

Determining neutrino oscillation parameters in presence of Non-Standard neutrino Interaction (NSI)

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Reference

- P. B. and Y. Farzan, “CP-Violation and Non-Standard Interactions at the MOMENT,” JHEP **1607** (2016) 109 [arXiv:1602.07099 [hep-ph]].
- P. B. and A. N. Khan, “Sensitivities to charged-current nonstandard neutrino interactions at DUNE,” arXiv:1607.00065 [hep-ph].

Overview

- 1 Introduction
- 2 T2K, NO ν A, MOMENT, DUNE
- 3 Results
- 4 Summary and Conclusion

Introduction

Neutrino Oscillation

There is a mixing between mass and flavor states

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle, \quad \alpha = e, \mu, \tau, \quad i = 1, 2, 3 \quad (1)$$

PMNS mixing matrix

$$U = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \quad (2)$$

Neutrino Oscillation in Vacuum

$$H |\nu_k\rangle = E_k |\nu_k\rangle \quad (3)$$

$$i \frac{d}{dt} |\nu_k(t)\rangle = H |\nu_k(t)\rangle \quad (4)$$

$$|\nu_k(t)\rangle = e^{-iE_k t} |\nu_k\rangle \quad (5)$$

$$|\nu_\alpha(t)\rangle = \sum_k U_{\alpha k}^* e^{-iE_k t} |\nu_k\rangle \quad (6)$$

$$|\nu_\alpha(t)\rangle = \sum_\beta \sum_k U_{\alpha k}^* e^{-iE_k t} U_{\beta k} |\nu_\beta\rangle \quad (7)$$

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* e^{-i(E_k - E_j)t} \quad (8)$$

Neutrino Oscillation in Vacuum

$$E_k \simeq E + \frac{m_k^2}{2E} \quad (9)$$

$$E_k - E_j \simeq \frac{\Delta m_{kj}^2}{2E} \equiv \frac{m_k^2 - m_j^2}{2E} \quad (10)$$

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* e^{-i \frac{\Delta m_{kj}^2 L}{2E}} \quad (11)$$

Neutrino Oscillation parameters

NuFIT 2.2 (2016)

| | Normal Ordering (best fit) | | Inverted Ordering ($\Delta\chi^2 = 0.56$) | | Any Ordering |
|---|---------------------------------|-------------------------------|---|-------------------------------|--|
| | bfp $\pm 1\sigma$ | 3σ range | bfp $\pm 1\sigma$ | 3σ range | 3σ range |
| $\sin^2 \theta_{12}$ | $0.308^{+0.013}_{-0.012}$ | $0.273 \rightarrow 0.348$ | $0.308^{+0.013}_{-0.012}$ | $0.273 \rightarrow 0.349$ | $0.273 \rightarrow 0.348$ |
| $\theta_{12}/^\circ$ | $33.72^{+0.79}_{-0.76}$ | $31.52 \rightarrow 36.18$ | $33.72^{+0.79}_{-0.76}$ | $31.52 \rightarrow 36.18$ | $31.52 \rightarrow 36.18$ |
| $\sin^2 \theta_{23}$ | $0.440^{+0.023}_{-0.019}$ | $0.388 \rightarrow 0.630$ | $0.584^{+0.018}_{-0.022}$ | $0.398 \rightarrow 0.634$ | $0.388 \rightarrow 0.632$ |
| $\theta_{23}/^\circ$ | $41.5^{+1.3}_{-1.1}$ | $38.6 \rightarrow 52.5$ | $49.9^{+1.1}_{-1.3}$ | $39.1 \rightarrow 52.8$ | $38.6 \rightarrow 52.7$ |
| $\sin^2 \theta_{13}$ | $0.02163^{+0.00074}_{-0.00074}$ | $0.01938 \rightarrow 0.02388$ | $0.02175^{+0.00075}_{-0.00074}$ | $0.01950 \rightarrow 0.02403$ | $0.01938 \rightarrow 0.02396$ |
| $\theta_{13}/^\circ$ | $8.46^{+0.14}_{-0.15}$ | $8.00 \rightarrow 8.89$ | $8.48^{+0.15}_{-0.15}$ | $8.03 \rightarrow 8.92$ | $8.00 \rightarrow 8.90$ |
| $\delta_{\text{CP}}/^\circ$ | 289^{+38}_{-51} | $0 \rightarrow 360$ | 269^{+39}_{-45} | $146 \rightarrow 377$ | $0 \rightarrow 360$ |
| $\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$ | $7.49^{+0.19}_{-0.17}$ | $7.02 \rightarrow 8.08$ | $7.49^{+0.19}_{-0.17}$ | $7.02 \rightarrow 8.08$ | $7.02 \rightarrow 8.08$ |
| $\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$ | $+2.526^{+0.039}_{-0.037}$ | $+2.413 \rightarrow +2.645$ | $-2.518^{+0.038}_{-0.037}$ | $-2.634 \rightarrow -2.406$ | $\begin{bmatrix} +2.413 \rightarrow +2.645 \\ -2.630 \rightarrow -2.409 \end{bmatrix}$ |

