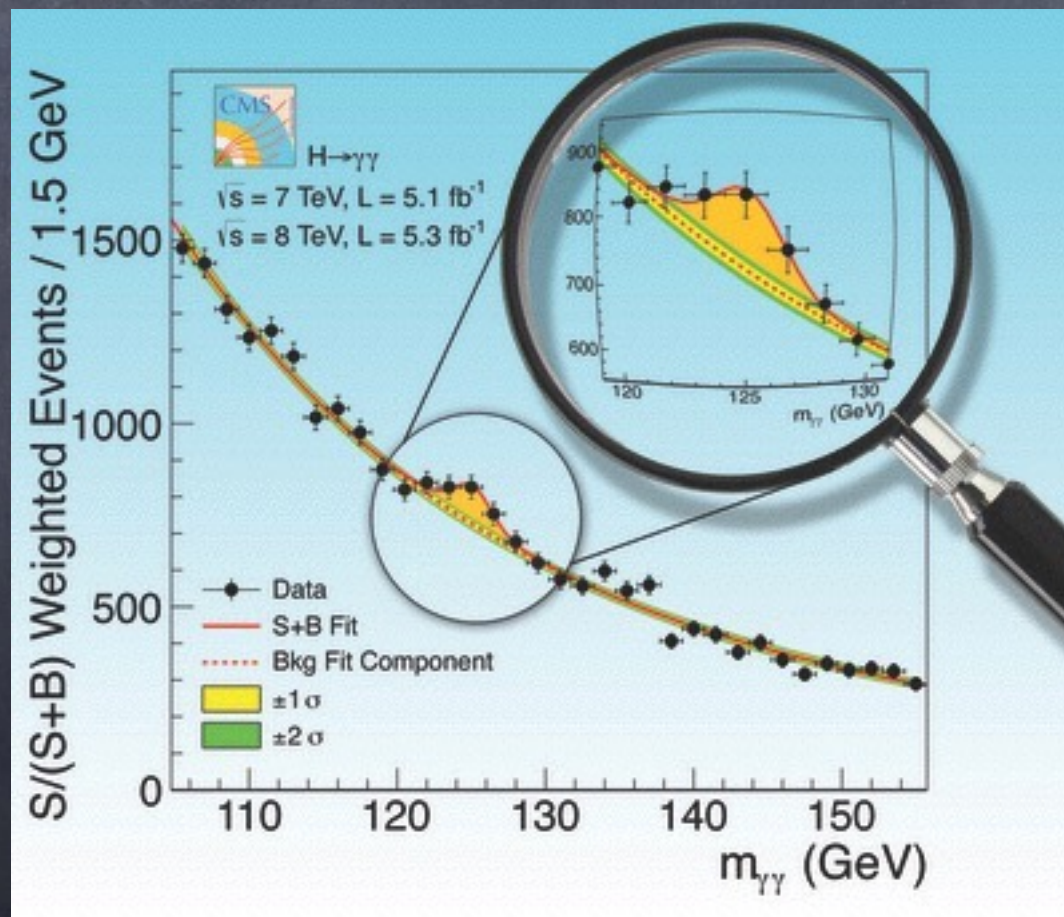


Higgs, Dark Matter, di-photon:
a confining dynamics
to rule them all!

G. Cacciapaglia (IPNL)
@ Montpellier

Monsieur Le Higgs: enfin!

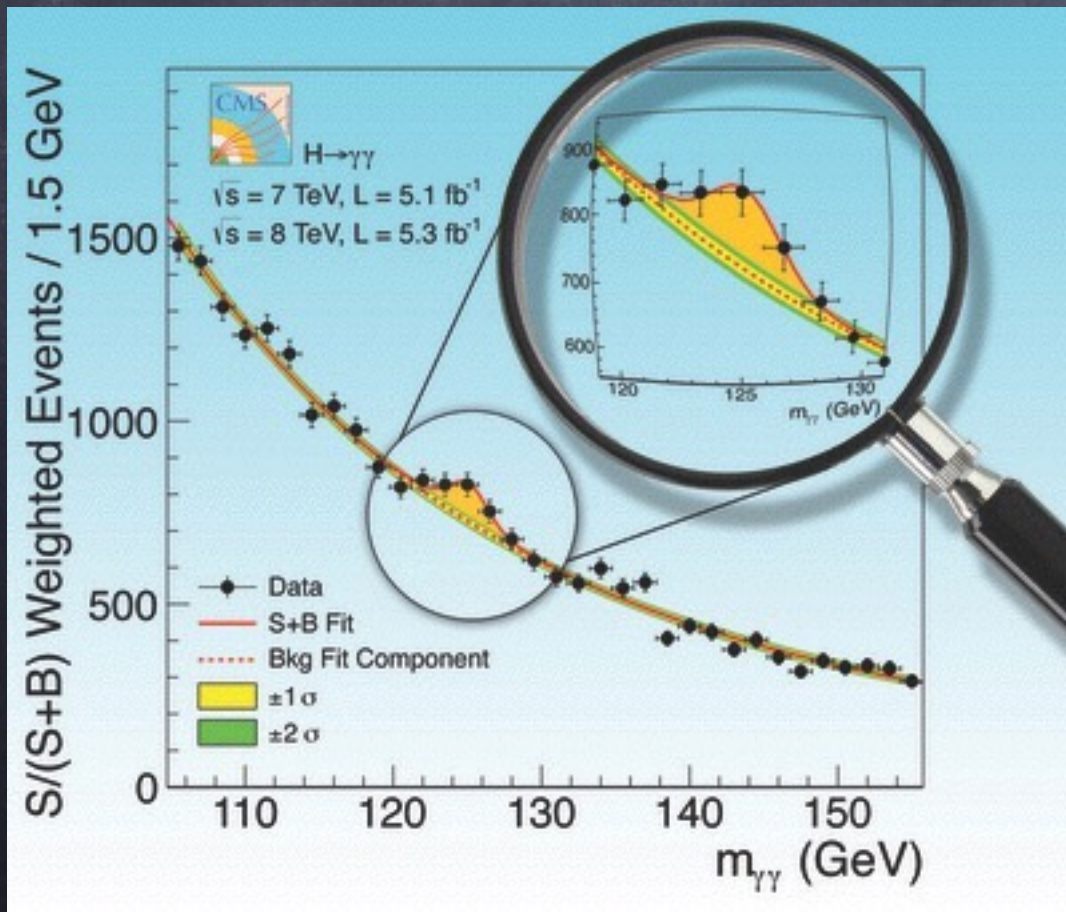


The Standard Model
is finally complete!

(is it?)

Monsieur Le Higgs: enfin!

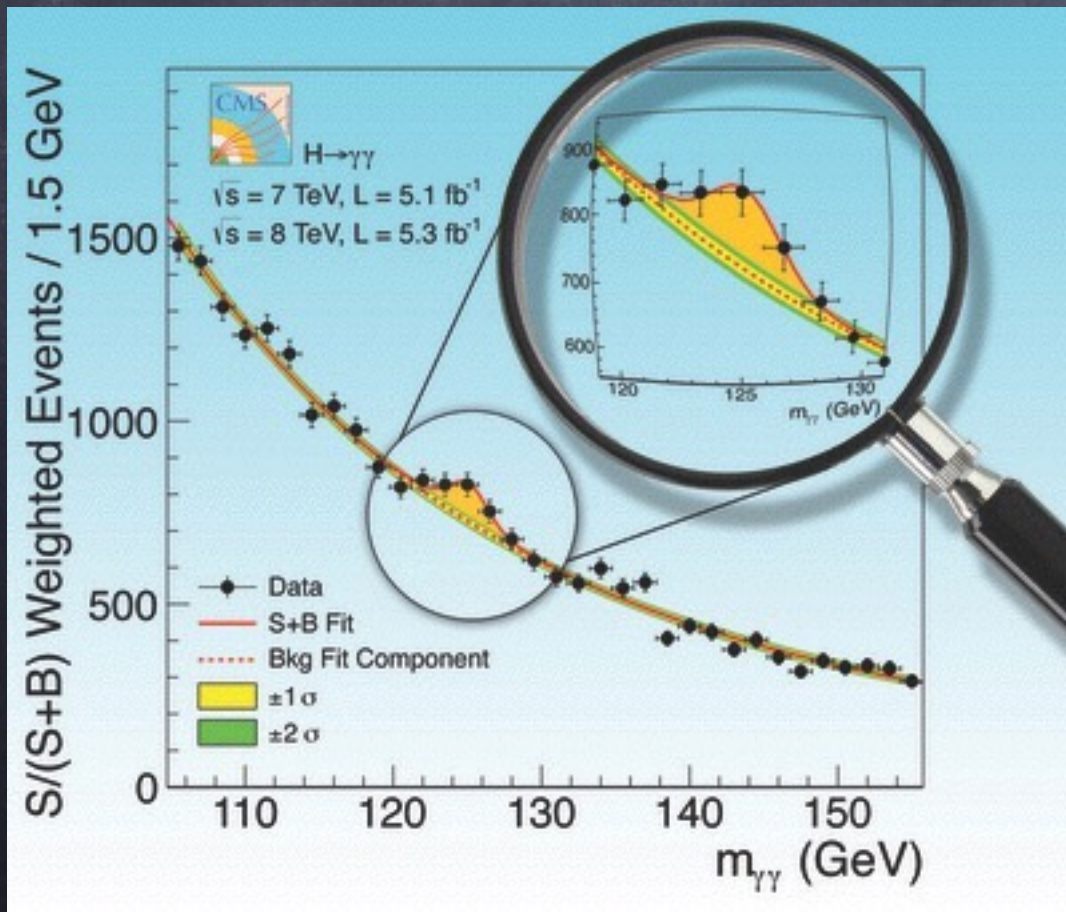
...and the LHC crew
is happy!



Monsieur Le Higgs: enfin!

For a theorist, this is the beginning of a new era!

We have a new toy,



a new probe in the EW sector!

Do we still need BSM?



We have a pretty good idea of the mechanism

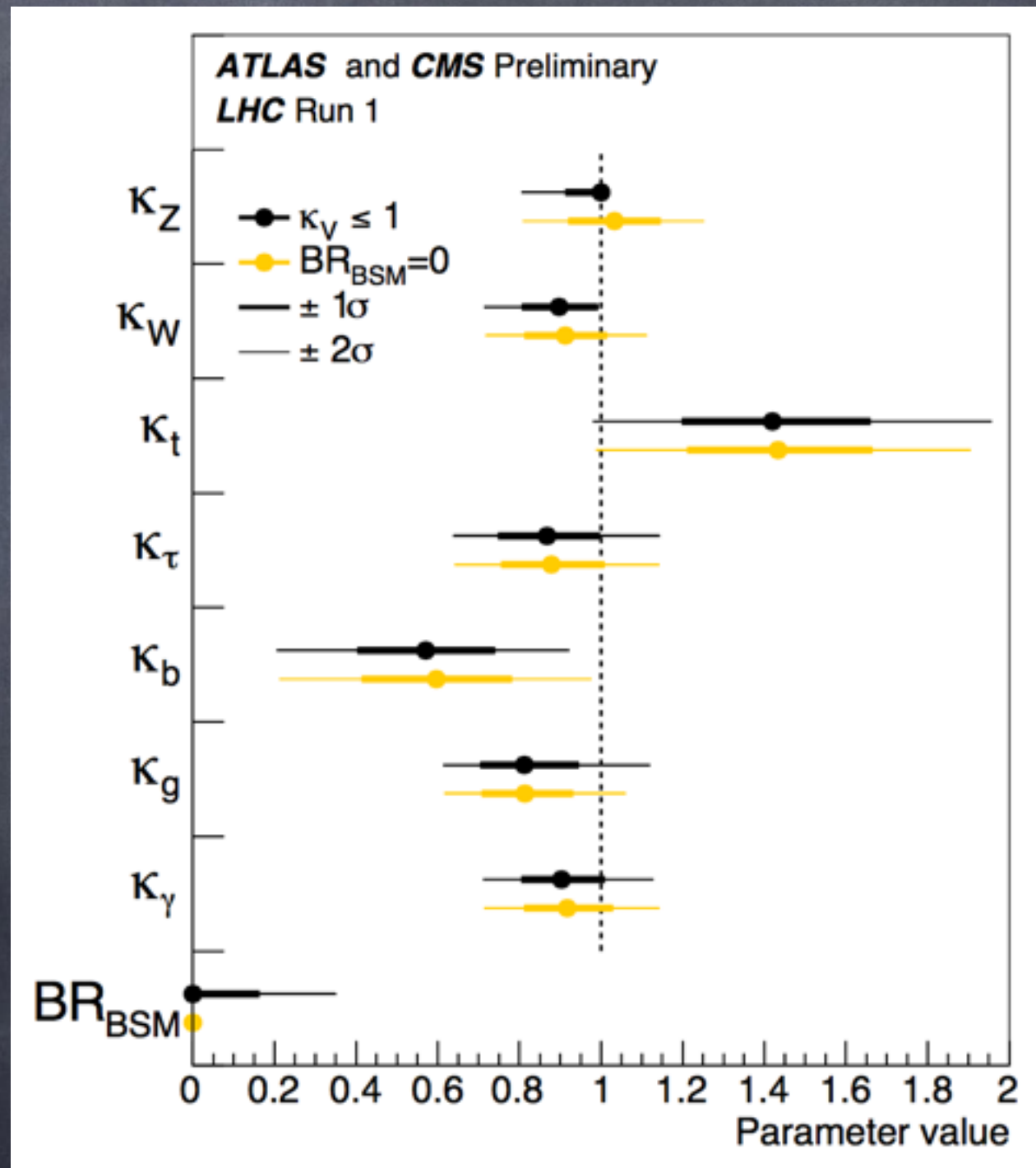
But, we don't know how to protect it:



