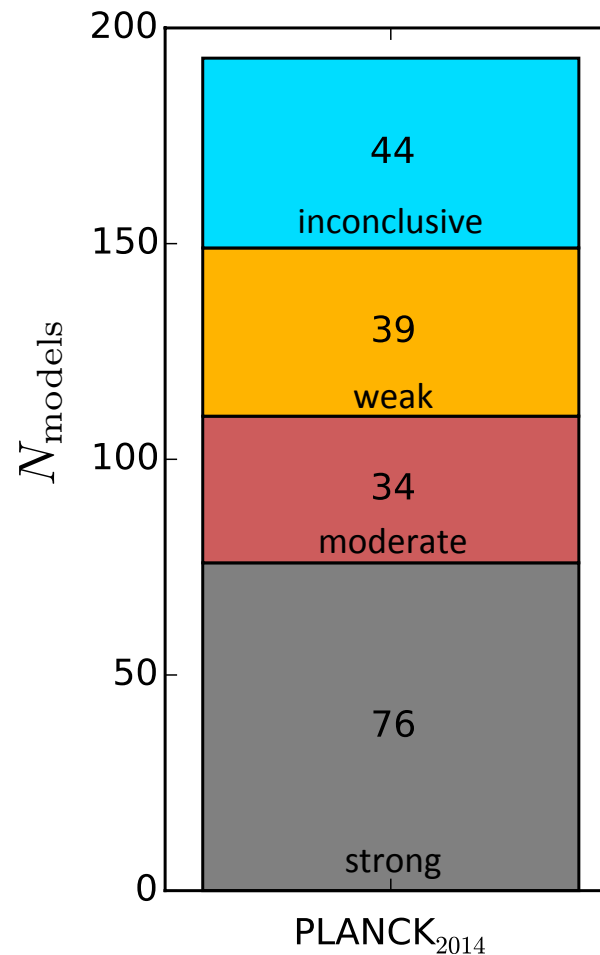
- One third of the models are “ruled out”

# Bayesian evidences computed with Planck

Martin, Ringeval, V.V (2014)

*Summary of the results*

- One third of the models are “ruled out” according to the Jeffreys scale



# Bayesian evidences computed with Planck

Martin, Ringeval, V.V (2014)

*Summary of the results*

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- Planck favors “Plateau Inflation”

# Bayesian evidences computed with Planck

Martin, Ringeval, V.V (2014)

*Summary of the results*

- One third of the models are “ruled out”
- Planck favors “Plateau Inflation”
- Some models are killed by “fine-tuning”

# Constraints on Reheating

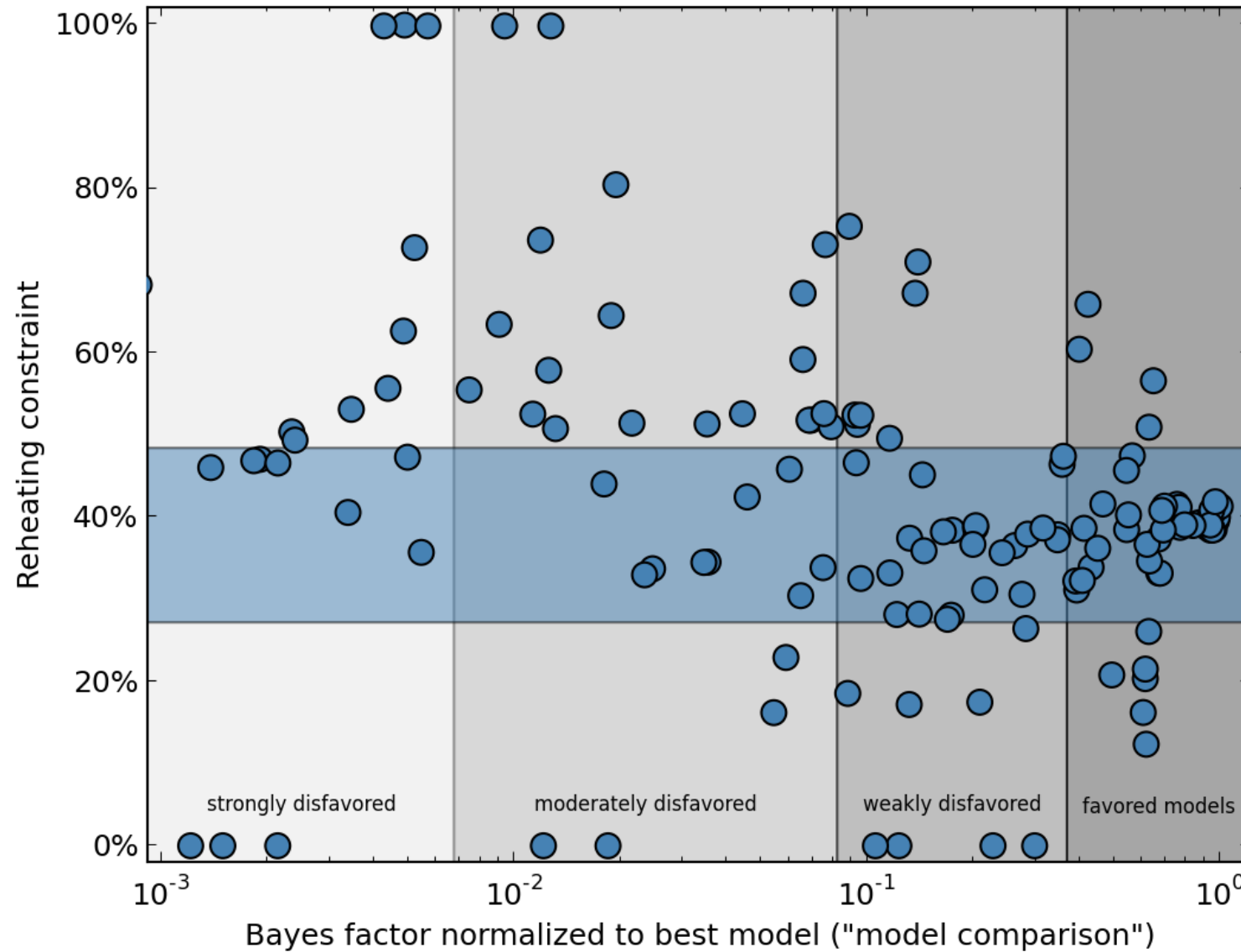
Martin, Ringeval, V.V (2014)

$$\ln R_{\text{reh}} = \frac{1 - 3\bar{w}_{\text{reh}}}{12(1 + \bar{w}_{\text{reh}})} \ln \left( \frac{\rho_{\text{reh}}}{\rho_{\text{end}}} \right) + \ln \left( \frac{\rho_{\text{end}}^{1/4}}{M_{\text{Pl}}} \right)$$



# Constraints on Reheating

Martin, Ringeval, V.V (2014)