Neutrino Oscillation in Matter

- Forward elastic scattering processes affect neutrino oscillation

$$H = H_{vac} + H_{mat} \quad (12)$$

$$H_{vac} = U_{PMNS} \cdot \text{Diag}(0, \Delta m_{21}^2, \Delta m_{31}^2) \cdot U_{PMNS}^\dagger \quad (13)$$

$$H_{mat} = \sqrt{2} G_F N_e \text{diag}(1, 0, 0) \quad (14)$$

- Assuming neutral-current non-standard neutrino interaction

$$H_{mat} = \sqrt{2} G_F N_e \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \quad (15)$$

- NC NSI

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \varepsilon_{\alpha\beta}^{fC} (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{f} \gamma_\mu P_C f) , \quad (19)$$

$$\epsilon_{\alpha\beta} = (N_d/N_e)\epsilon_{\alpha\beta}^d + (N_u/N_e)\epsilon_{\alpha\beta}^u$$

- CC NSI (pion decay)

$$\mathcal{L}^S = -2\sqrt{2}G_F \sum_{\beta} (\delta_{\mu\beta} + \varepsilon_{\alpha\beta}^{udL}) (\bar{\mu} \gamma_\lambda P_L U_{\beta a} \nu_a) (\bar{d} \gamma^\lambda P_L u)^\dagger \quad (20)$$

- CC NSI (muon decay)

$$\mathcal{L}^L = -2\sqrt{2}G_F \sum_{\alpha,\beta} (\delta_{\alpha\beta} + \varepsilon_{\alpha\beta}^L) (\bar{e} \gamma_\lambda P_L U_{\alpha a} \nu_a) (\bar{\mu} \gamma^\lambda P_L U_{\beta b} \nu_b)^\dagger \quad (21)$$

- Current constraints on NC NSI parameters (arXiv:1307.3092)

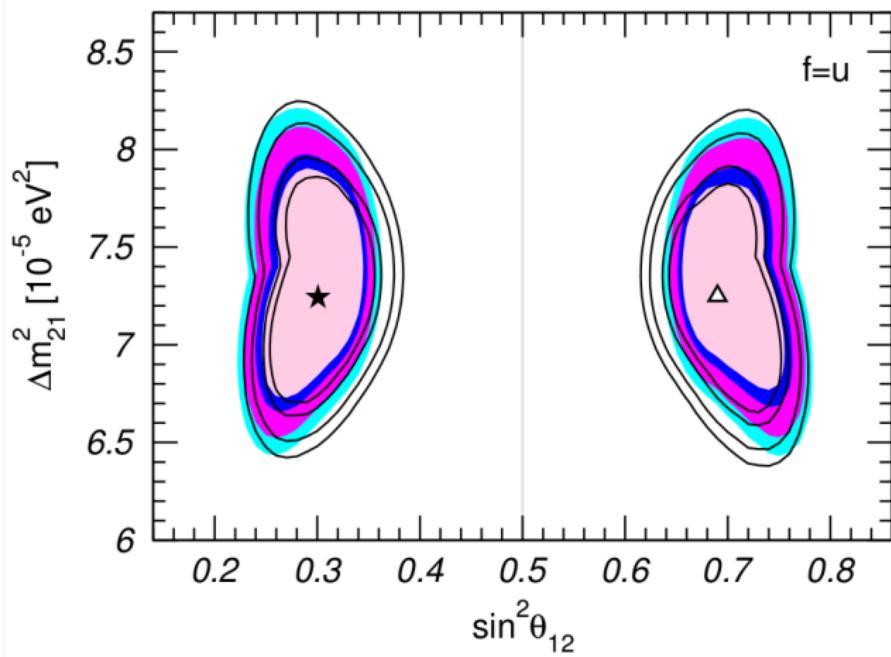
$$\begin{aligned} |\epsilon_{e\mu}| &< 0.16 \\ |\epsilon_{e\tau}| &< 0.26 \\ |\epsilon_{\mu\tau}| &< 0.02 \end{aligned} \tag{22}$$

and

$$\begin{aligned} -0.018 &< \epsilon_{\tau\tau} - \epsilon_{\mu\mu} &< 0.054 \\ 0.35 &< \epsilon_{ee} - \epsilon_{\mu\mu} &< 0.93. \end{aligned} \tag{23}$$