$$\delta m_h^2 \sim \frac{g^2}{16\pi^2} M_{\text{NPh}}^2$$

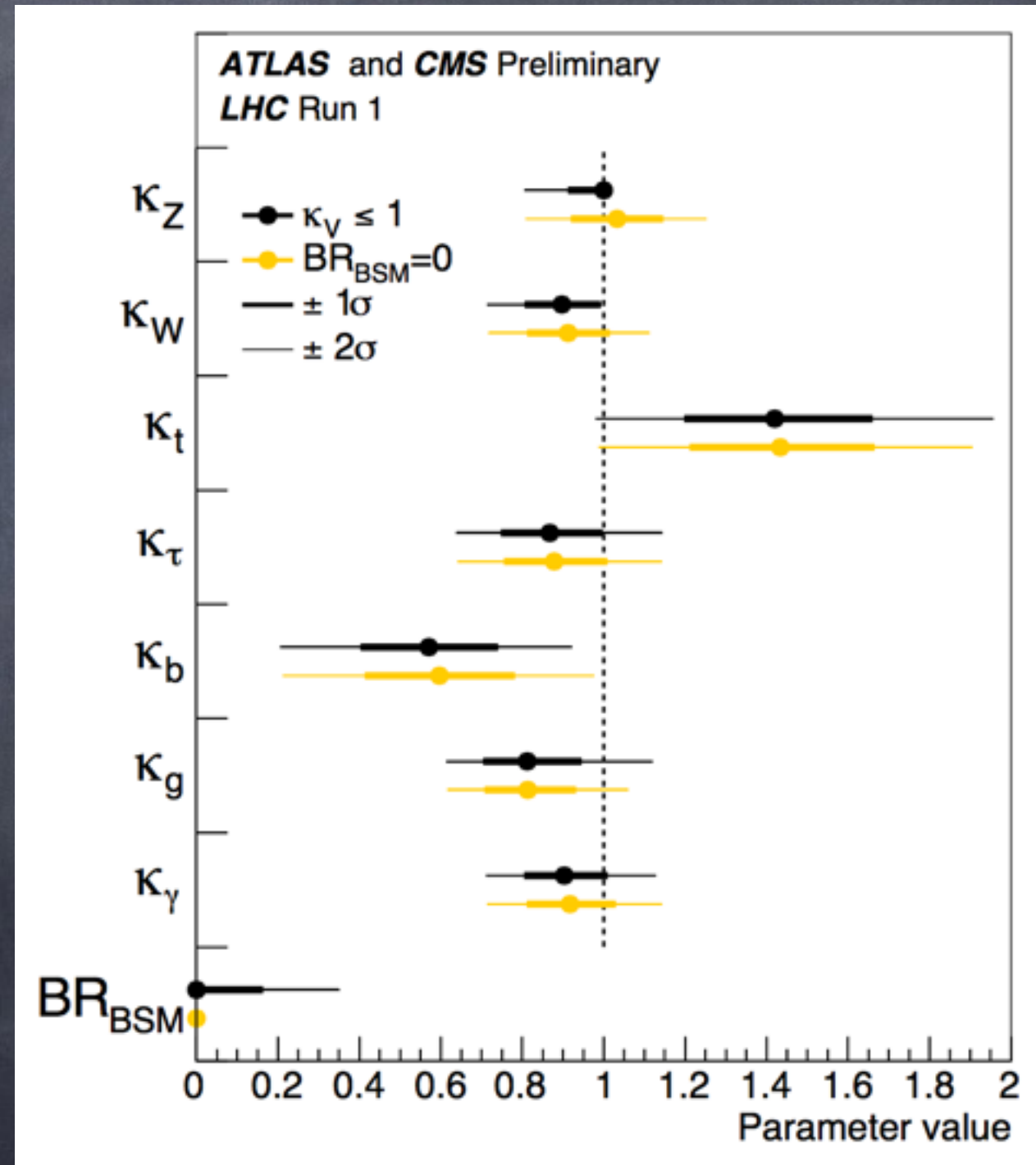
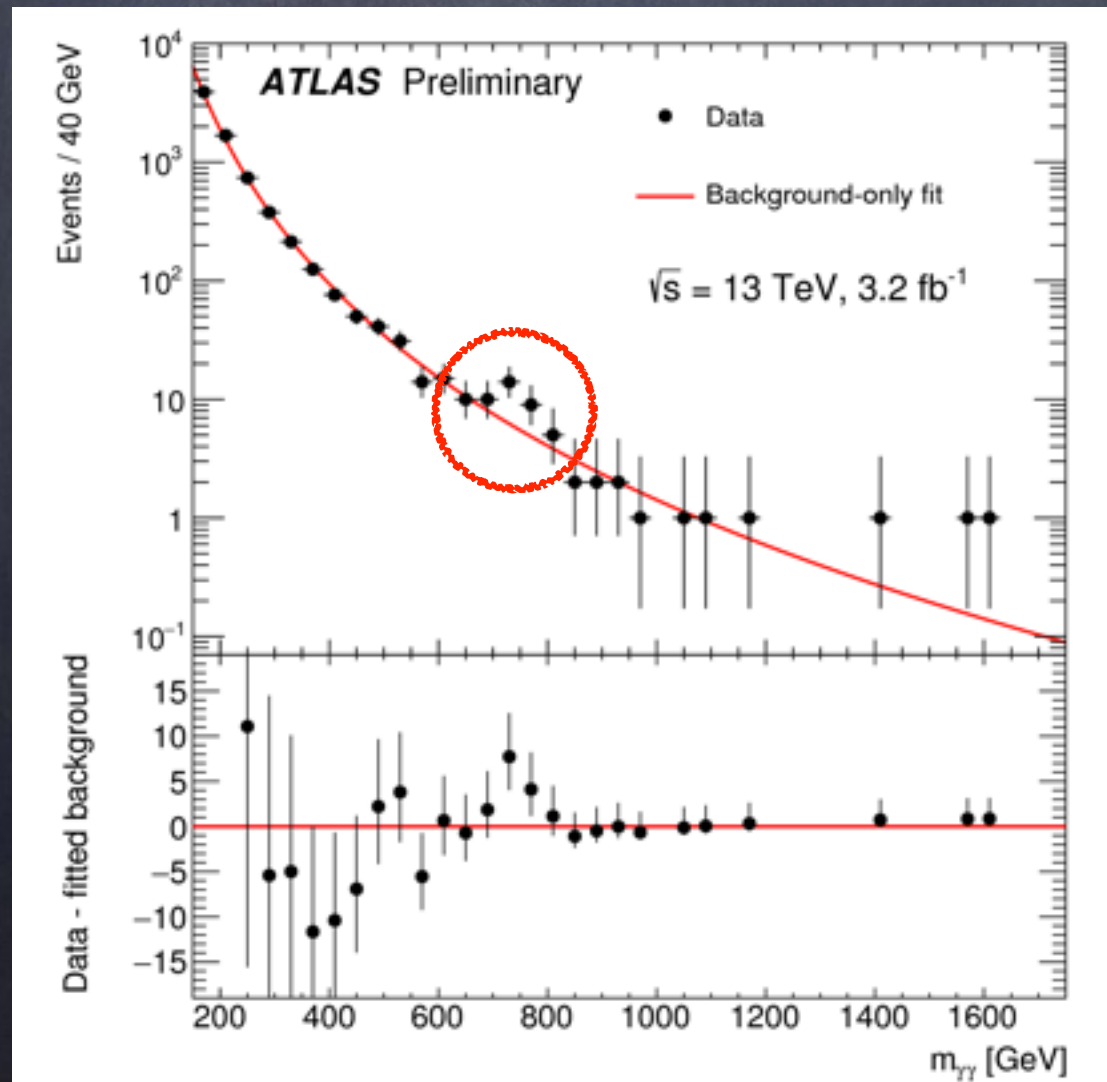
Do we still need BSM?

- Still large uncertainties on the Higgs coupling measurements.
- Ample space for BSM!!!!

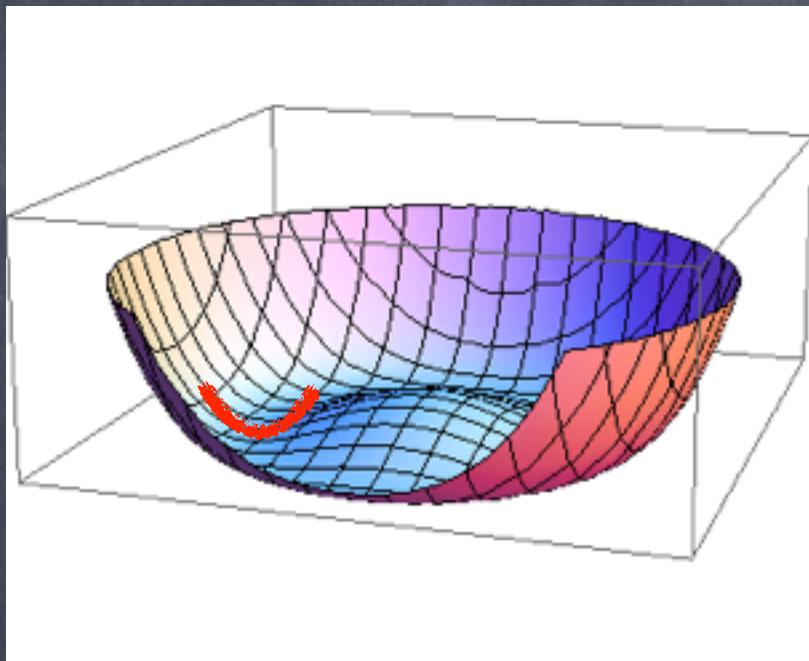


Do we still need BSM?

- Still large uncertainties on the Higgs coupling measurements.
- Ample space for BSM!!!!



Compositeness, and the Higgs boson



Global symmetry of the system:

$$G \rightarrow H$$

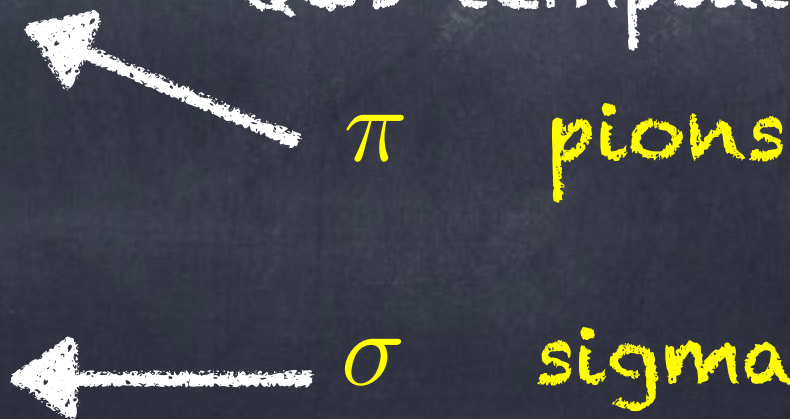


$$SU(2) \times U(1) \rightarrow U(1)_{em}$$

SM gauge symmetry!

- Goldstones include the longitudinal d.o.f. of W and Z
- the Higgs is a heavy bound state (singlet under H)

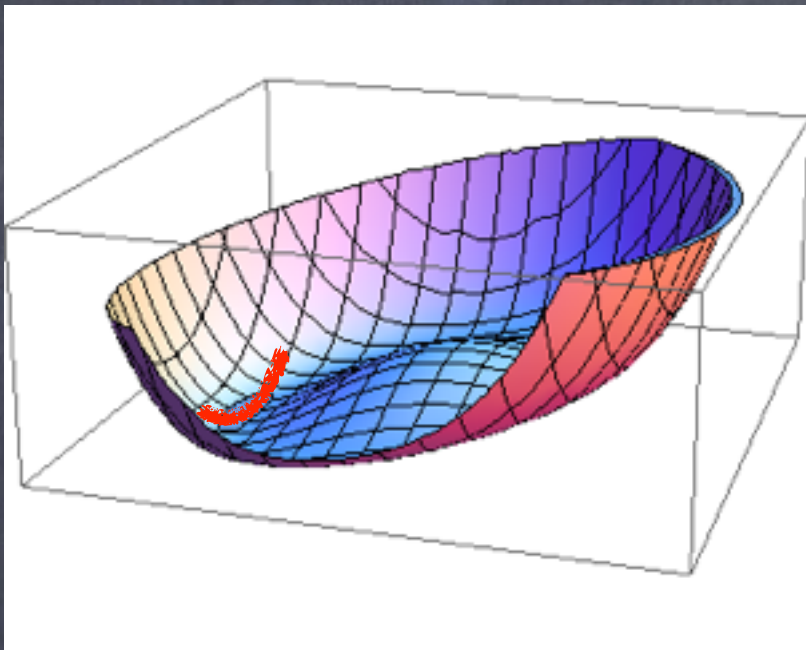
QCD template:



π pions

σ sigma

Compositeness, and the Higgs boson



Global symmetry of the system:

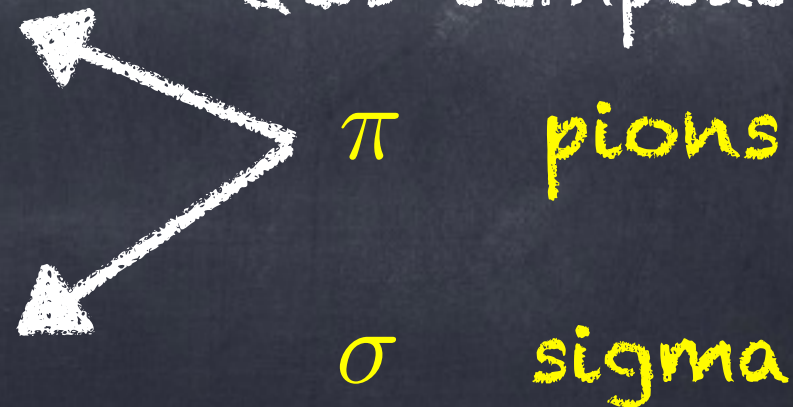
$$G \rightarrow \mathcal{H}$$

$$SU(2) \times U(1) \rightarrow U(1)_{em}$$

SM gauge symmetry!