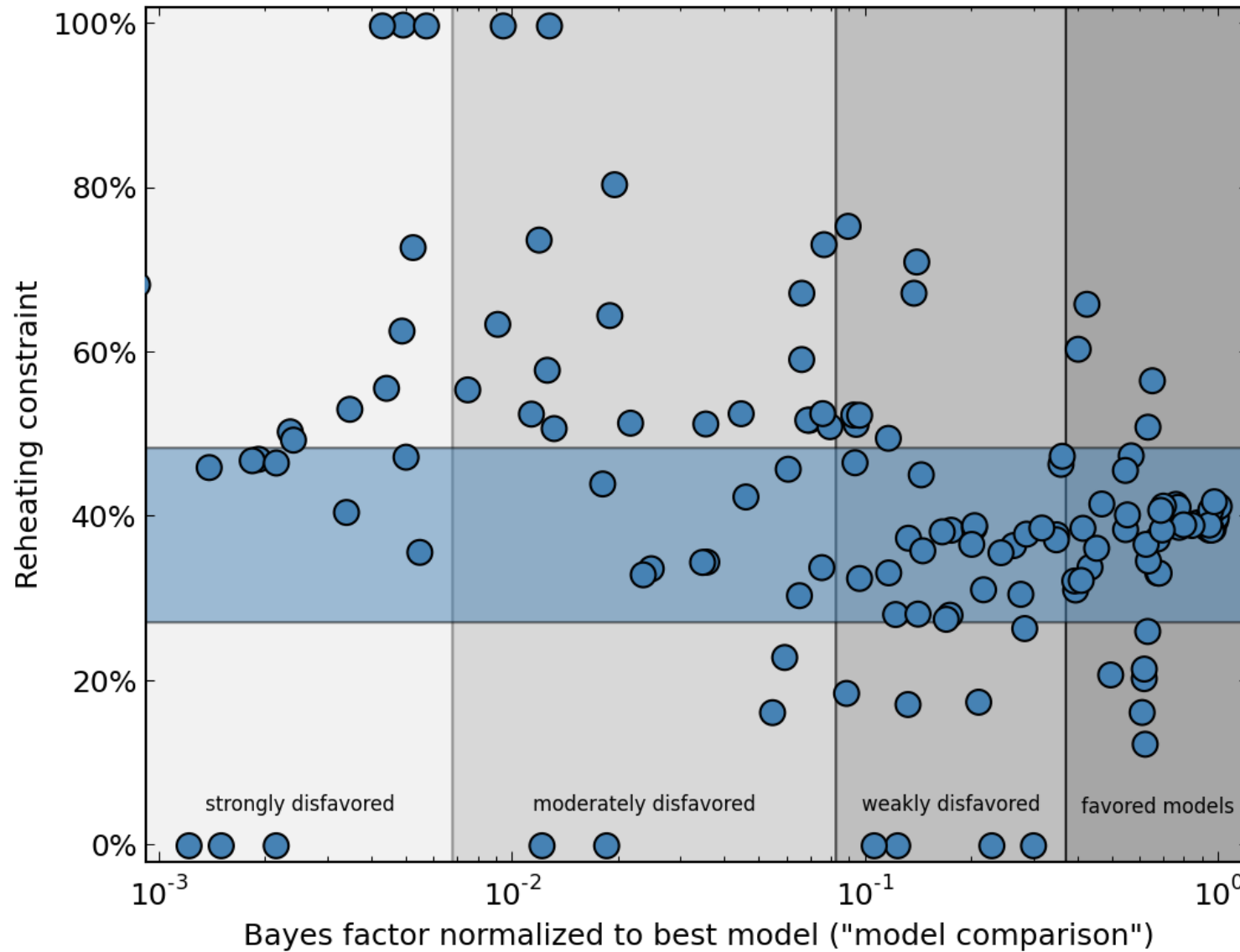
# Constraints on Reheating

Martin, Ringeval, V.V (2014)

**WORST**



**BEST**



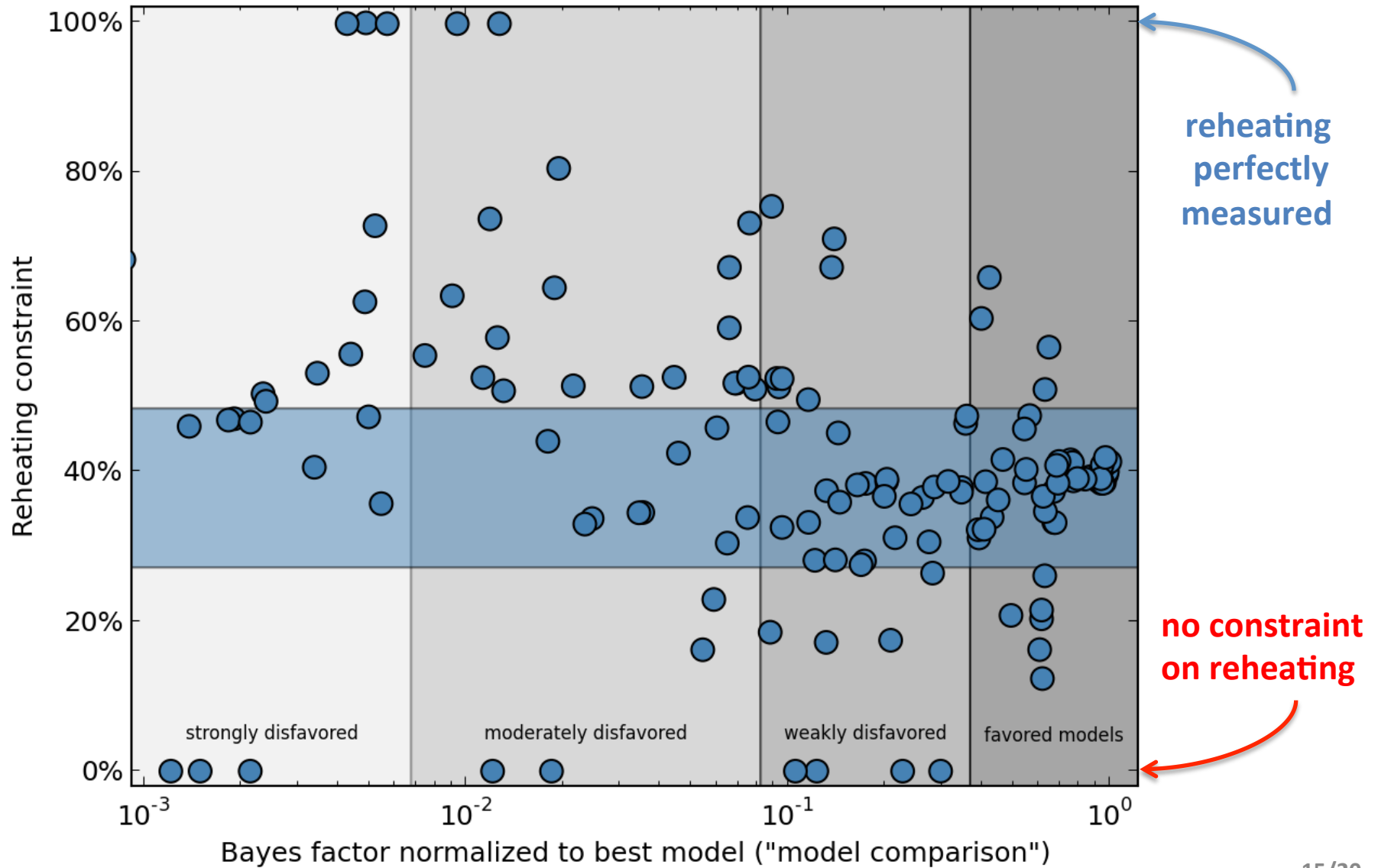
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Martin, Ringeval, V.V (2014)