- NC NSI affect neutrino oscillation during propagation in matter
- $a_{CC} = 2\sqrt{2}G_F N_e E$
- CC NSI affect flavour change in neutrino production and detection

LMA-Dark



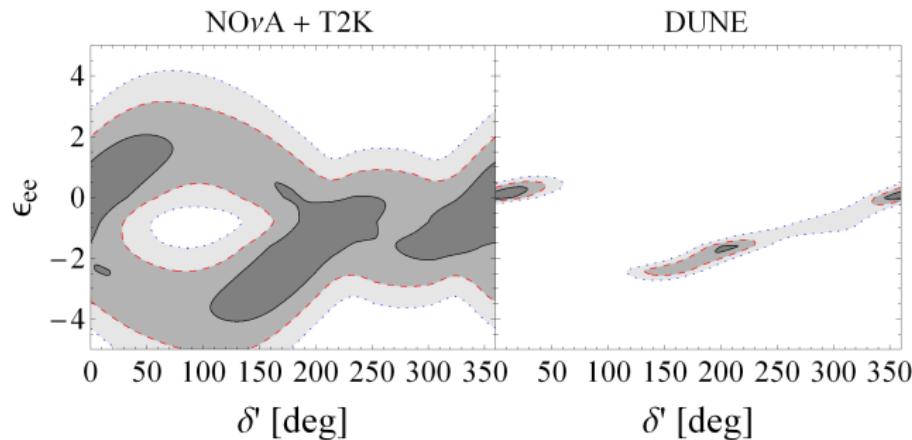
M. C. Gonzalez-Garcia and M. Maltoni, JHEP **1309** (2013) 152,
arXiv:1307.3092

Current Constraints on CC NSI

| parameter | Current constraints (90%C.L.) |
|---------------------------|-------------------------------|
| $ \epsilon_{\mu e}^s $ | 0.026 |
| $ \epsilon_{\mu \mu}^s $ | 0.078 |
| $ \epsilon_{\mu \tau}^s $ | 0.013 |
| $ \epsilon_{ee}^d $ | 0.041 |
| $ \epsilon_{e\mu}^d $ | 0.026 |
| $ \epsilon_{\mu e}^d $ | 0.025 |
| $ \epsilon_{\mu \mu}^d $ | 0.078 |
| $ \epsilon_{\tau e}^d $ | 0.041 |
| $ \epsilon_{\tau \mu}^d $ | 0.013 |

C. Biggio, M. Blennow and E. Fernandez-Martinez, JHEP **0908** (2009) 090 [arXiv:0907.0097 [hep-ph]].

Degeneracy between NSI and δ_{CP}



A. de Gouvea and K. J. Kelly, Nucl. Phys. B **908** (2016) 318
[arXiv:1511.05562 [hep-ph]].

- Muon disappearance probability in presence of NC NSI

$$\begin{aligned}
 P_{\nu_\mu \rightarrow \nu_\mu} = & 1 - \sin^2 2\theta_{23} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \\
 & - |\epsilon_{\mu\tau}| \cos \phi_{\mu\tau} \sin 2\theta_{23} \left[\sin^2 2\theta_{23} \frac{a_{CC} L}{2E} \sin \frac{\Delta m_{31}^2 L}{2E} \right. \\
 & + 4 \cos^2 2\theta_{23} \frac{a_{CC}}{\Delta m_{31}^2} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \left. \right] \\
 & + \frac{1}{2} (|\epsilon_{\mu\mu}| - |\epsilon_{\tau\tau}|) \sin^2 2\theta_{23} \cos 2\theta_{23} \left[\frac{a_{CC} L}{2E} \sin \frac{\Delta m_{31}^2 L}{2E} \right. \\
 & \left. - 4 \frac{a_{CC}}{\Delta m_{31}^2} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \right]
 \end{aligned} \tag{24}$$