- Goldstones include the longitudinal d.o.f. of W and Z
- the Higgs is a pseudo-Goldstone

QCD template:

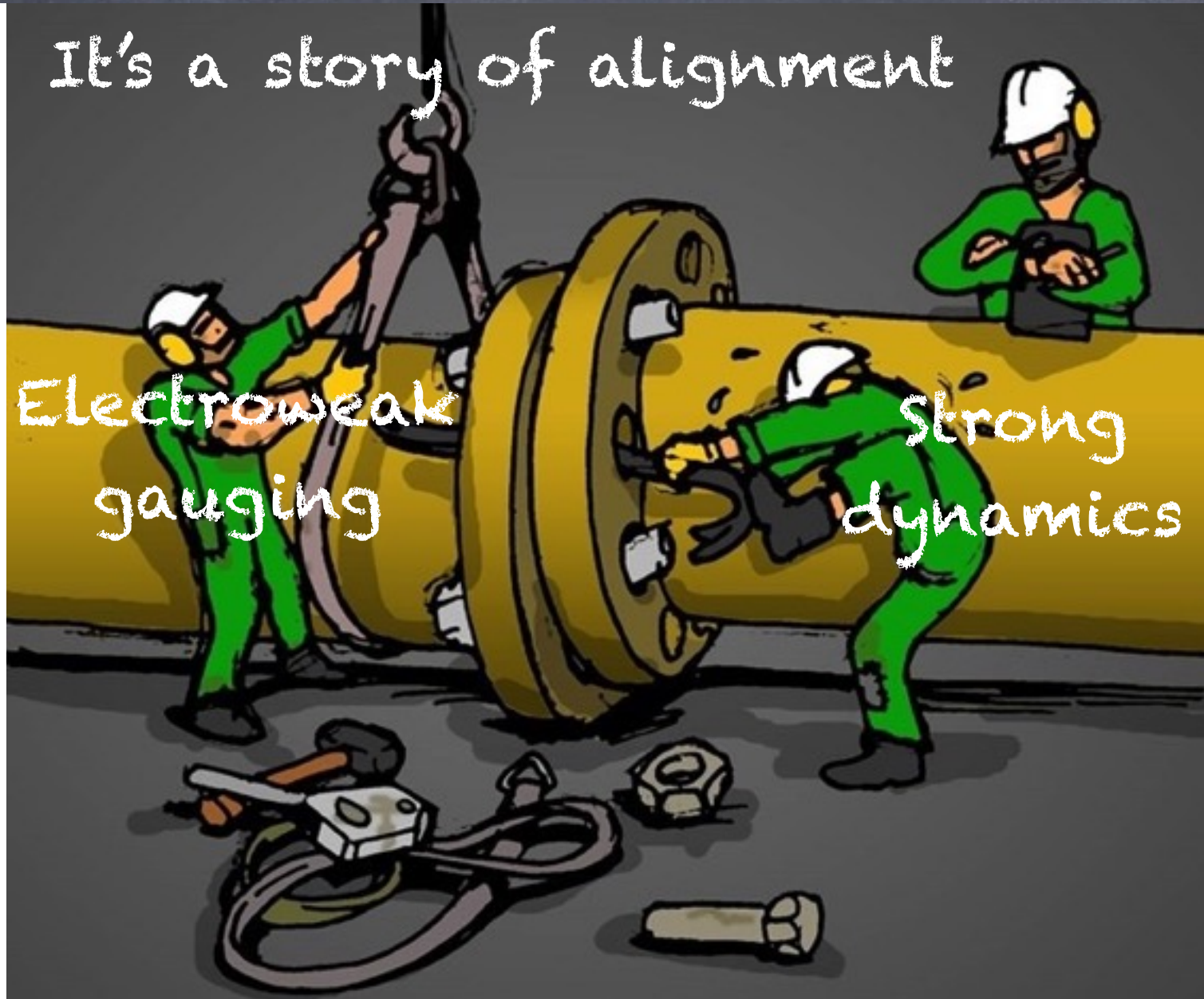


Compositeness, and the Higgs boson

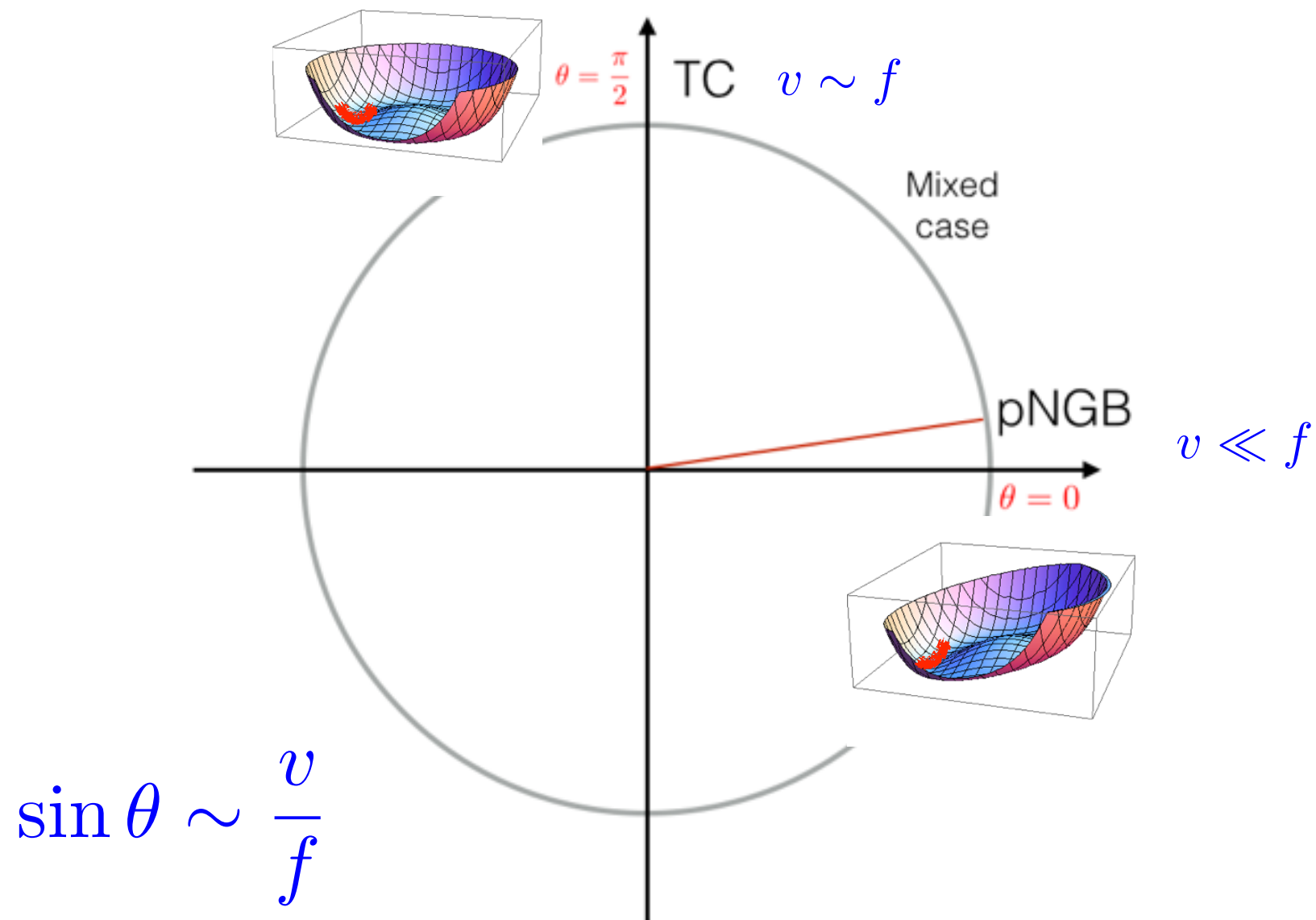
It's a story of alignment

Electroweak
gauging

Strong
dynamics



Compositeness, and the Higgs boson



Higgs: p-Goldstone vs. sigma

👍 Mass is param. lighter than the compositeness scale

👍 The couplings to SM states are naturally close to SM

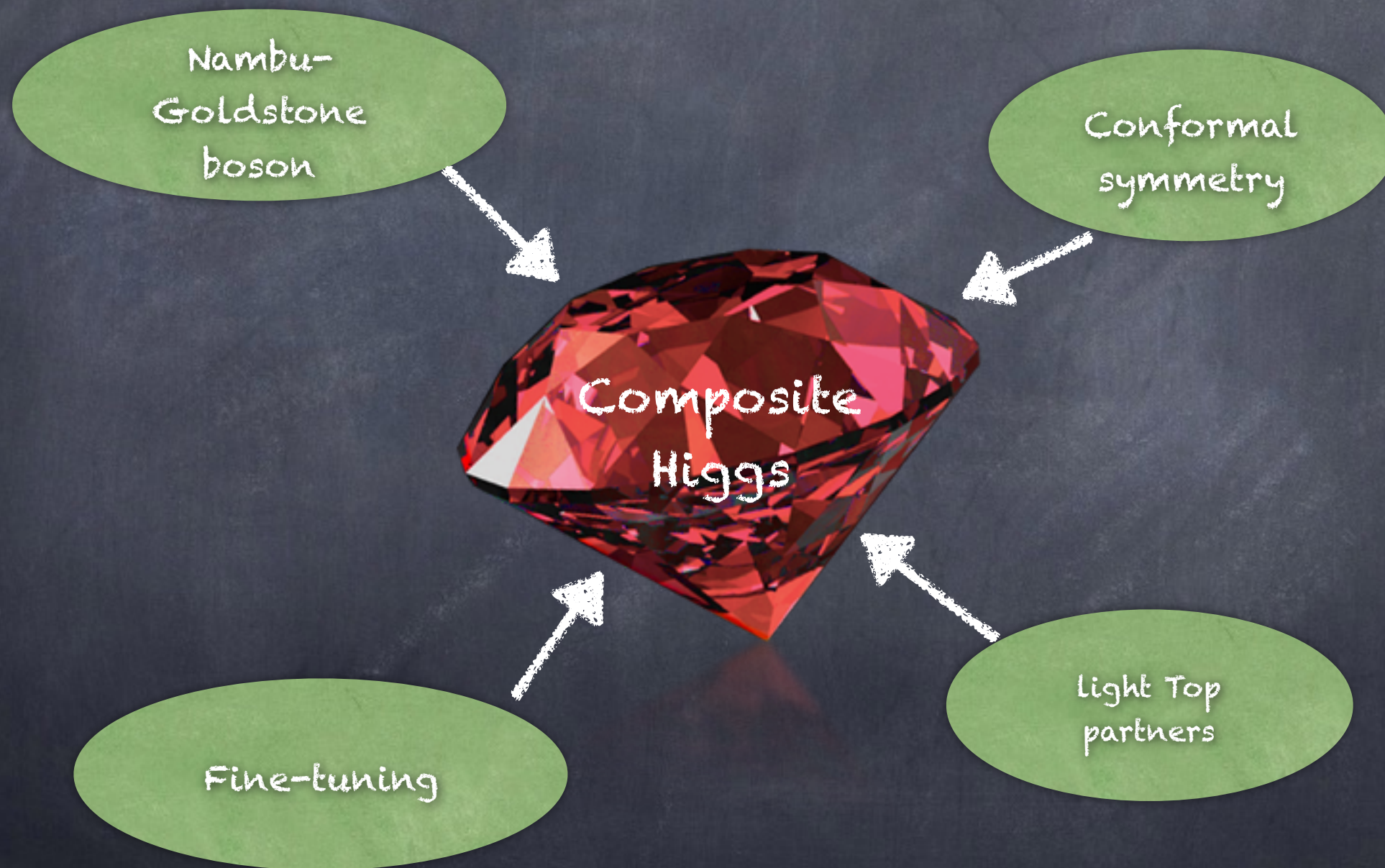
👎 Tuning the tilt in the potential!