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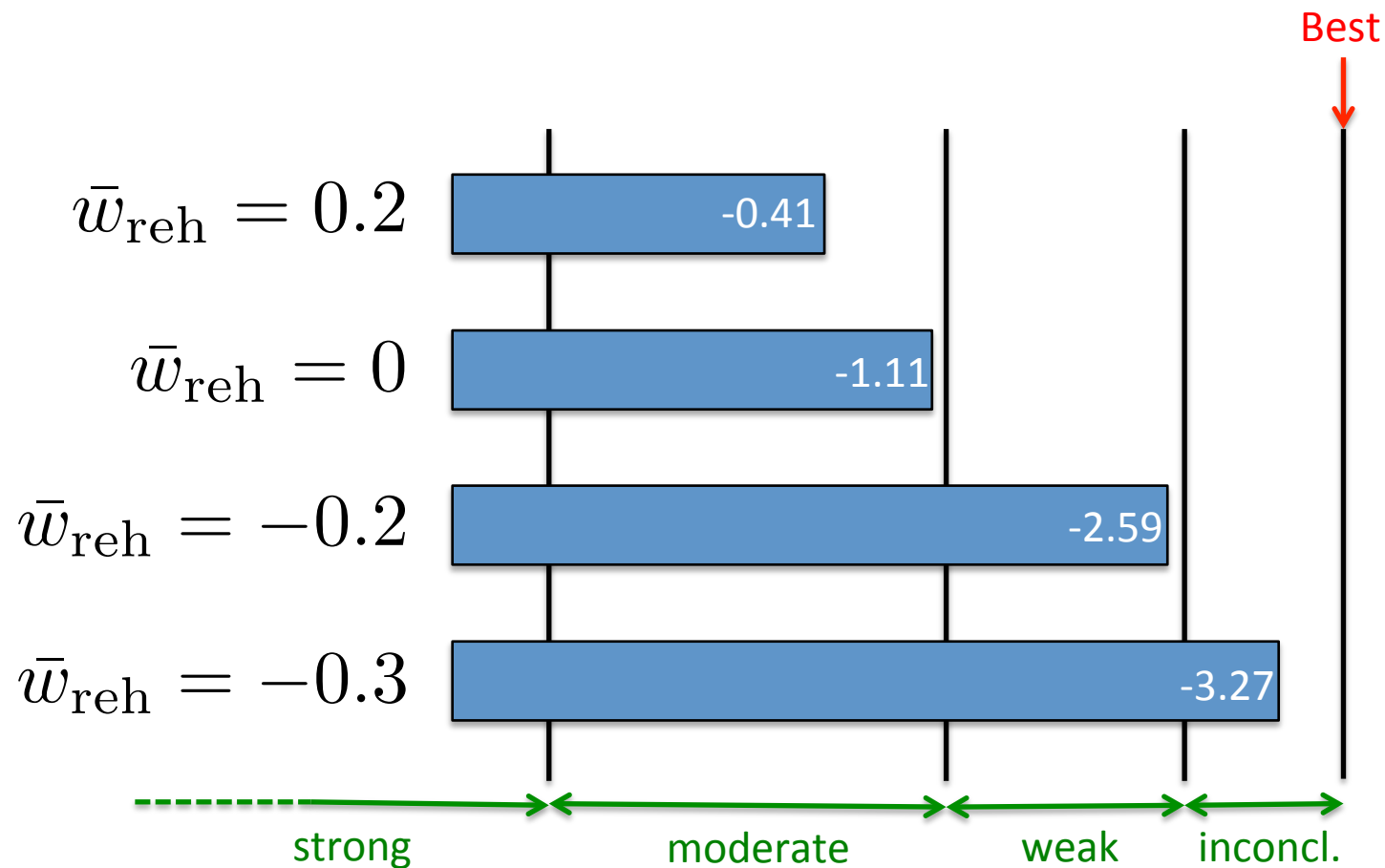
**BEST**



# Reheating does Matters!

Martin, Ringeval, V.V (2014)

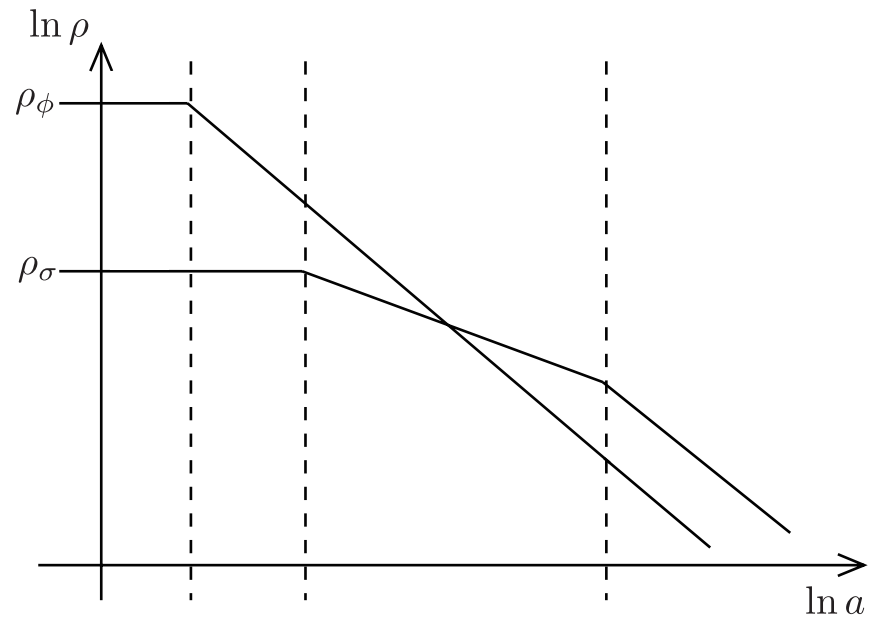
Example :  $\text{LI}_{\alpha>0}$        $V(\phi) = M^4 \left[ 1 + \alpha \ln \left( \frac{\phi}{M_{\text{Pl}}} \right) \right]$



# Adding a Light Scalar Field

## Curvaton Scenarios

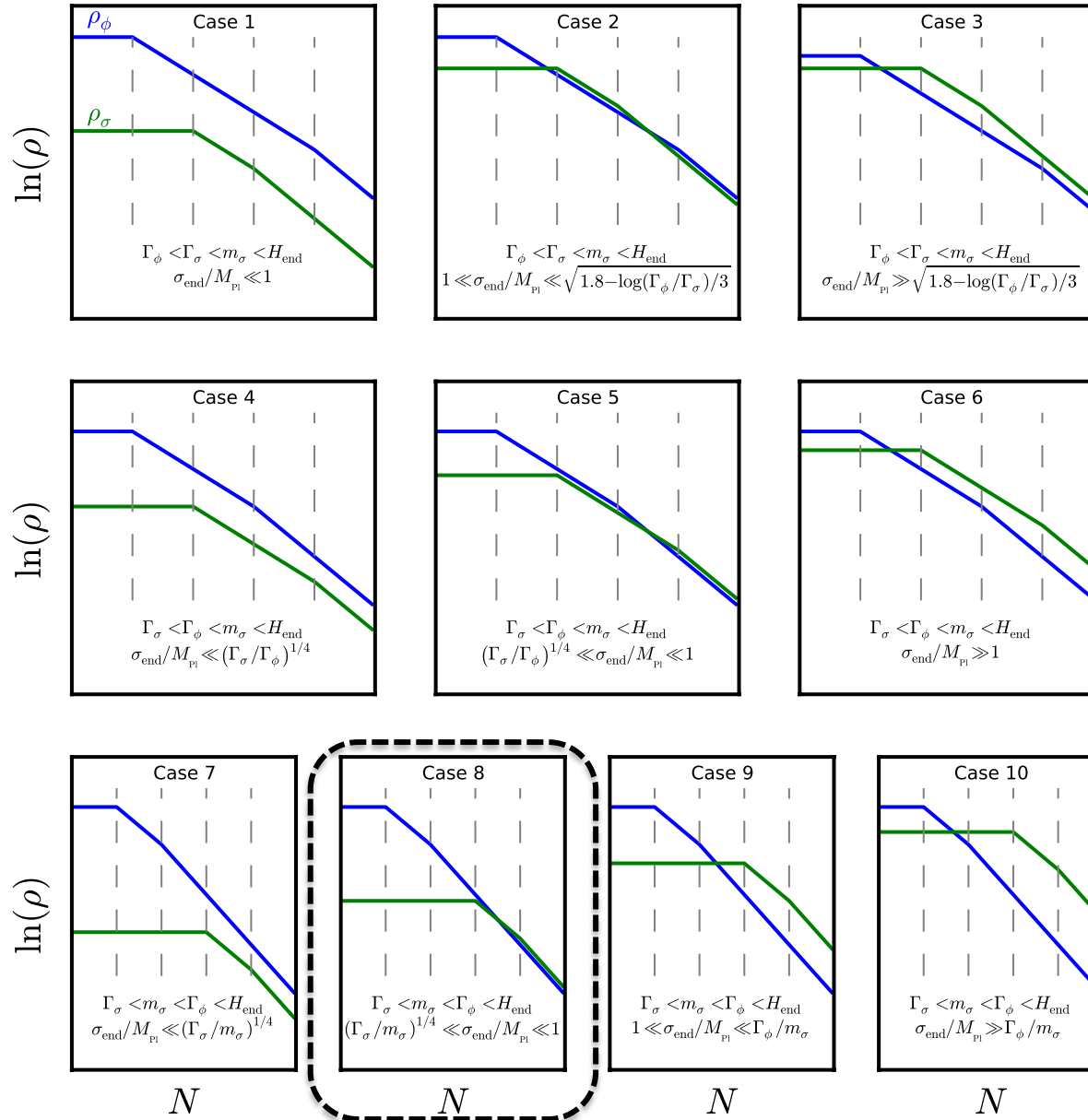
- Linde and Mukhanov, 1997
- Enqvist and Sloth, 2001
- Lyth and Wands, 2001
- Moroi and Takahashi, 2001
- Langlois and Vernizzi, 2004