$$\begin{aligned}
P_{\nu_\mu^s \rightarrow \nu_e^d}^{\text{mat}} &= 4\tilde{s}_{13} \sin^2 \theta_{23} \left(\tilde{s}_{13} + 2(|\epsilon_{e\mu}| \cos(\phi_{e\mu} + \delta_{\text{CP}}) + |\epsilon_{e\tau}| \cos(\phi_{e\tau} + \delta_{\text{CP}})) \sin \theta_{23} \frac{a_{\text{CC}}}{\Delta m_{31}^2 - a_{\text{CC}}} \right) \\
&\quad \sin^2 \frac{(\Delta m_{31}^2 - a_{\text{CC}})L}{4E} + \left(\frac{\Delta m_{21}^2}{a_{\text{CC}}} \right)^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2 \frac{a_{\text{CC}}L}{4E} \\
&\quad + 4(|\epsilon_{e\mu}| \cos \phi_{e\mu} - |\epsilon_{e\tau}| \cos \phi_{e\tau}) \frac{\Delta m_{21}^2}{a_{\text{CC}}} \sin 2\theta_{12} \cos^3 \theta_{23} \sin^2 \frac{a_{\text{CC}}L}{4E} \\
&\quad + \left(-\frac{\Delta m_{21}^2}{a_{\text{CC}}} \tilde{s}_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} - 4|\epsilon_{e\mu}| |\tilde{s}_{13} \sin \theta_{23} \cos^2 \theta_{23} \cos(\phi_{e\mu} + \delta_{\text{CP}}) \right. \\
&\quad \left. + 4|\epsilon_{e\tau}| |\tilde{s}_{13} \sin^2 \theta_{23} \cos \theta_{23} \cos(\phi_{e\tau} + \delta_{\text{CP}}) - 2|\epsilon_{e\mu}| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sin 2\theta_{12} \sin^2 \theta_{23} \cos \theta_{23} \cos \phi_{e\mu} \right. \\
&\quad \left. \frac{\Delta m_{31}^2}{\Delta m_{31}^2 - a_{\text{CC}}} - 2|\epsilon_{e\tau}| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sin 2\theta_{12} \sin \theta_{23} \cos^2 \theta_{23} \cos \phi_{e\tau} \frac{\Delta m_{31}^2}{\Delta m_{31}^2 - a_{\text{CC}}} \right) \left[\sin^2 \frac{a_{\text{CC}}L}{4E} \right. \\
&\quad \left. - \sin^2 \frac{\Delta m_{31}^2 L}{4E} + \sin^2 \frac{(\Delta m_{31}^2 - a_{\text{CC}})L}{4E} \right] + \left(-\frac{1}{2} \frac{\Delta m_{21}^2}{a_{\text{CC}}} \tilde{s}_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \right. \\
&\quad \left. - 2|\epsilon_{e\mu}| |\tilde{s}_{13} \sin \theta_{23} \cos^2 \theta_{23} \sin(\phi_{e\mu} + \delta_{\text{CP}}) + 2|\epsilon_{e\tau}| |\tilde{s}_{13} \sin^2 \theta_{23} \cos \theta_{23} \sin(\phi_{e\tau} + \delta_{\text{CP}}) \right. \\
&\quad \left. + |\epsilon_{e\mu}| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sin 2\theta_{12} \sin^2 \theta_{23} \cos \theta_{23} \sin \phi_{e\mu} \frac{\Delta m_{31}^2}{\Delta m_{31}^2 - a_{\text{CC}}} + |\epsilon_{e\tau}| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sin 2\theta_{12} \sin \theta_{23} \right. \\
&\quad \left. \cos^2 \theta_{23} \sin \phi_{e\tau} \frac{\Delta m_{31}^2}{\Delta m_{31}^2 - a_{\text{CC}}} \right) \left[\sin \frac{a_{\text{CC}}L}{2E} - \sin \frac{\Delta m_{31}^2 L}{2E} + \sin \frac{(\Delta m_{31}^2 - a_{\text{CC}})L}{2E} \right]
\end{aligned}$$

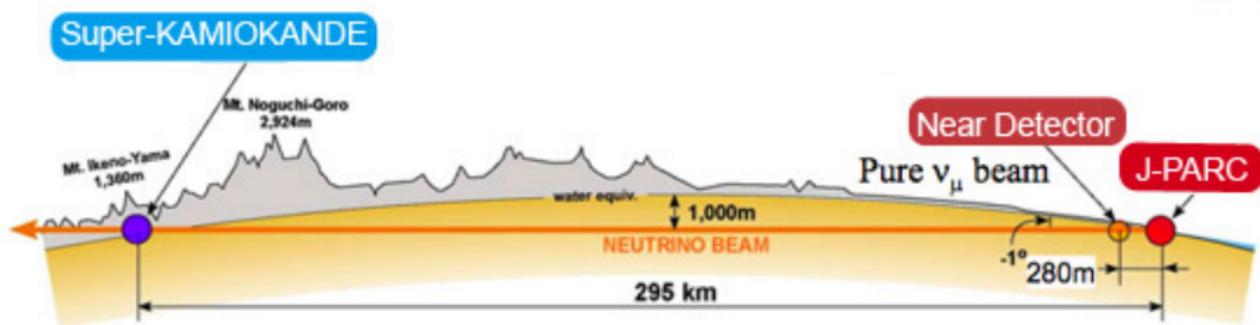
(25)