👎 Mass is expected to be heavy (close to the rho)

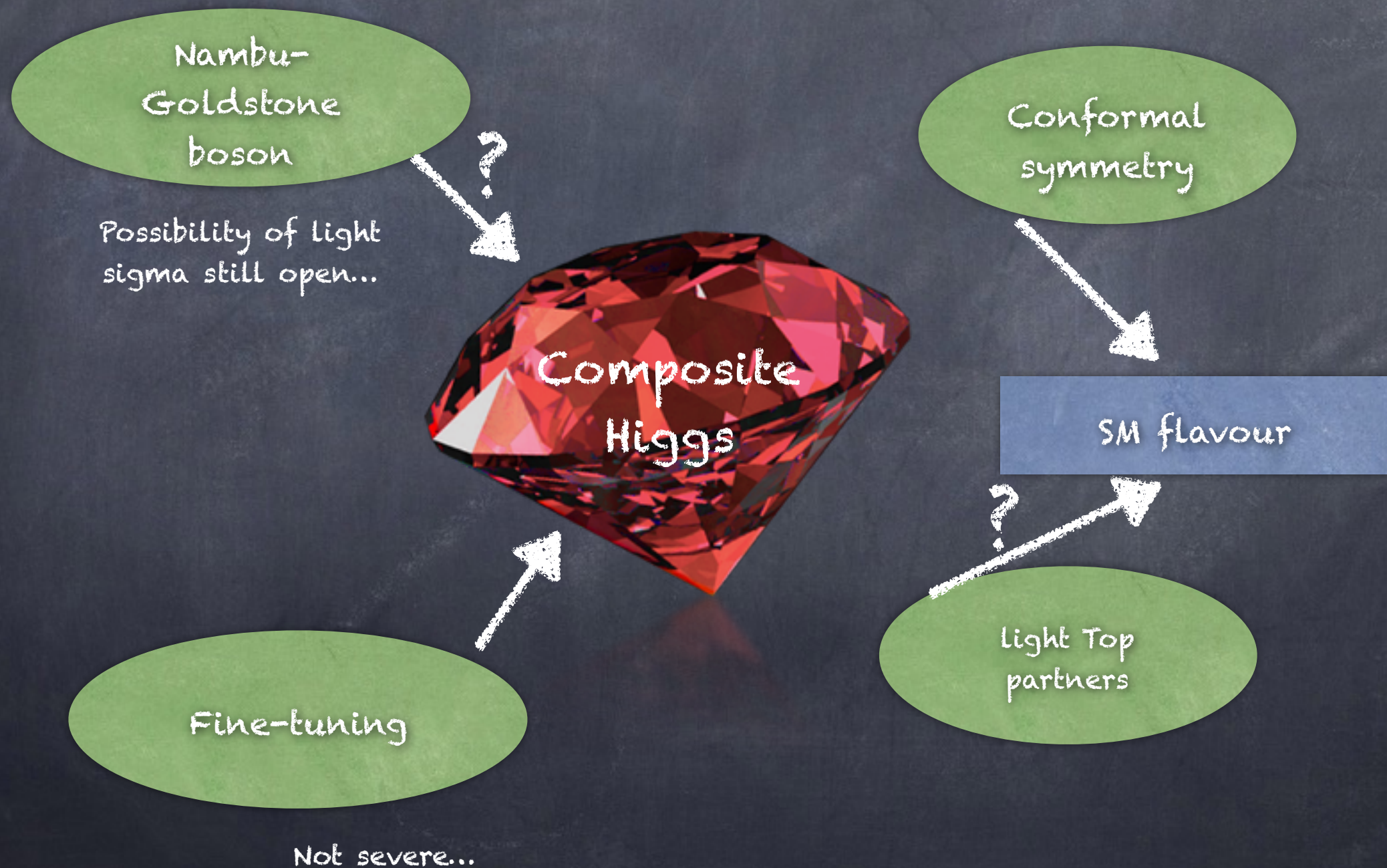
👐 The couplings to SM states are Unknown!

👍 No tuning is necessary!

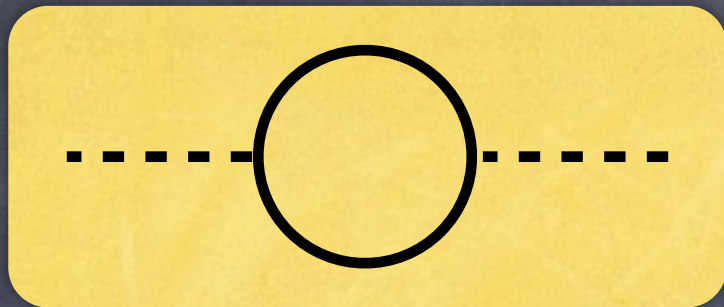
Provocative slide (for the composite alchemists)



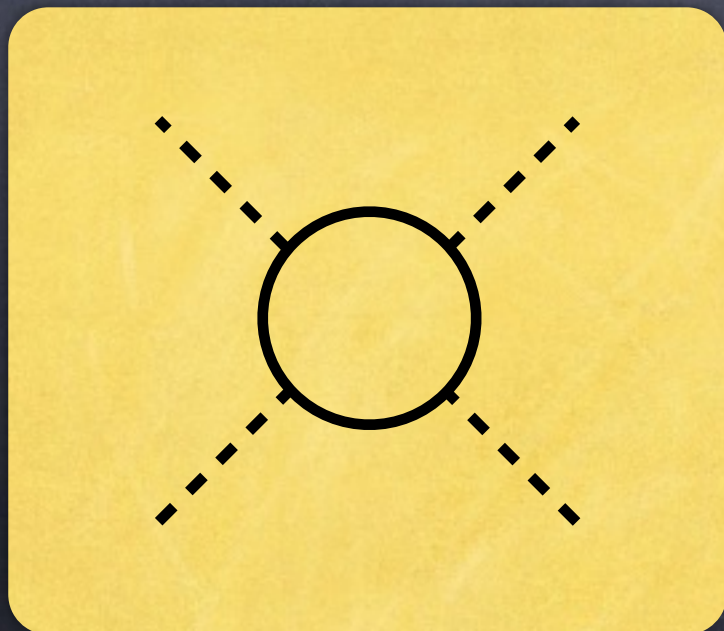
Provocative slide (for the composite alchemists)



Anatomy of the potential

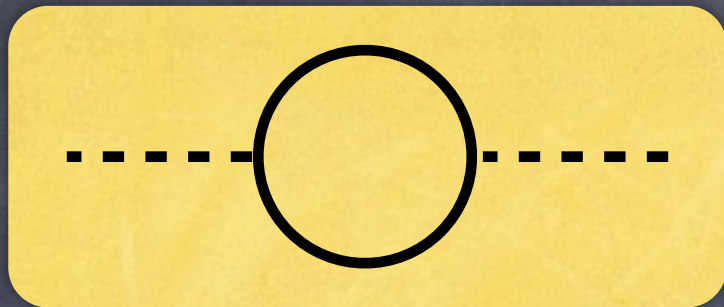


$$\delta m_h^2 \sim \frac{y^2}{16\pi^2} \Lambda^2$$



$$\delta \lambda \sim \frac{y^4}{16\pi^2} \log \Lambda$$

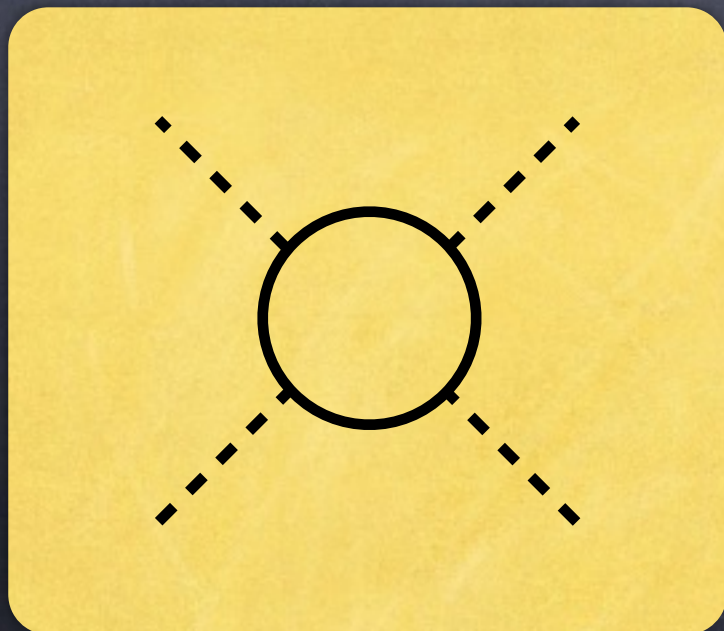
Anatomy of the potential



$$\delta m_h^2 \sim \frac{y^2}{16\pi^2} (4\pi f)^2 \sim y^2 f^2$$

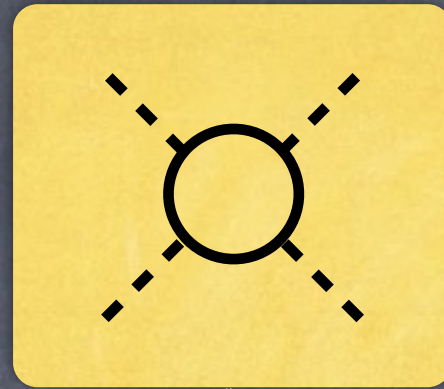
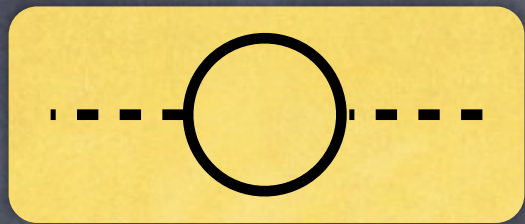
$$\Lambda \sim 4\pi f$$

Strong dynamics
estimate



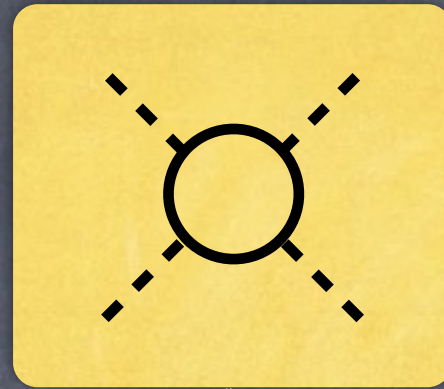
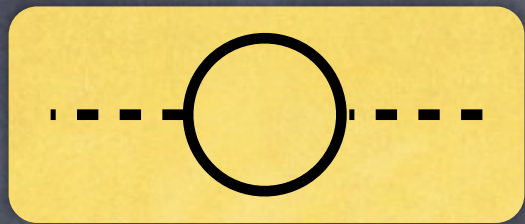
$$\delta\lambda \sim \frac{y^4}{16\pi^2} \log \Lambda$$



Anatomy of the potential



$$V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$$

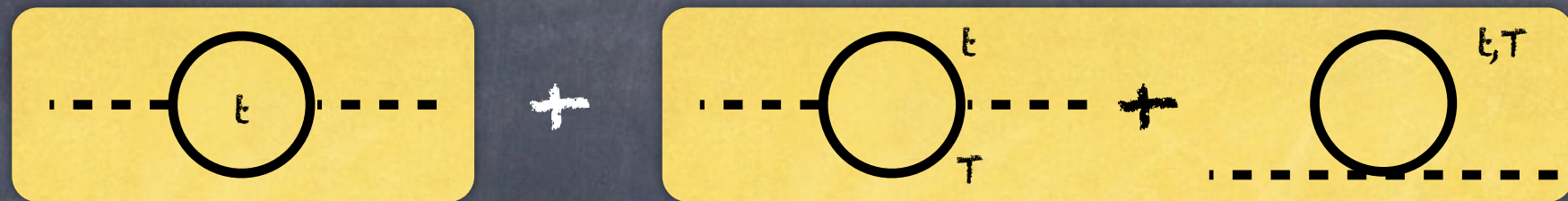
Anatomy of the potential




 $V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$


Minimum: $\theta \sim \frac{\pi}{2}$

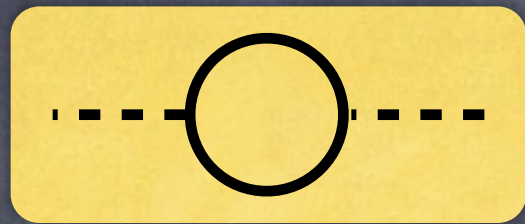
Anatomy of the potential



$$V \sim \alpha \sin^2 \theta + \beta \sin^4 \theta$$

Minimum: $\theta \sim \epsilon$

Anatomy of the potential



Explicit
symmetry
breaking?

$$V \sim \alpha \sin^2 \theta + \gamma \cos \theta$$

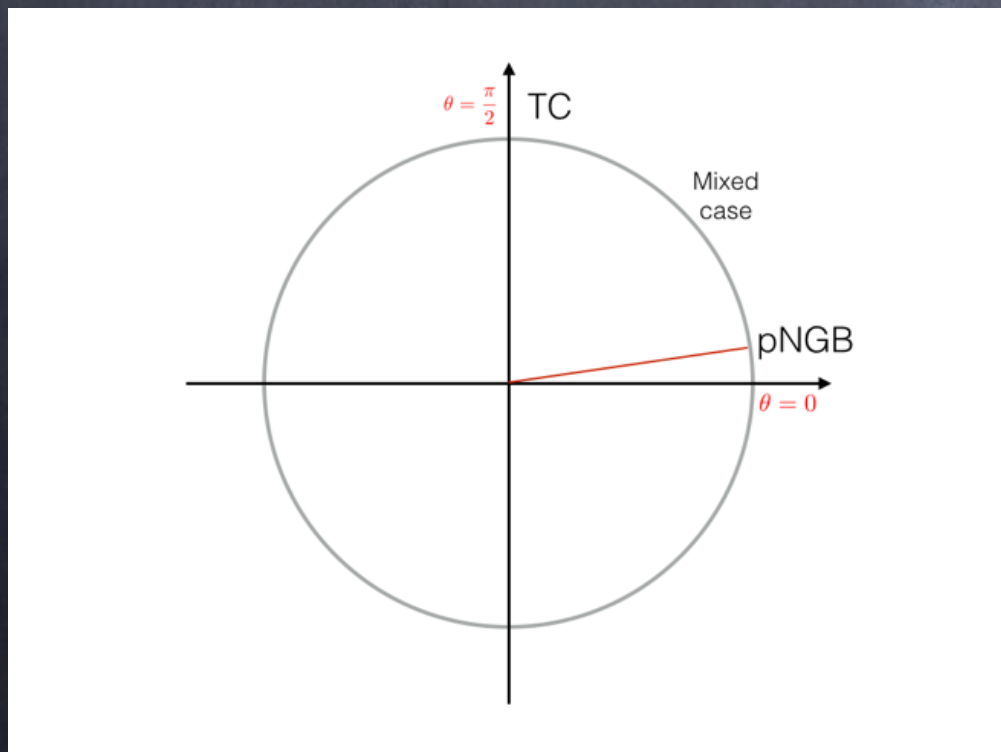
Minimum: $\theta \sim \epsilon$

Possible!