$$\zeta = \zeta_\phi + \zeta_\sigma \quad \Downarrow \quad \bar{w}_{\text{reh}}(m_\sigma, \sigma_{\text{end}}, \Gamma_\phi, \Gamma_\sigma)$$

# Adding a Light Scalar Field

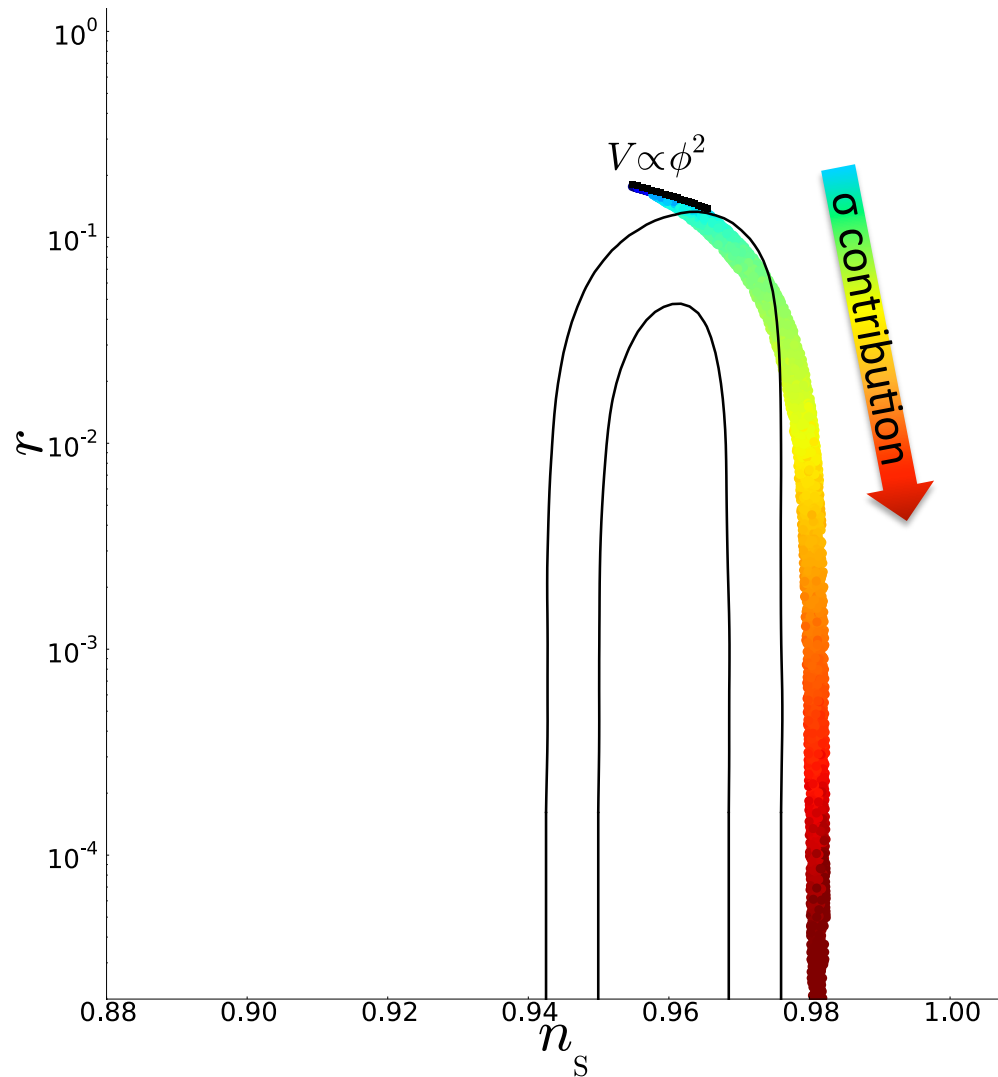
V.V, Koyama and Wands (2015)



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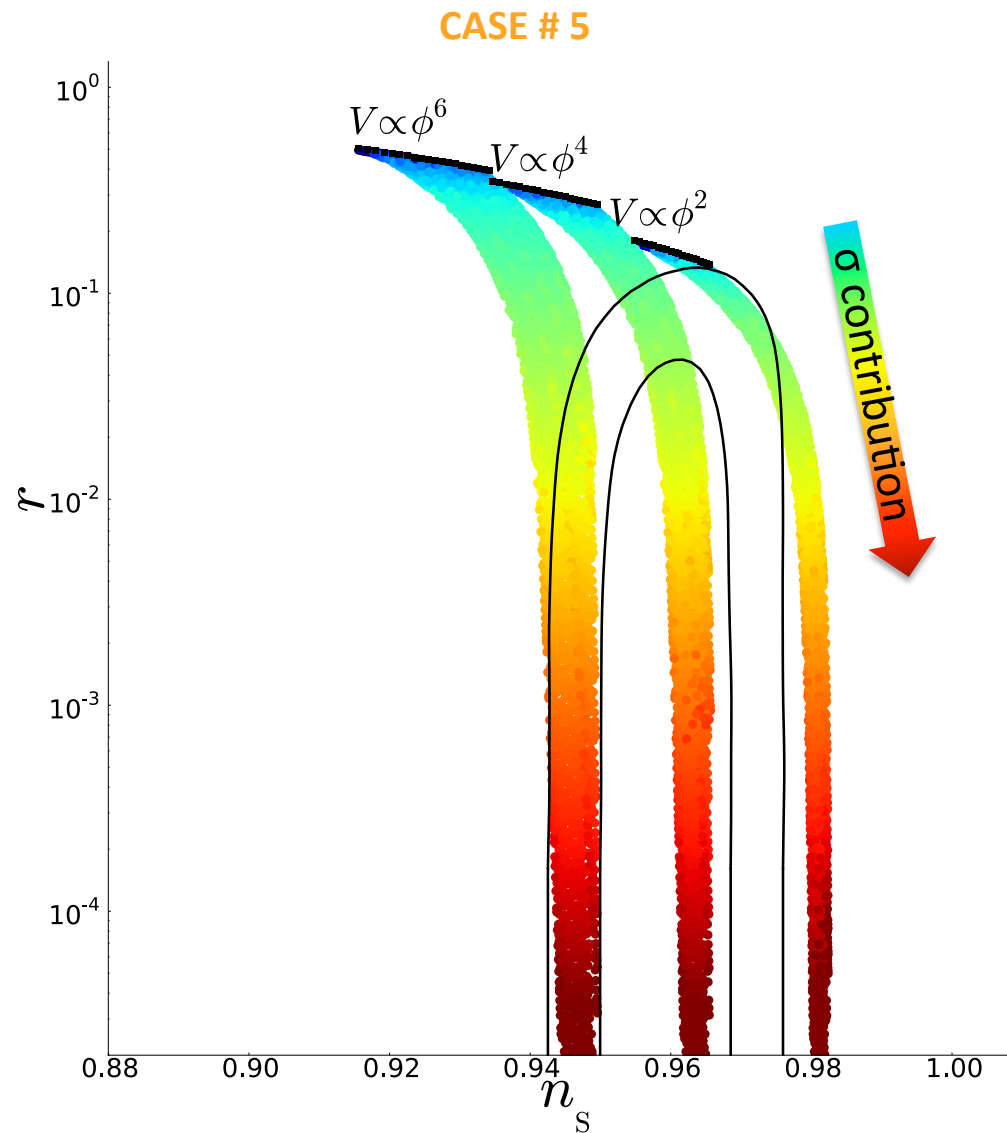
V.V, Koyama and Wands (2015)