- If a_{CC} is negligible

$$\begin{aligned} P_{\nu_\mu \rightarrow \nu_e} = & 4 \sin^2 \theta_{13} \sin^2 \theta_{23} \sin^2 \frac{\Delta m_{31}^2 L}{4E} \\ & + \left(\frac{\Delta m_{21}^2 L}{4E} \right)^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \\ & + \frac{\Delta m_{21}^2 L}{4E} \sin \theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} (\cos \delta_{CP} \sin \frac{\Delta m_{31}^2 L}{2E} \\ & - 2 \sin \delta_{CP} \sin^2 \frac{\Delta m_{31}^2 L}{4E}). \end{aligned} \tag{26}$$

T2K, NO ν A, MOMENT, DUNE

T2K

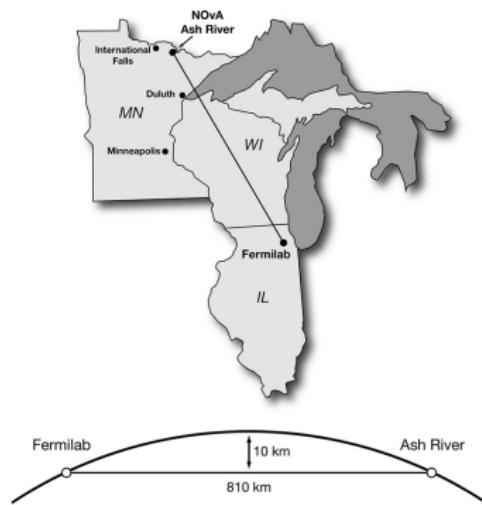


<http://www2.warwick.ac.uk/fac/sci/physics/research/epp/exp/t2k/>

T2K

- 2 years in neutrino and 6 years in anti-neutrino mode
- fiducial mass of SK: 22.5 kton
- Energy resolution is 85 MeV
- $400\text{MeV} < E < 1.2\text{GeV}$
- 2.5% normalization uncertainty for disappearance, 5% appearance
- free normalization for quasi-elastic events
- Background: lepton flavor misidentification, neutral-current events, charge misidentification and intrinsic background
- Background normalization uncertainty is 20%
- 5% uncertainties of matter density profile, (PREM)
- We used GLoBES for simulation