Anatomy of the potential

Higgs mass in the small theta limit:

$$m_h \sim y f \sin \theta \sim y v_{SM}$$



Naturally in the
right ballpark,
without further
fine tuning!

Composite Higgses: which model?

\mathcal{G}	\mathcal{H}	C	N_G	$\mathbf{r}_\mathcal{H} = \mathbf{r}_{\text{SU}(2) \times \text{SU}(2)} (\mathbf{r}_{\text{SU}(2) \times \text{U}(1)})$	Ref.
SO(5)	SO(4)	✓	4	$\mathbf{4} = (\mathbf{2}, \mathbf{2})$	[11]
SU(3) × U(1)	SU(2) × U(1)		5	$\mathbf{2}_{\pm 1/2} + \mathbf{1}_0$	[10, 35]
SU(4)	Sp(4)	✓	5	$\mathbf{5} = (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$	[29, 47, 64]
SU(4)	[SU(2)] ² × U(1)	✓*	8	$(\mathbf{2}, \mathbf{2})_{\pm 2} = 2 \cdot (\mathbf{2}, \mathbf{2})$	[65]
SO(7)	SO(6)	✓	6	$\mathbf{6} = 2 \cdot (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$	—
SO(7)	G ₂	✓*	7	$\mathbf{7} = (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$	[66]
SO(7)	SO(5) × U(1)	✓*	10	$\mathbf{10}_0 = (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$	—
SO(7)	[SU(2)] ³	✓*	12	$(\mathbf{2}, \mathbf{2}, \mathbf{3}) = 3 \cdot (\mathbf{2}, \mathbf{2})$	—
Sp(6)	Sp(4) × SU(2)	✓	8	$(\mathbf{4}, \mathbf{2}) = 2 \cdot (\mathbf{2}, \mathbf{2})$	[65]
SU(5)	SU(4) × U(1)	✓*	8	$\mathbf{4}_{-5} + \bar{\mathbf{4}}_{+5} = 2 \cdot (\mathbf{2}, \mathbf{2})$	[67]
SU(5)	SO(5)	✓*	14	$\mathbf{14} = (\mathbf{3}, \mathbf{3}) + (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{1})$	[9, 47, 49]
SO(8)	SO(7)	✓	7	$\mathbf{7} = 3 \cdot (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$	—
SO(9)	SO(8)	✓	8	$\mathbf{8} = 2 \cdot (\mathbf{2}, \mathbf{2})$	[67]
SO(9)	SO(5) × SO(4)	✓*	20	$(\mathbf{5}, \mathbf{4}) = (\mathbf{2}, \mathbf{2}) + (\mathbf{1} + \mathbf{3}, \mathbf{1} + \mathbf{3})$	[34]
[SU(3)] ²	SU(3)		8	$\mathbf{8} = \mathbf{1}_0 + \mathbf{2}_{\pm 1/2} + \mathbf{3}_0$	[8]
[SO(5)] ²	SO(5)	✓*	10	$\mathbf{10} = (\mathbf{1}, \mathbf{3}) + (\mathbf{3}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$	[32]
SU(4) × U(1)	SU(3) × U(1)		7	$\mathbf{3}_{-1/3} + \bar{\mathbf{3}}_{+1/3} + \mathbf{1}_0 = 3 \cdot \mathbf{1}_0 + \mathbf{2}_{\pm 1/2}$	[35, 41]
SU(6)	Sp(6)	✓*	14	$\mathbf{14} = 2 \cdot (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{3}) + 3 \cdot (\mathbf{1}, \mathbf{1})$	[30, 47]
[SO(6)] ²	SO(6)	✓*	15	$\mathbf{15} = (\mathbf{1}, \mathbf{1}) + 2 \cdot (\mathbf{2}, \mathbf{2}) + (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3})$	[36]

Table 1: Symmetry breaking patterns $\mathcal{G} \rightarrow \mathcal{H}$ for Lie groups. The third column denotes whether the breaking pattern incorporates custodial symmetry. The fourth column gives the dimension N_G of the coset, while the fifth contains the representations of the GB's under \mathcal{H} and $\text{SO}(4) \cong \text{SU}(2)_L \times \text{SU}(2)_R$ (or simply $\text{SU}(2)_L \times \text{U}(1)_Y$ if there is no custodial symmetry). In case of more than two SU(2)'s in \mathcal{H} and several different possible decompositions we quote the one with largest number of bi-doublets.

The FCD approach

G.C., F.Sannino

1402.0233

- Define a confining gauge group (GTC)
- Add in N fermions charged under the confining group GTC
- Assign SM quantum numbers to the fermions
- Couple them to SM fermions

The FCD approach

coset	GTC	TF	doublets	pNGBs	
$SU(4)/Sp(4)$	$Sp(2N)$	fund	1	5	← Minimal!
$SU(5)/SO(5)$	$SU(4)$	6	1	14	Dugan, Georgi, Kaplan 1985!!!
$SU(4) \times SU(4) / SU(4)$	$SU(N)$	fund	2	15	G.C., T.Ma 1508.07014
$SU(6)/Sp(6)$	$Sp(2N)$	fund	2	14	G.C., M.Lespinasse in prep.

The minimal case

T.Ryttov, F.Sannino 0809.0713
Galloway, Evans, Luty, Tacchi 1001.1361

Anti-symmetric

$$\langle \psi^i \psi^j \rangle = 6_{\text{SU}(4)} \rightarrow 5_{\text{Sp}(4)} \oplus 1_{\text{Sp}(4)}$$

$\text{Sp}(4)$ contains a $\text{SO}(4)$ subgroup:
identify with custodial symmetry!

Pions: $5_{\text{Sp}(4)} \rightarrow (2, 2) \oplus (1, 1)$

The minimal case

Galloway, Evans, Luty, Tacchi 1001.1361
G.C., F.Sannino 1402.0233

$$V_m = C_m f^4 \text{Tr}(\Sigma_B \cdot \Sigma)$$
$$\sim C_m \left(-4f^4 c_\theta + \sqrt{2} f^3 s_\theta h + \frac{1}{4} f^2 c_\theta (h^2 + \eta^2) + \dots \right)$$

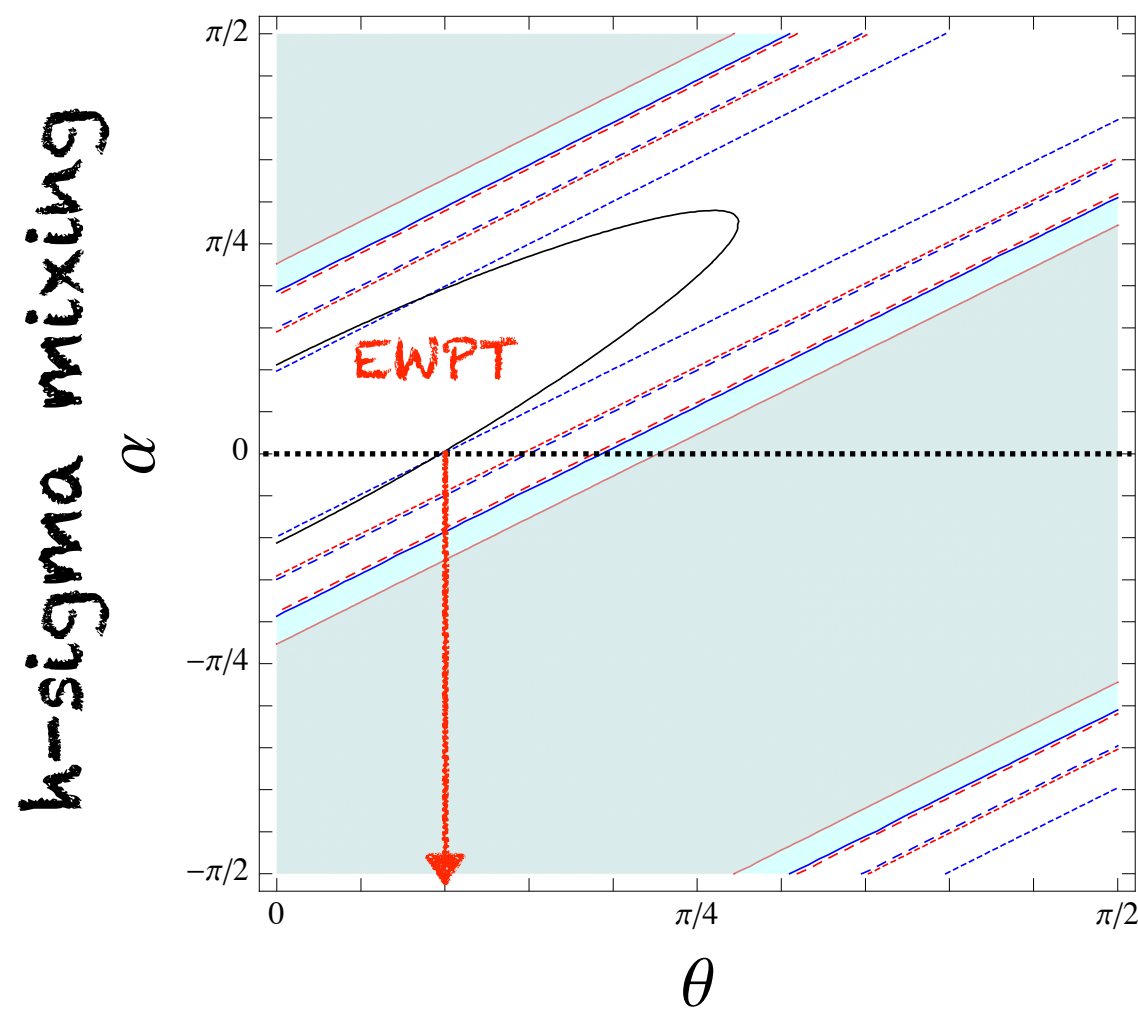
$$V_{top} = -C_t y_t'^2 f^4 \sum_{\alpha=1}^2 [\text{Tr}(P^\alpha \Sigma)]^2$$
$$\sim -C_t y_t'^2 \left[f^4 s_\theta^2 + \frac{1}{\sqrt{2}} f^3 c_\theta s_\theta h + \frac{1}{8} f^2 (c_{2\theta} h^2 - s_\theta^2 \eta^2) + \dots \right]$$

$$m_h \sim \frac{\sqrt{C_t}}{2} y_t f \sin \theta \sim \frac{\sqrt{C_t}}{2} m_t \quad m_\eta \sim \frac{\sqrt{C_t}}{2} y_t f \sim \frac{m_h}{\sin \theta}$$

$C_t \sim 2$ reproduces data!