CASE # 5



# Adding a Light Scalar Field

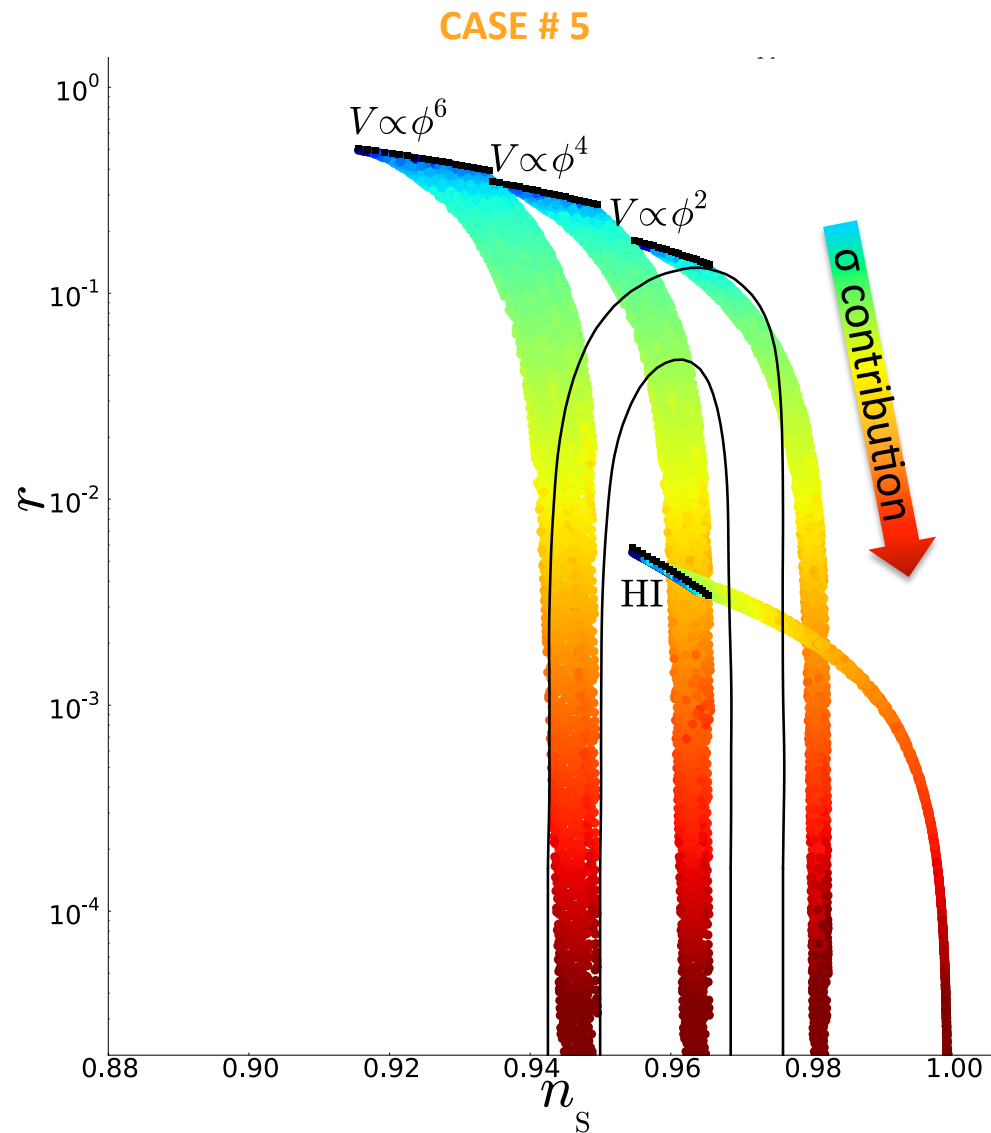
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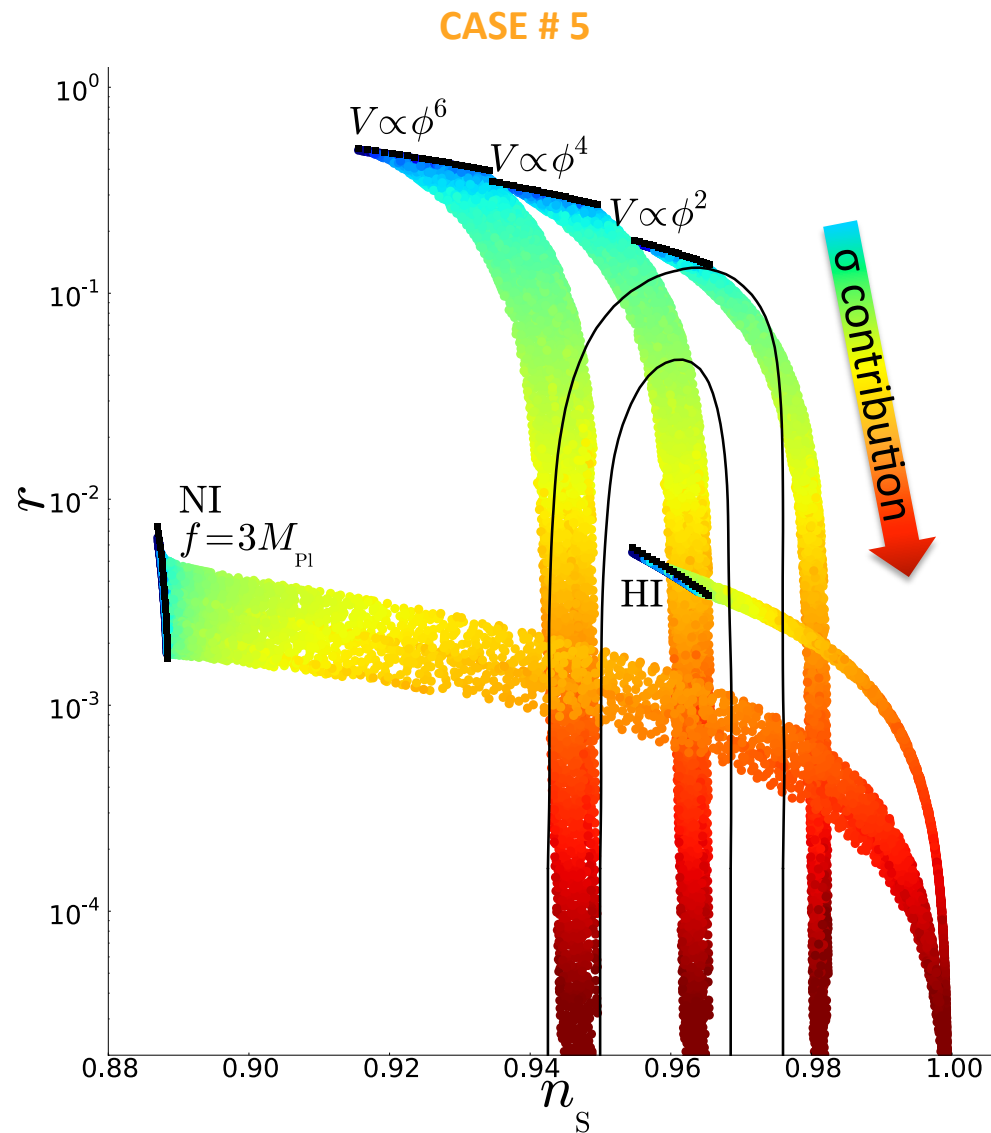
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V.V, Koyama and Wands (2015)