NO ν A

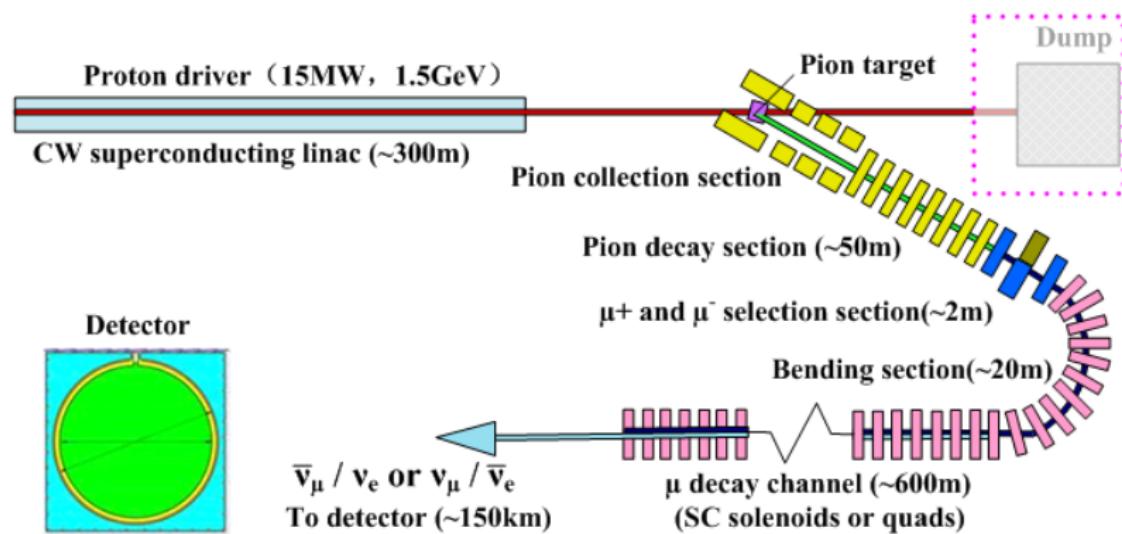


<http://www.universetoday.com/109444/nova-experiment-nabs-its-first-neutrinos/>

NO ν A

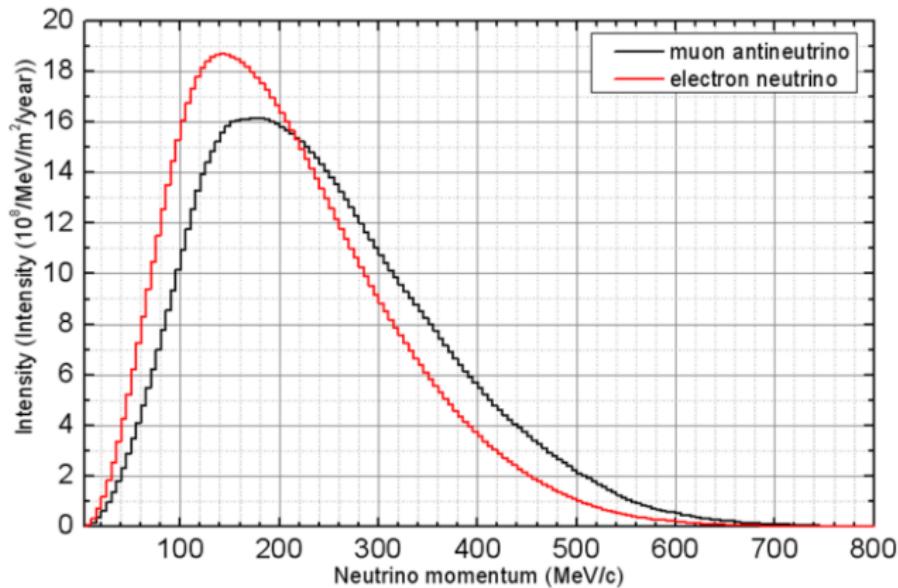
- 3 years in each mode
- 25 kton fiducial mass
- $1\text{GeV} < E < 3.5\text{GeV}$
- energy resolution is $10\%\sqrt{E}$ for electron neutrino and $5\%\sqrt{E}$ for muon neutrino
- normalization uncertainty is 5%
- Background: neutral current interaction, lepton flavor misidentification and intrinsic background

MOMENT



<https://indico.cern.ch/event/351600/contributions/1754021/>

MOMENT



<https://indico.cern.ch/event/351600/contributions/1754021/>

- MuOn decay MEdium baseline NeuTrino beam facility

MOMENT

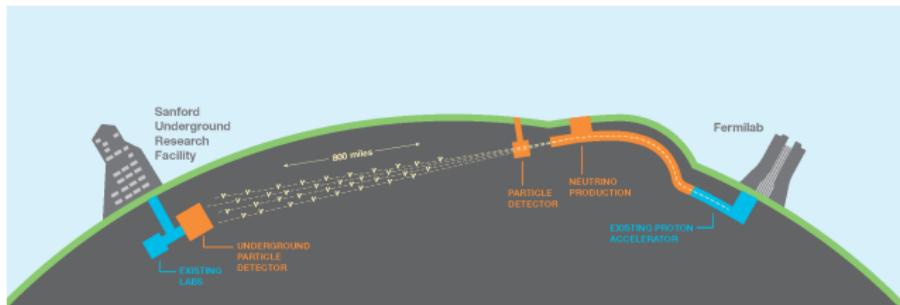
- baseline is 150 km
- 500 kton Gd-doped water Cherenkov detector
- muon mode, $\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$; and anti-muon mode, $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$
- $50\text{MeV} < E < 700\text{MeV}$ with peak around 150 MeV
- dominant interaction modes are quasi-elastic interactions:

$$\nu_e + n \rightarrow p + e^- \quad \bar{\nu}_\mu + p \rightarrow n + \mu^+$$

$$\bar{\nu}_e + p \rightarrow n + e^+ \quad \nu_\mu + n \rightarrow p + \mu^-$$

- Charge Identification is 80 %
- background: charge misidentification, atmospheric neutrinos, flavor misidentification
- neutrino flux of each flavor mode at the detector equal to $4.7 \times 10^{11} \text{ m}^{-2} \text{ year}^{-1}$
- 5 years of data taking in each mode