The minimal case

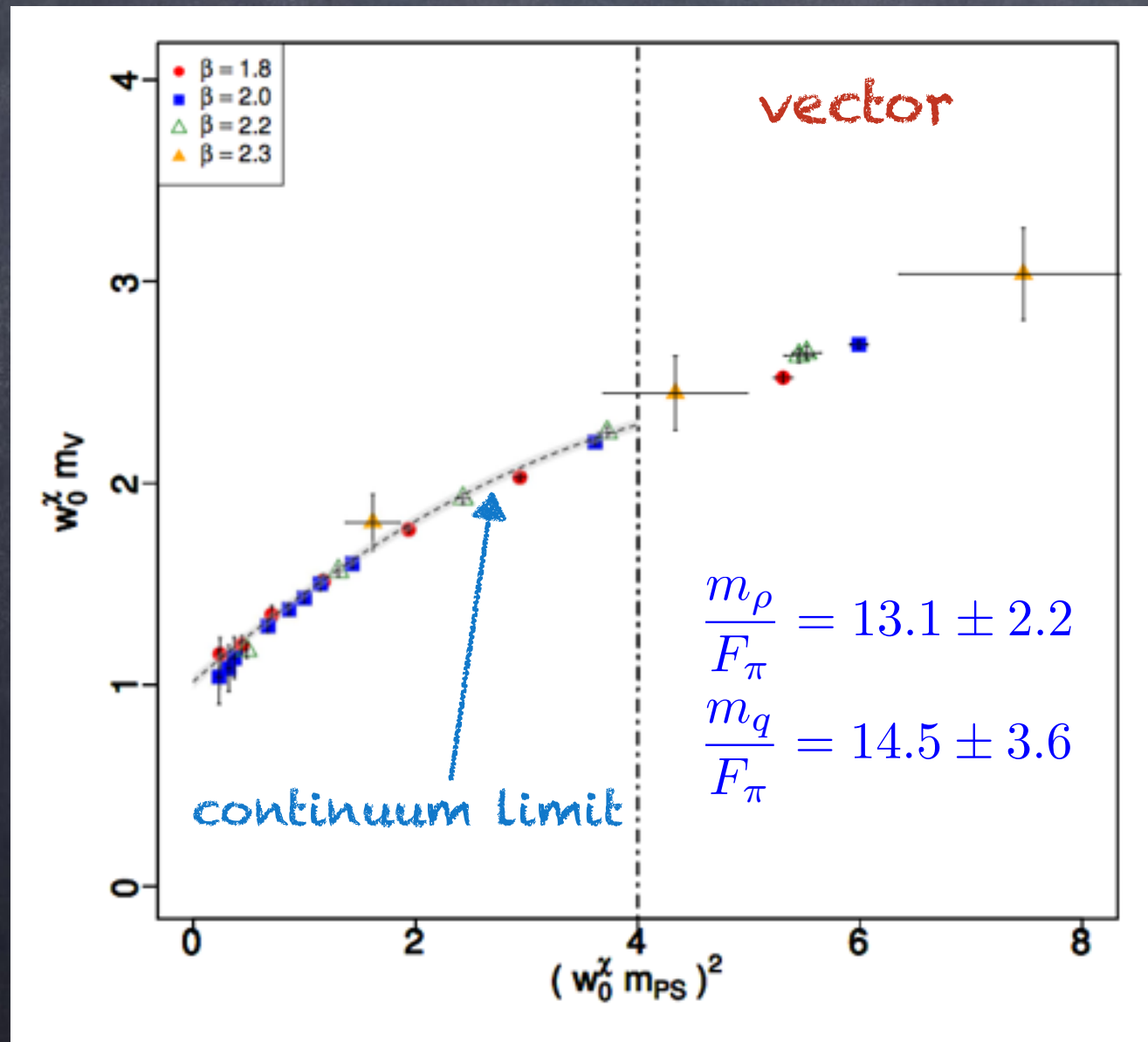
Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718



$$\sin \theta \lesssim 0.2$$

The minimal case

Lattice results:



$$m_a = \frac{3.6 \pm 0.9 \text{ TeV}}{\sin \theta} \gtrsim 18 \text{ TeV}$$

$$m_\rho = \frac{3.2 \pm 0.5 \text{ TeV}}{\sin \theta} \gtrsim 16 \text{ TeV}$$

$$m_\sigma \sim ???$$

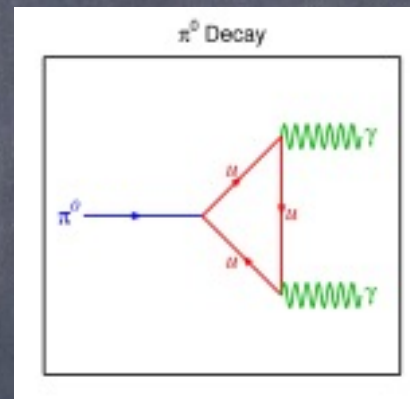
$$m_\eta \sim \frac{m_h}{\sin \theta} \gtrsim 600 \text{ GeV}$$

$$m_h = 125 \text{ GeV}$$

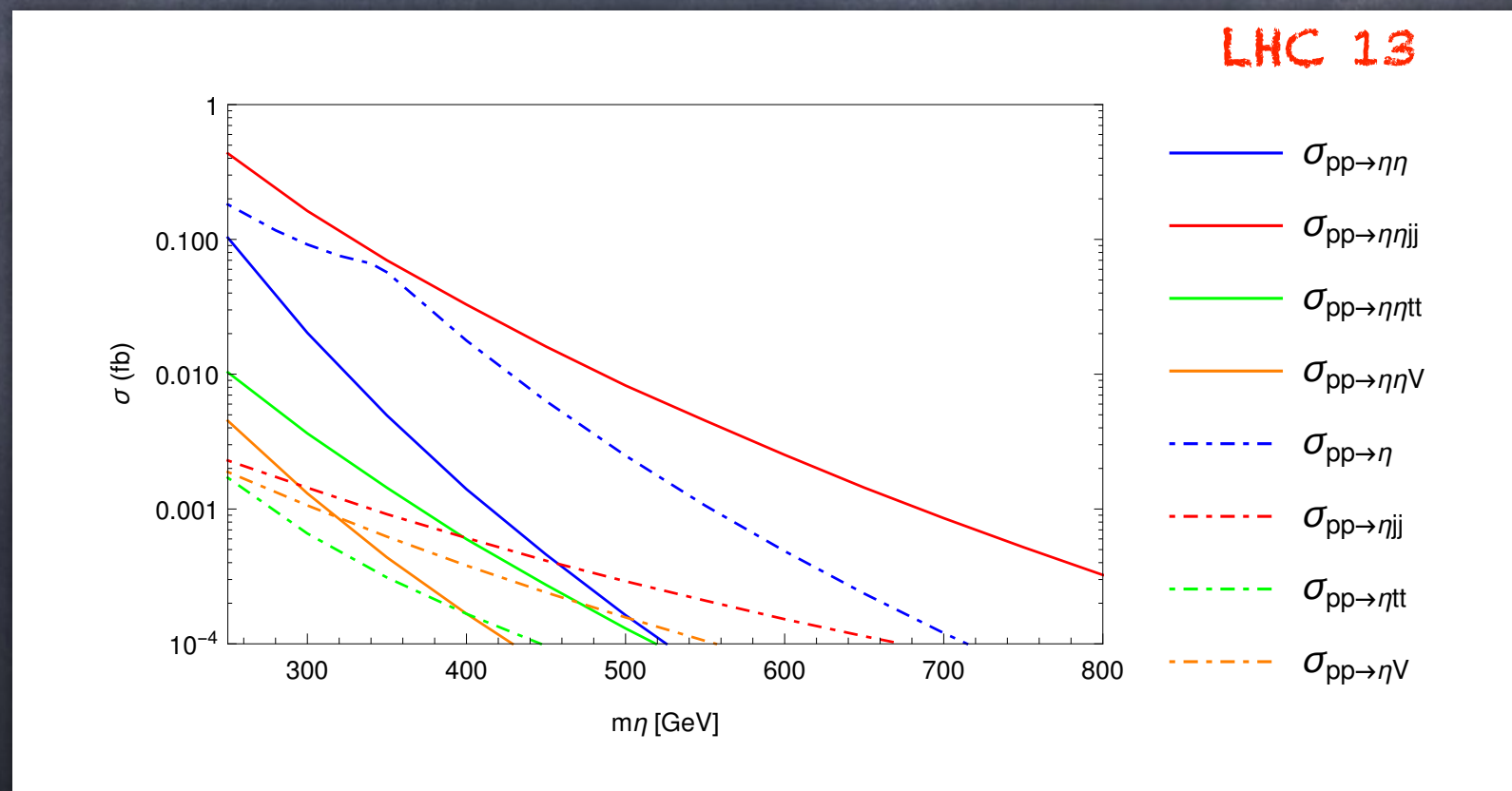
The minimal case

Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718

Singlet cannot be
Dark Matter!



Anomalous coupling to:
WW, ZZ, Z gamma



A composite 2HDM

G.C., T.Ma
1508.07014

	SU(N)	SU(2) _L	U(1) _Y
$\psi_L = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$	□	2	0
$\psi_R = \begin{pmatrix} \psi_3 \\ \psi_4 \end{pmatrix}$	□	1 1	1/2 -1/2

$$SU(4) \times SU(4) \rightarrow SU(4)$$

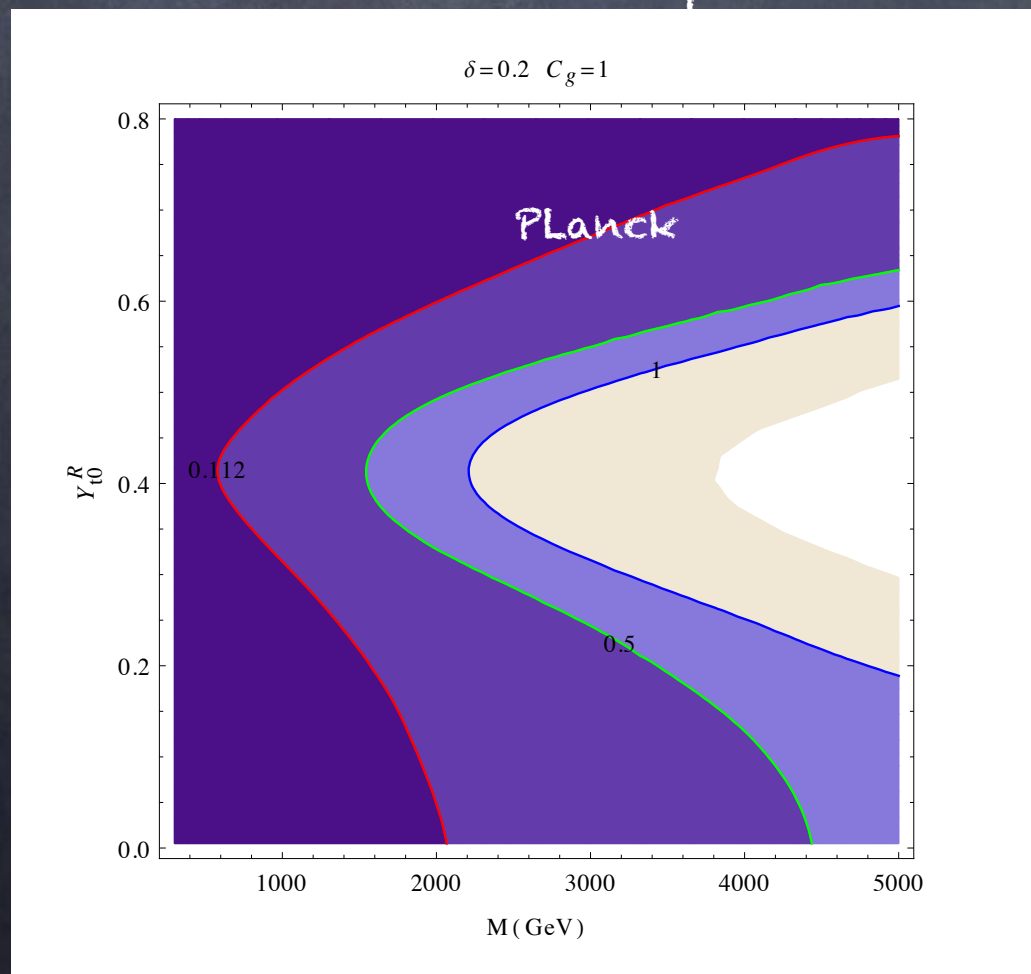
$$\Pi = \frac{1}{2} \begin{pmatrix} \sigma_i \Delta^i + s/\sqrt{2} & -i\Phi_H \\ i\Phi_H^\dagger & \sigma_i N^i - s/\sqrt{2} \end{pmatrix}$$

Triplet Complex bi-doublet (2HDM)
SU(2)_R Triplet

A composite 2HDM: Dark-Matter

- Singlet pheno similar to $SU(4)/Sp(4)$ case
- Other pNGBs can be stable:

DM = $SU(2)_R$ triplet



Preliminary results
(G.C., T.Ma, S.Zhao, B.Zhang)

Good DM thermal
relic densities possible.