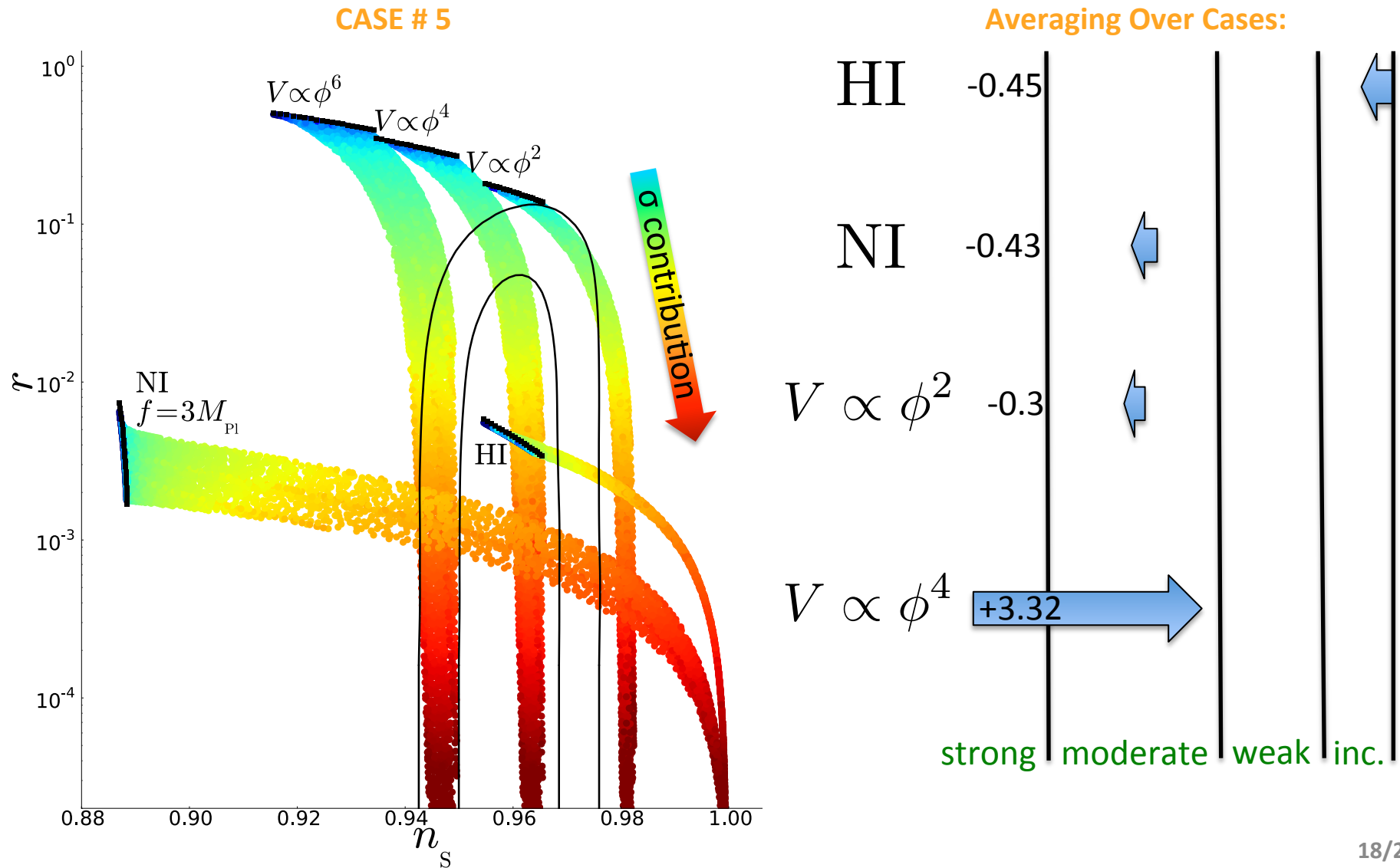
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V.V, Koyama and Wands (2015)



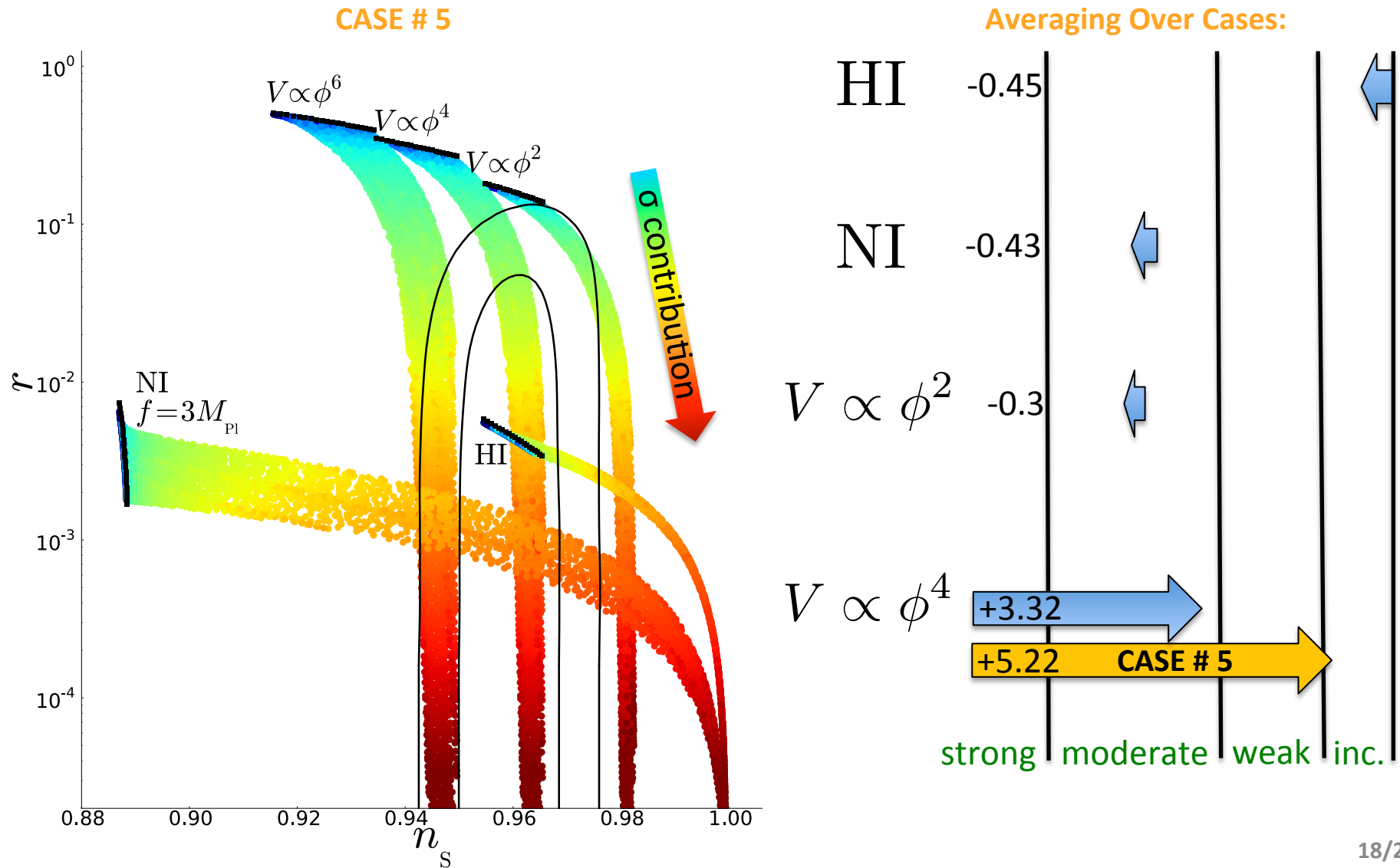
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V.V, Koyama and Wands (2015)