DUNE



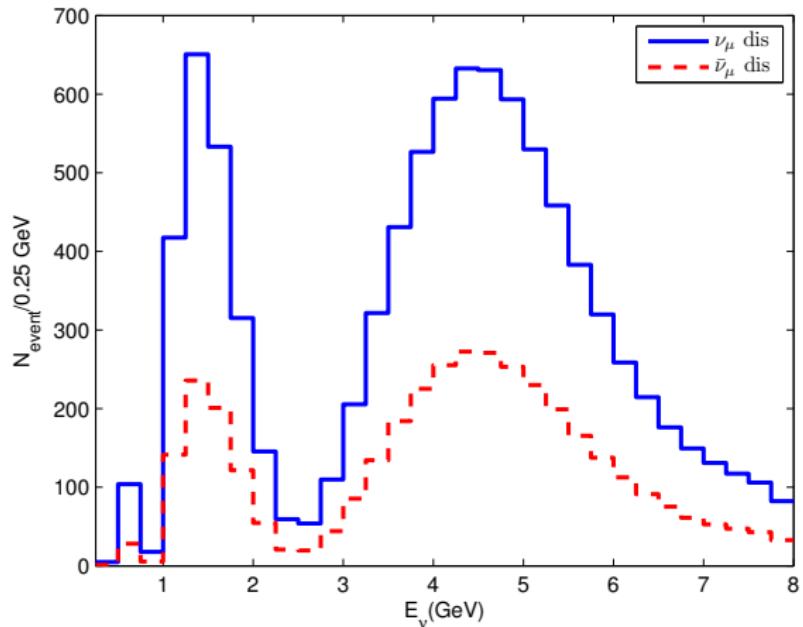
<http://lbnf.fnal.gov/>

- Deep Underground Neutrino Experiment (DUNE)
- 1300 km baseline
- $250\text{MeV} < E < 8\text{GeV}$ peak around 3 GeV
- 3 years of data taking in each mode

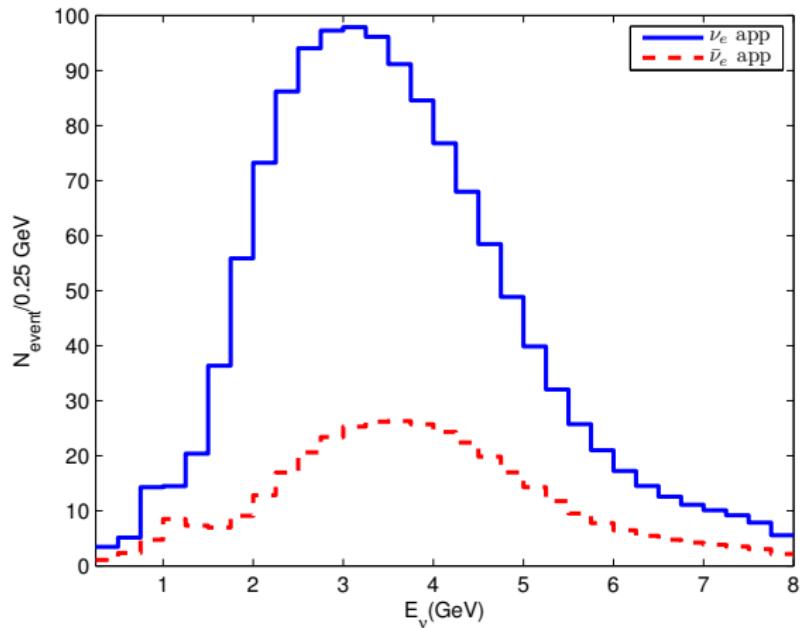
DUNE

- Liquid Argon Time-Projection Chamber(LArTPC) with 34 kton fiducial mass
- near detector: 5 ton at 460 m
- NC events, lepton flavor misidentification and intrinsic background
- $15\%/\sqrt{E(\text{GeV})}$ and $20\%/\sqrt{E(\text{GeV})}$ energy resolution for CC detection of electron neutrino and muon neutrino
- flux uncertainty is 5% and flux uncertainty of background is 10%