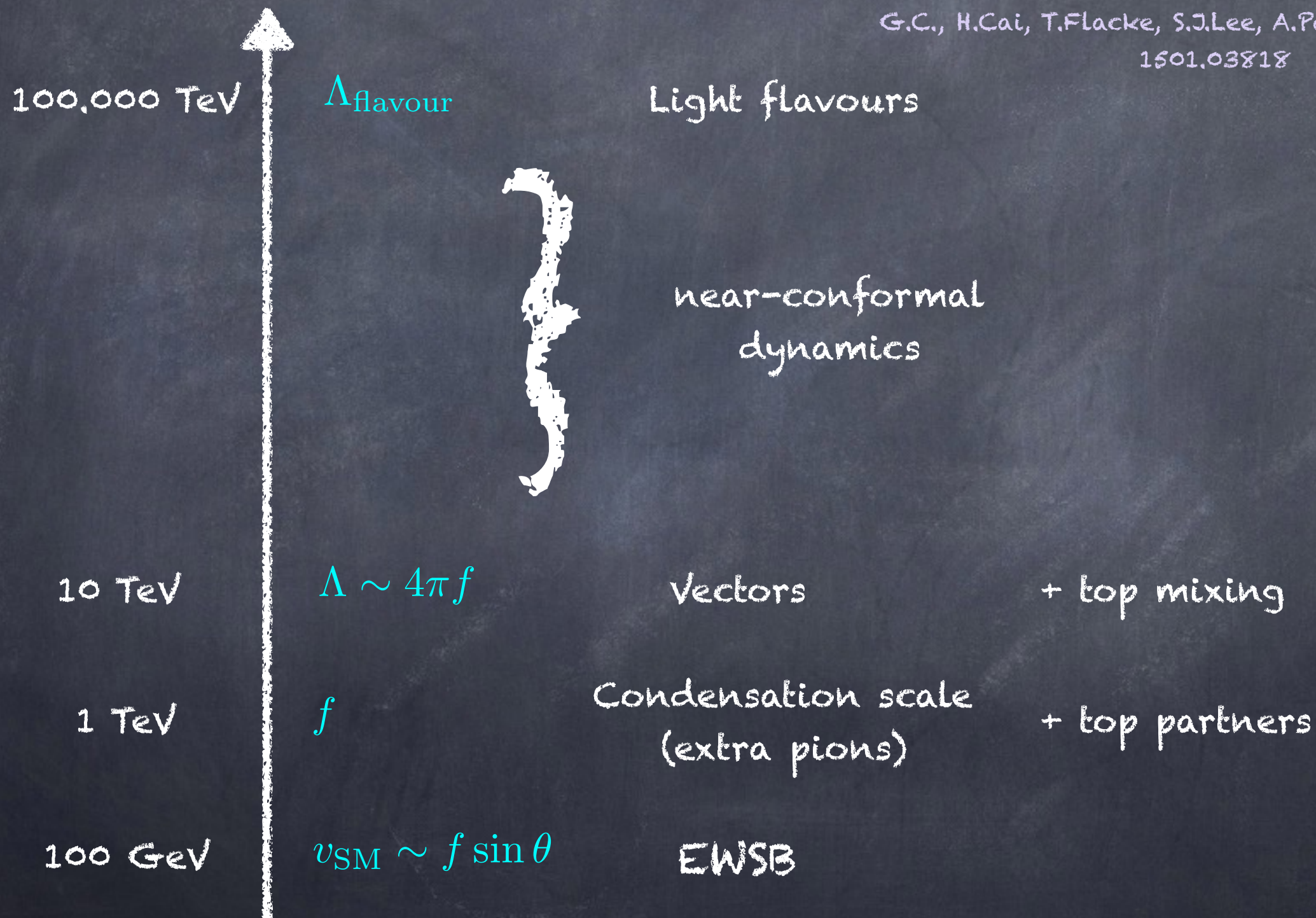
$$m_{DM} \sim 1 \div 2 \text{ TeV}$$

$$\sin \theta \sim 0.1$$

$SU(6)/Sp(6)$ and others
under investigation!

The hot potato: flavour!

G.C., H.Cai, T.Flacke, S.J.Lee, A.Parolini, H.Serodio
1501.03818



The hot potato: flavour!

- In FCD, the origin of the top partial compositeness can be probed!

$$\frac{c}{\Lambda_{\text{top}}^2} t\psi\psi \rightarrow y_t tT$$

- Why are the top partners lighter than other resonances? Conformal dynamics, or symmetry reason?
- "Theory of flavour" can be addressed!

The hot potato: flavour!

Typical scenario:

G_{TC} : rep R rep R'

Q

χ

G.Ferretti, D.Karateev
1312.5330, 1604.06467

SM :

EW

colour + hypercharge

global : $\langle QQ \rangle \neq 0$



pNGB Higgs

DM?

a) $\langle \chi\chi \rangle \neq 0$

coloured pNGBs
di-boson

b) $\langle \chi\chi \rangle = 0$

Exception: 1506.00623

Light top partners

Light t Hooft Top Partners

G.C., A. Parolini
1511.05163

	gauged	global symmetries		
	$Sp(2N_c)$	$SU(N_Q)$	$SU(N_\chi)$	$U(1)$
Q	\square	N_Q	1	1
χ	\square	1	N_χ	$-\frac{N_Q}{2N_\chi(N_c-1)}$

TABLE I: Fermionic field content of the first model.

Barnard, Gherghetta, Ray 1311.6562
Ferretti, Karateev 1312.5330

- $\langle QQ \rangle \neq 0$ breaks $SU(N_Q) \rightarrow Sp(N_Q)$, giving a pNGB Higgs.
- In the vacuum $\langle \chi\chi \rangle = 0$, $SU(N_\chi)$ is unbroken.

$$A_{SU(N_\chi)^2} \propto \dim(\chi) = (2N_c + 1)(N_c - 1)$$

Light t Hooft Top Partners

G.C., A. Parolini
1511.05163

global symmetries

	$SU(N_Q) \times SU(N_\chi)$	$Sp(N_Q) \times SU(N_\chi)$	$d_{Sp(N_Q)}$
χQQ	(\mathbf{A}, \mathbf{F})	$(\mathbf{1}, \mathbf{F})$	1
$\chi \bar{Q} \bar{Q}$	(\mathbf{S}, \mathbf{F})	(\mathbf{A}, \mathbf{F})	$\frac{N_Q(N_Q-1)}{2} - 1$
$\bar{\chi} \bar{Q} Q$	$(\mathbf{1}, \bar{\mathbf{F}})$	$(\mathbf{S}, \bar{\mathbf{F}})$	$\frac{N_Q(N_Q+1)}{2}$
	$(\mathbf{Adj}, \bar{\mathbf{F}})$	$(\mathbf{A}, \bar{\mathbf{F}})$	1
		$(\mathbf{S}, \bar{\mathbf{F}})$	$\frac{N_Q(N_Q-1)}{2} - 1$
			$\frac{N_Q(N_Q+1)}{2}$

$$A_{SU(N_\chi)^2} \propto \dim(\chi) = (2N_c + 1)(N_c - 1)$$

- Matched if $N_Q = 2N_c$. For $N_Q = 4$, $A = \mathfrak{so}(5)$ of $Sp(4) \sim SO(5)$!!

Predicting di-boson resonances

More precisely, the global symmetries are:

$$SU(N_Q) \times SU(N_X) \times U(1)_Q \times U(1)_X$$

$$\mathcal{L} \supset \frac{g_i^2}{32\pi^2} \frac{\kappa_i}{f_a} \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu}^i G_{\alpha\beta}^i,$$

$$\frac{\kappa_i}{f_a} = \frac{q_\psi \kappa_i^\psi + q_X \kappa_i^X}{\sqrt{q_\psi^2 f_\psi^2 + q_X^2 f_X^2}},$$

$$G = A, W, Z, g \quad !!!$$

Cai, Flacke, Lespinasse 1512.04508

Anomalous U(1) \rightarrow heavy η'

Orthogonal U(1) \rightarrow pNGB a

Decays and production only via WZW anomaly.

Predicting di-boson resonances

G_{HC}	ψ	χ	EW	Colour	X
$Sp(2N_c), 2 \leq N_c \leq 18$	F	A	$\frac{SU(4)}{Sp(4)}$	$\frac{SU(6)}{SO(6)}$	2/3
SO(11), SO(13)	Spin	F			2/3
$Sp(2N_c), N_c \geq 2$	A	F			1/3
$Sp(2N_c), N_c \geq 6$	Adj	F	$\frac{SU(5)}{SO(5)}$	$\frac{SU(6)}{Sp(6)}$	1/3
SO(11), SO(13)	F	Spin			1/3
SO(7), SO(9)	Spin	F			2/3
SO(7), SO(9)	F	Spin	$\frac{SU(5)}{SO(5)}$	$\frac{SU(6)}{SO(6)}$	1/3
$SO(N_c), N_c \geq 15$	Adj	F			1/3
$SO(N_c), N_c \geq 55$	S	F			1/3
SU(4)	A	F	$\frac{SU(5)}{SO(5)}$	$\frac{SU(3)^2}{SU(3)}$	1/3
SO(10), SO(14)	F	Spin			1/3
SU(4)	F	A	$\frac{SU(4)^2}{SU(4)}$	$\frac{SU(6)}{SO(6)}$	2/3
SO(10)	Spin	F			2/3
SU(7)	F	A₃			1/12
$SU(N_c), N_c \geq 5$	F	A			2/3
$SU(N_c), N_c \geq 5$	F	S	$\frac{SU(4)^2}{SU(4)}$	$\frac{SU(3)^2}{SU(3)}$	2/3
$SU(N_c), N_c \geq 5$	A	F			1/12
$SU(N_c), N_c \geq 8$	S	F			1/12

TABLE I: The complete list of theories. The HC representations are: **F** fundamental, **S** 2-index symmetric, **A** 2-index anti-symmetric, **A₃** 3 index anti-symmetric, **Adj** adjoint, **Spin** spinorial of SO. The last column contains the $U(1)_X$ charge assignment.

Model scan for the diphoton excess

Belyaev, G.C. Cai, Flacke, Parolini, Serodio 1512.07242

$$R_{XY} = \frac{\text{BR}(\sigma \rightarrow XY)}{\text{BR}(\sigma \rightarrow \gamma\gamma)}$$

$$R_{gg} \lesssim 1400, R_{WW} \lesssim 19, R_{ZZ} \lesssim 6, R_{Z\gamma} \lesssim 2,$$

$SU(4)^2/SU(4)$

$SU(4)/Sp(4)$

$SU(5)/SO(5)$



		R_{WW}	R_{ZZ}	$R_{Z\gamma}$	R_{gg}	Γ_{tot}	f_a
SU(7)	(F, A ₃)	9.5	3.0	0.8	140	0.4	2900
SU(5)	(A, F)	10	3.2	0.91	1300	3.2	830
SO(11)	(Spin, F)	4.4	0.51	3.5	500	0.8	2330
SO(13)	(Spin, F)	2.6	0.2	2.6	400	1.0	4000
SU(4)	(A, F)	23	6.6	3.4	960	1.7	680
SO(7)	(F, Spin)	20	5.7	2.7	600	1.5	1300
SO(9)	(F, Spin)	16	4.8	2.0	300	0.8	2200
SO(10)	(F, Spin)	15	4.6	1.8	227	0.6	2500
SO(11)	(F, Spin)	15	4.3	1.7	180	0.4	2900
SO(13)	(F, Spin)	13	4.1	1.5	120	0.3	3500
SO(14)	(F, Spin)	13	4.0	1.4	99	0.2	3800



if $f_a \sim f$
 $\sin \theta \sim 0.1$

TABLE II: List of models that can explain the di-photon excess and are compatible with present data. The models are grouped according to the Higgs coset: $SU(4)^2/SU(4)$ for the top block, $SU(4)/Sp(4)$ for the second block, and $SU(5)/SO(5)$ for the bottom one. Values for Γ_{tot} and f_a are given in GeV.

Summary and outlook

- FCD is a guide to build composite Higgs models!
- A very simple Higgs sector is possible ($SU(4)/Sp(4)$)
- Less minimal models feasible: DM? LHC signatures?
- FCD allows to probe fermion mass generations (UV completion, conformal windows...)
- FCD models can be tested on the Lattice!
- Di-boson (photon) @ the LHC is a "standard candle" for compositeness (via the WZW anomaly)