# Adding a Light Scalar Field

V.V, Koyama and Wands (2015)



# Conclusions and Prospects

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- **Planck** also constraints reheating (excludes 40% of prior space)

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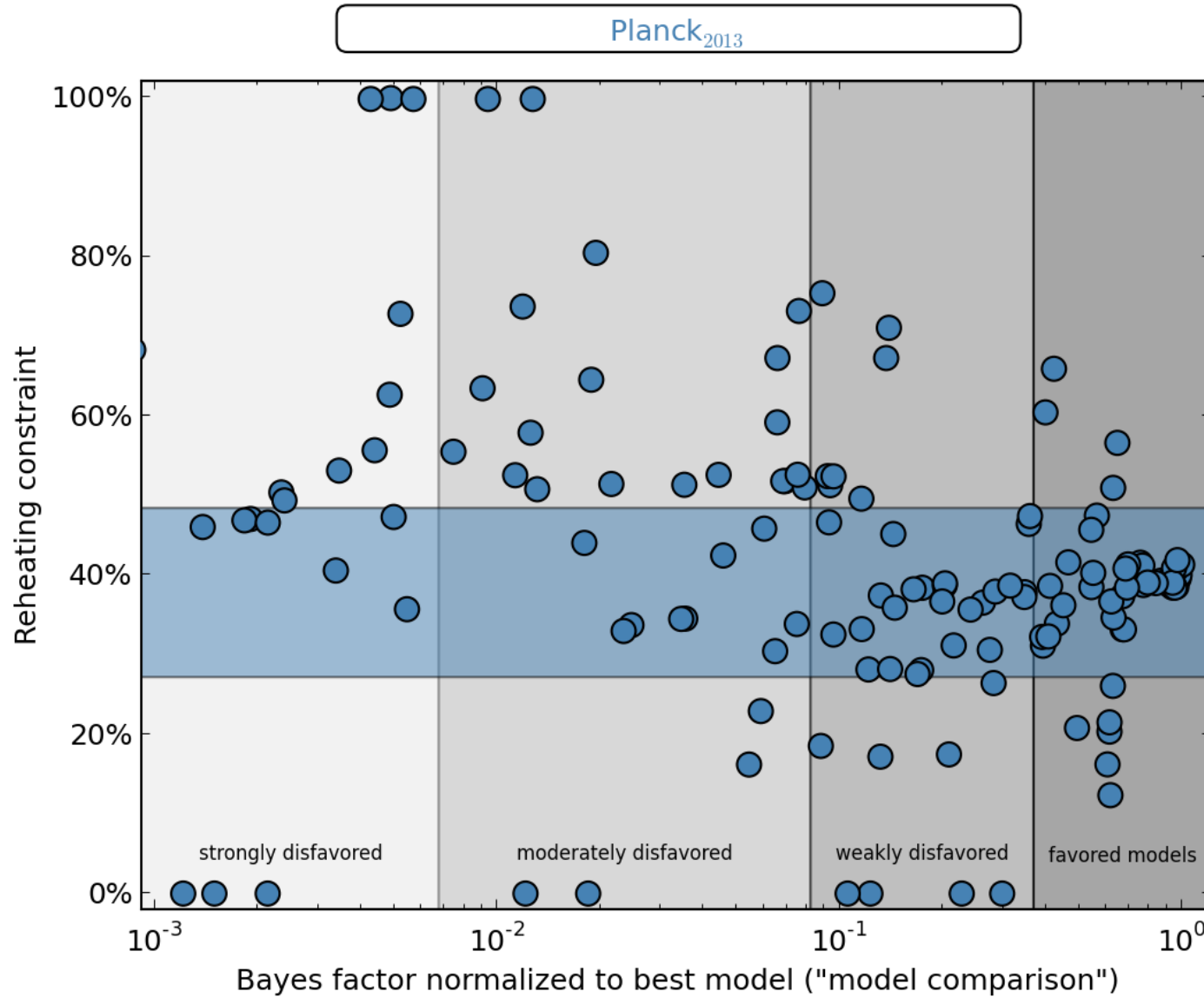
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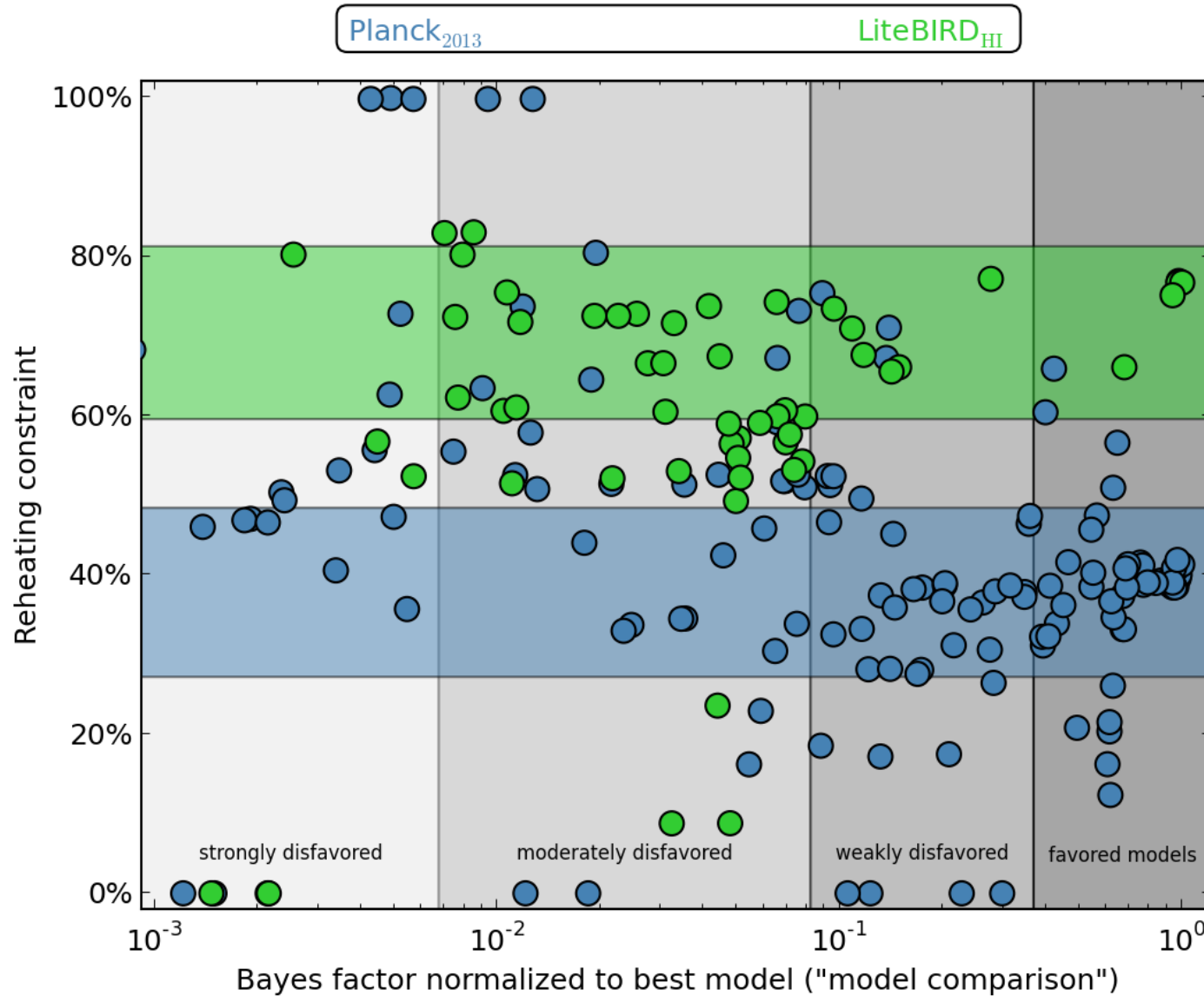
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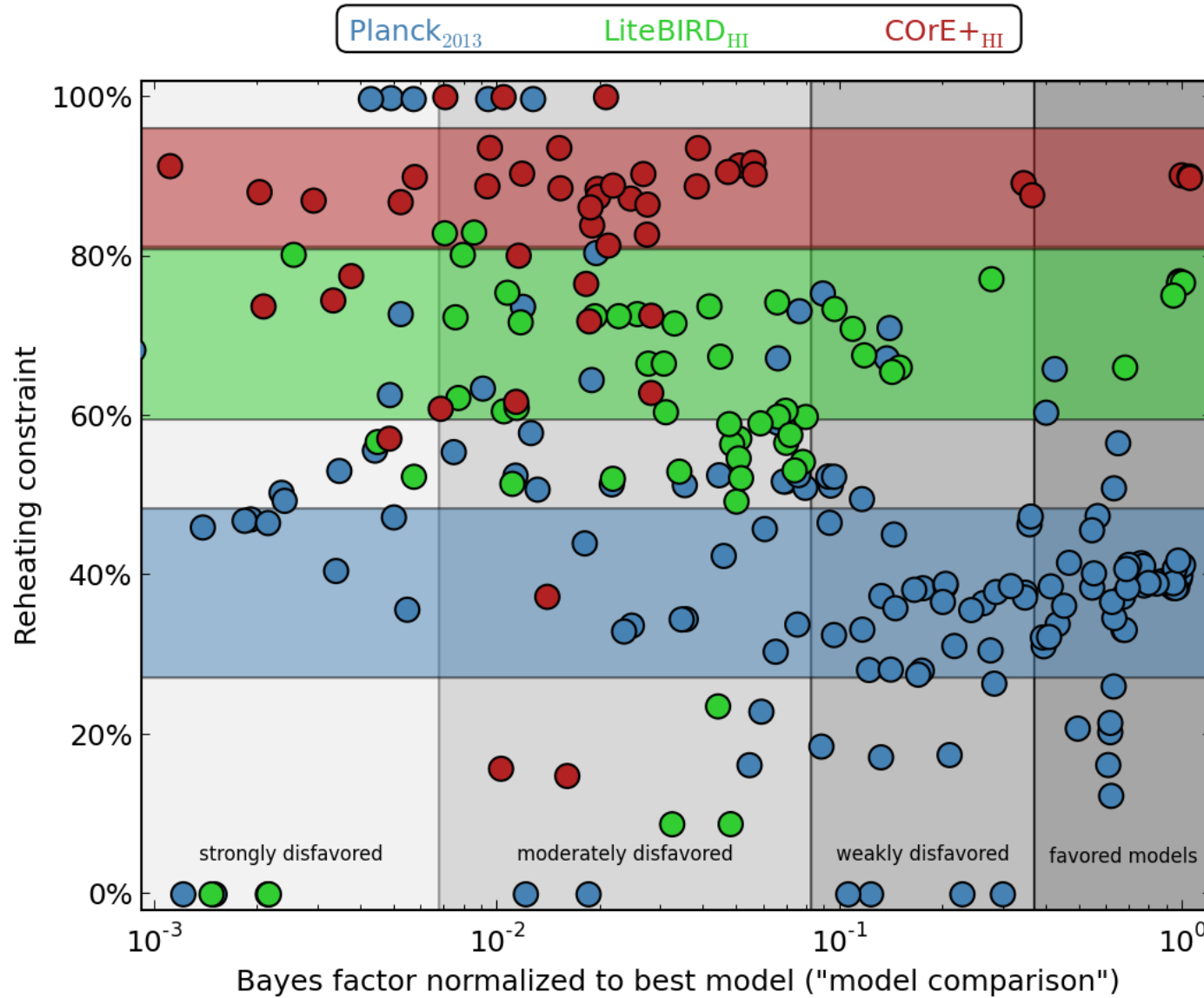
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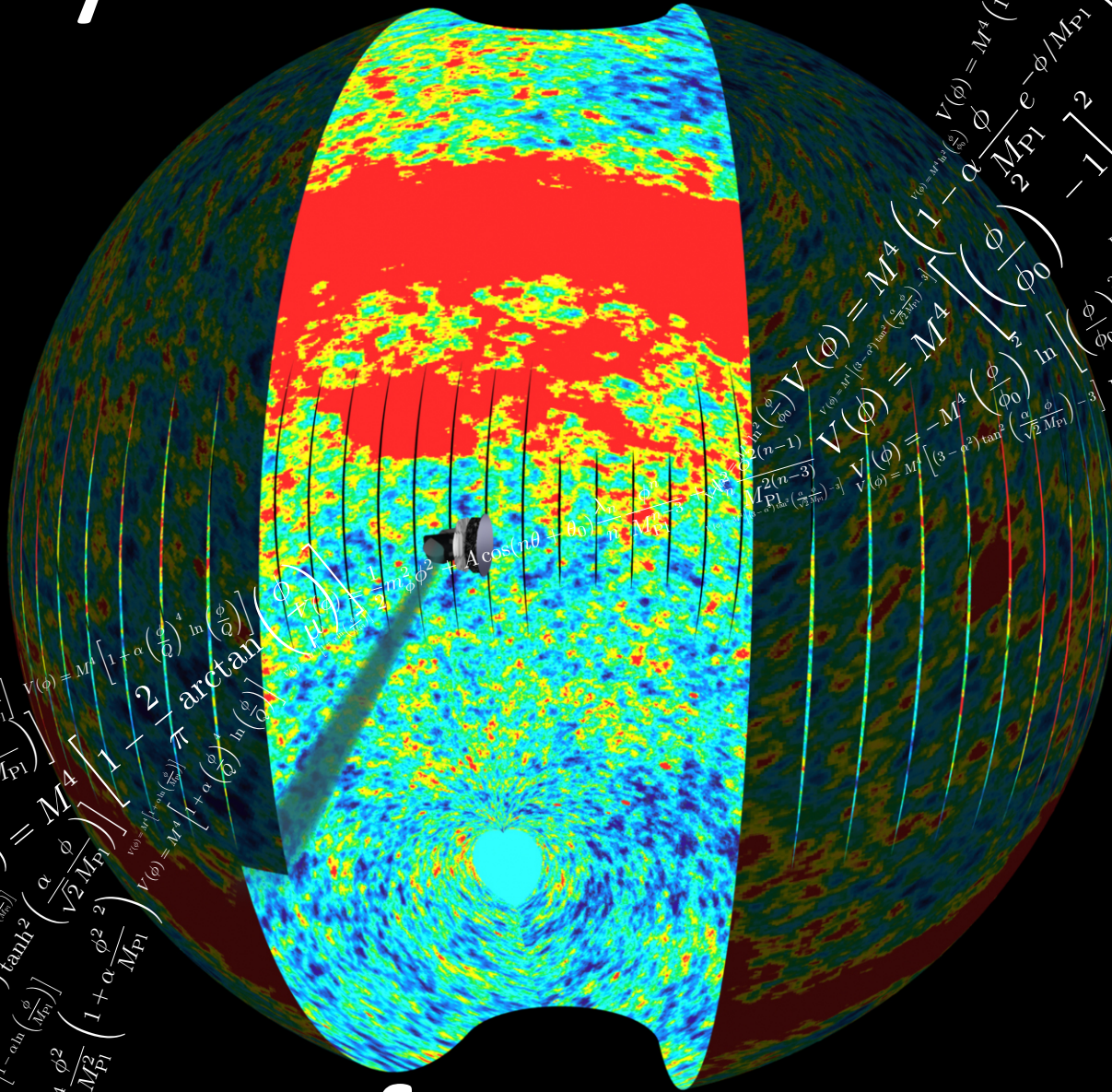
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# Future CMB Missions



# Thank you



# for your attention