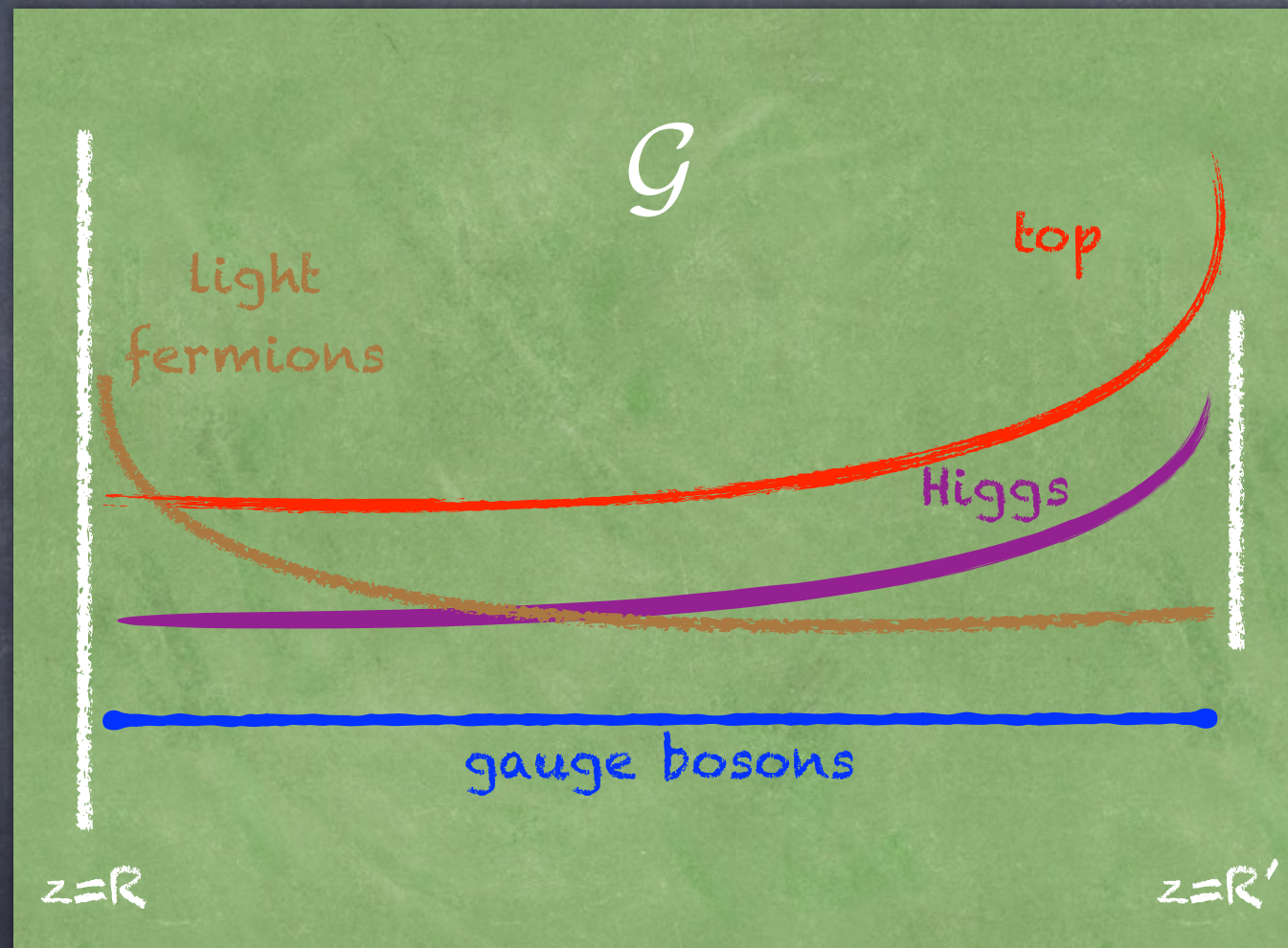
Backup

Rediscovery of composite Higgses

$$ds^2 = \left(\frac{R}{z}\right)^2 (dx^2 - dz^2)$$

Randall Sundrum

Elementary
sector

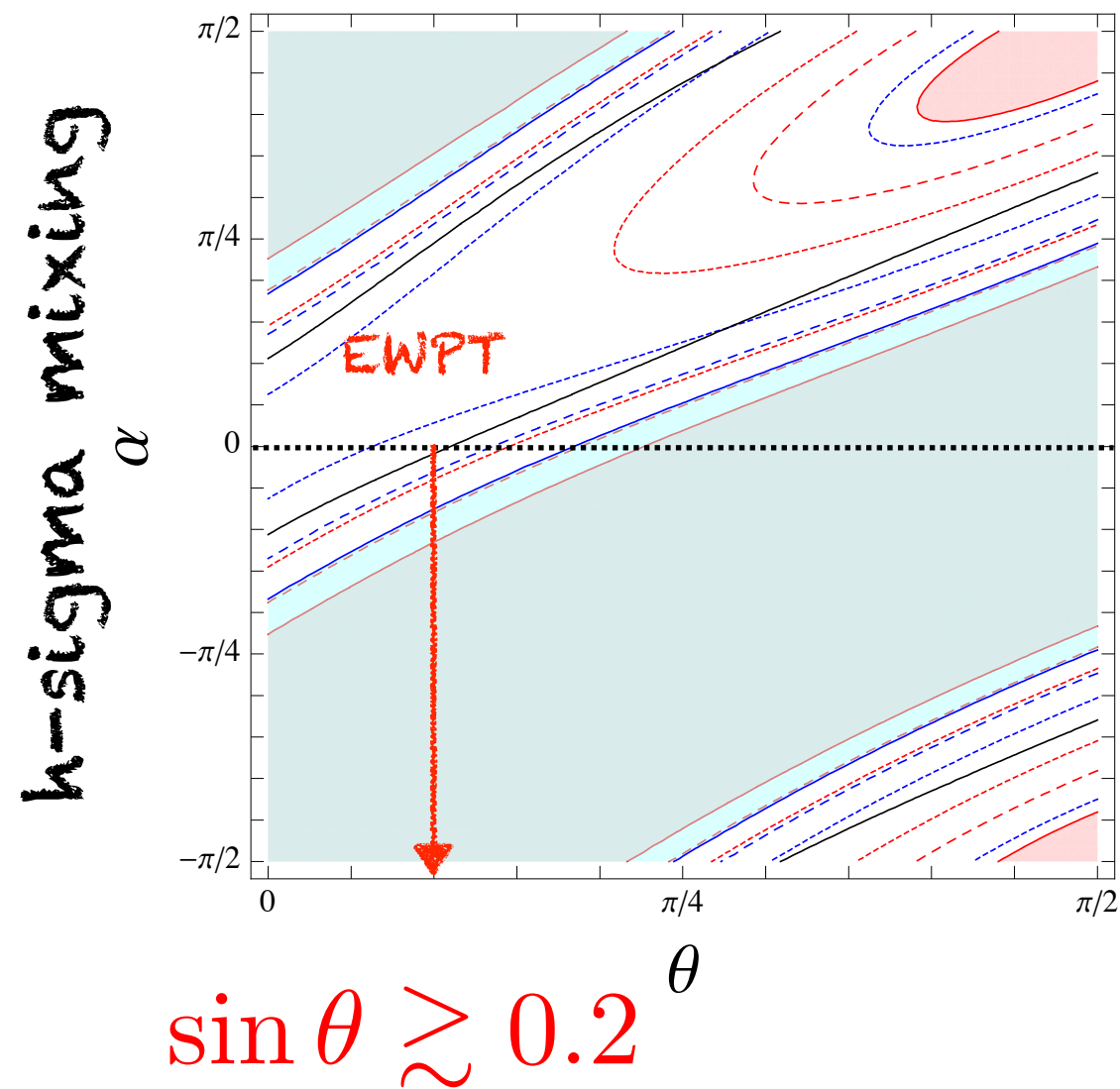


Breaking to H

- Conformal dynamics
- UV-insensitive potential
- Fermion partial compositeness
- Flavour under control

The minimal case

Arbey, G.C., Cai, Deandrea, Le Corre, Sannino
1502.04718



A composite 2HDM

G.C., T.Ma
1508.07014

$$\Pi = \frac{1}{2} \begin{pmatrix} \sigma_i \Delta^i + s/\sqrt{2} & -i\Phi_H \\ i\Phi_H^\dagger & \sigma_i N^i - s/\sqrt{2} \end{pmatrix}$$

$$\langle \Phi_H \rangle = \langle H_1 + iH_2 \rangle = \begin{pmatrix} ve^{i\beta} & 0 \\ 0 & ve^{i\beta} \end{pmatrix}$$

$$\Sigma = \begin{pmatrix} \cos \theta & 1 & e^{i\beta} \sin \theta & 1 \\ -e^{i\beta} \sin \theta & 1 & \cos \theta & 1 \end{pmatrix}$$

Beta can be removed by
an SU(4) rotation:

$$\Omega_\beta = \text{Exp} \left[-i\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right] = \begin{pmatrix} e^{-i\beta/2} & 0 \\ 0 & e^{i\beta/2} \end{pmatrix}$$

Beta = relative phase of the two T-quarks!

A composite 2HDM

G.C., T.Ma
1508.07014

$$\mathcal{L}_{\text{Yuk}} = -f (\bar{q}_L^\alpha t_R) \left[\text{Tr}[P_{1,\alpha}(y_{t1}\Sigma + y_{t2}\Sigma^\dagger)] + (i\sigma_2)_{\alpha\beta} \text{Tr}[P_2^\beta(y_{t3}\Sigma + y_{t4}\Sigma^\dagger)] \right] + h.c.$$

4 "Yukawa" couplings!

$$Y_t = \frac{y_{t1} - y_{t2} - (y_{t3} - y_{t4})}{2\sqrt{2}}, \quad Y_D = \frac{y_{t1} - y_{t2} + (y_{t3} - y_{t4})}{2\sqrt{2}},$$

$$Y_T = \frac{y_{t1} + y_{t2} + (y_{t3} + y_{t4})}{2\sqrt{2}}, \quad Y_0 = \frac{y_{t1} + y_{t2} - (y_{t3} + y_{t4})}{2\sqrt{2}}.$$

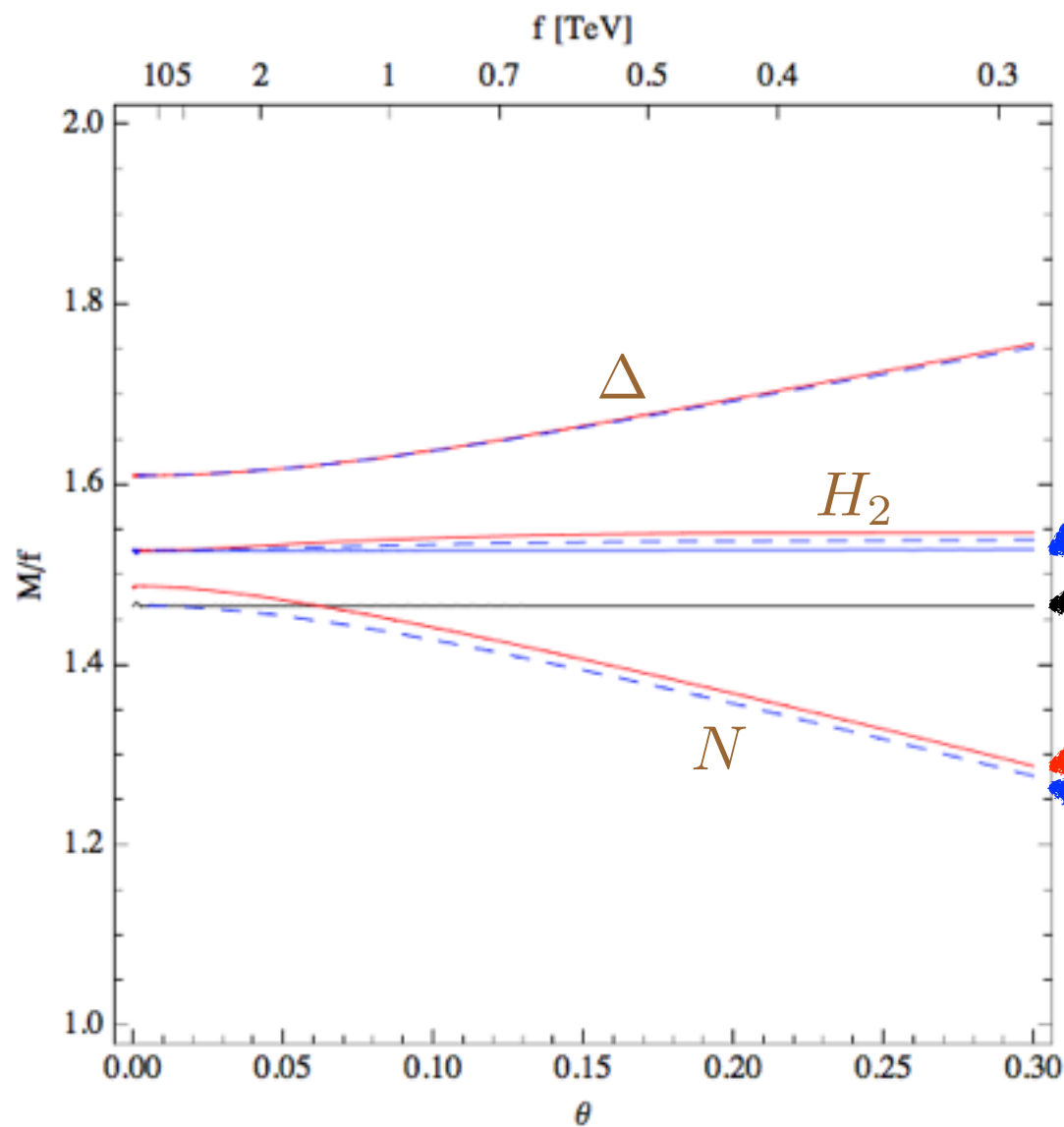
$$V_{\text{top}}(\theta) = -C_t f^4 \left[8|Y_t|^2 \sin^2 \theta + \leftarrow \text{Potential for theta} \right. \\ \left. 2\sqrt{2}|Y_t|^2 \sin(2\theta) \frac{h_1}{f} + \right.$$

$$\begin{array}{l} \text{Set to zero} \\ \text{by phase-shift} \end{array} \rightarrow +4\sqrt{2} \text{Im}(Y_D^* Y_t) \sin \theta \frac{h_2}{f}$$

$$\begin{array}{l} \text{Custodial} \\ \text{violating} \\ \text{VEVs!!!} \end{array} \rightarrow +2\sqrt{2} \text{Re}(Y_D^* Y_t) \sin(2\theta) \frac{A_0}{f}$$

$$\rightarrow +4 \text{Im}(Y_T^* Y_t) \sin^2 \theta \frac{N_0 + \Delta_0}{f} + \dots \left. \right]$$

A composite 2HDM: spectrum



$$F_\pi = 2\sqrt{2}f$$

CP-odd

singlet

charged

neutral
(DM?)

A composite 2HDM: EWPTs

$$\Delta S_{\text{Higgs}} = \frac{1 - \kappa_V^2}{6\pi} \ln \frac{\Lambda_{FCD}}{m_h}, \quad \Delta T_{\text{Higgs}} = -\frac{3(1 - \kappa_V^2)}{8\pi \cos^2 \theta_W} \ln \frac{\Lambda_{FCD}}{m_h},$$

$$\Delta S_{p\text{NGB}} = -\frac{\sin^2 \theta}{4\pi}, \quad \Delta T_{p\text{NGB}} = \frac{\sin^2 \theta}{8\pi \sin^2 \theta_W} \frac{m_{H^\pm}^2 - m_{A_0}^2}{m_W^2} \ln \frac{\Lambda_{FCD}}{m_{p\text{NGB}}} \sim 0,$$

$$\Delta S_{FCD} = \frac{\sin^2 \theta}{3\pi} N, \quad \Delta T_{FCD} \sim 0,$$

Bounds similar
to minimal cases.