Number of events by FD



Number of events by FD

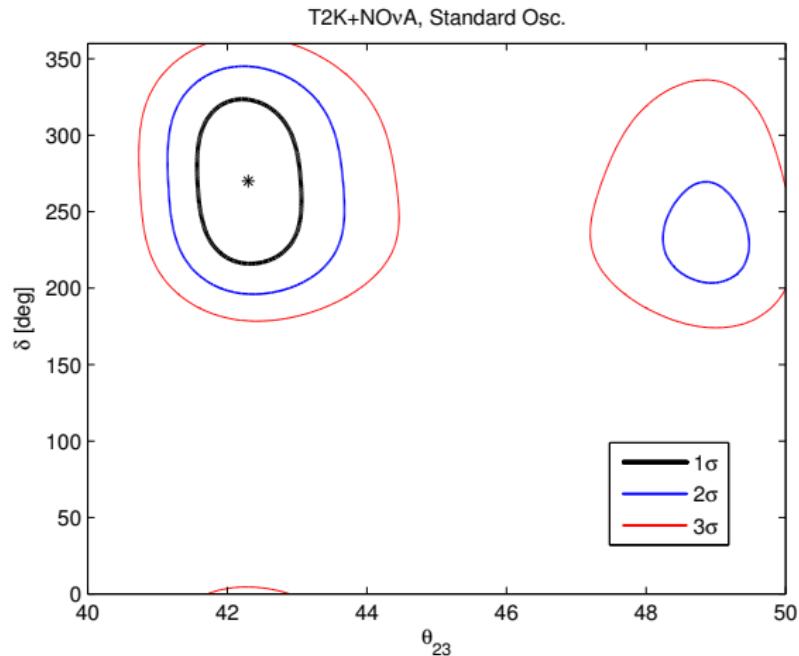


Results

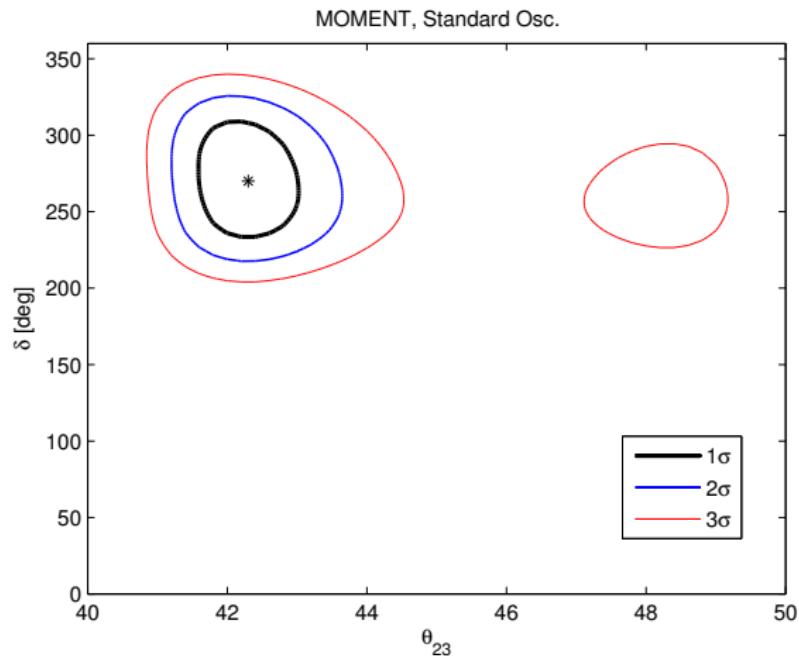
Details of our analysis

- at lower energies matter effect has less impact on neutrino oscillation ($a_{CC} = 2\sqrt{2}G_F N_e E$)
- Effect of NC NSI on determination of θ_{23} octant and δ_{CP} , for T2K, NO ν A and MOMENT
- Constraining NC NSI by T2K, NO ν A and MOMENT
- Effect of CC NSI on determining θ_{23} octant and δ_{CP} by DUNE
- Constraints on CC NSI by near and far detectors of DUNE

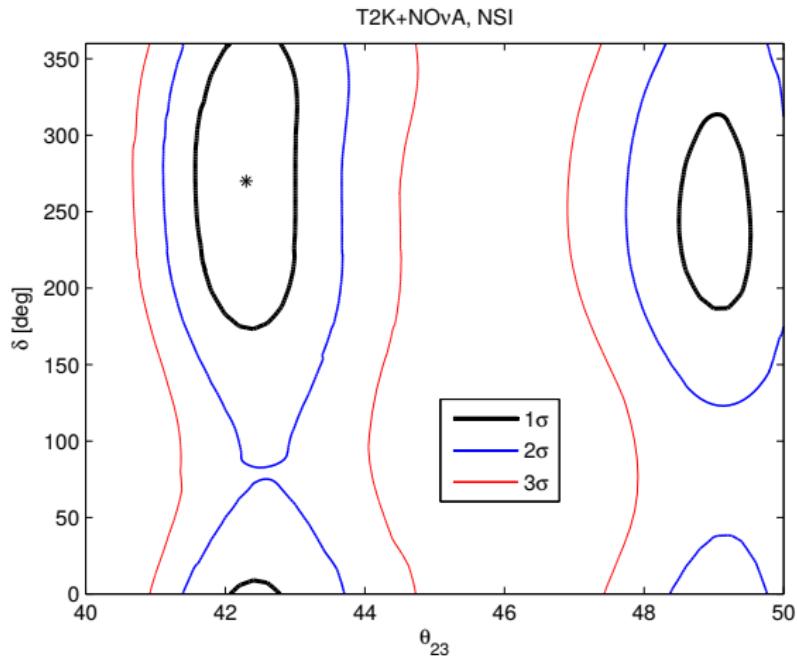
Determination of θ_{23} octant and δ_{CP}



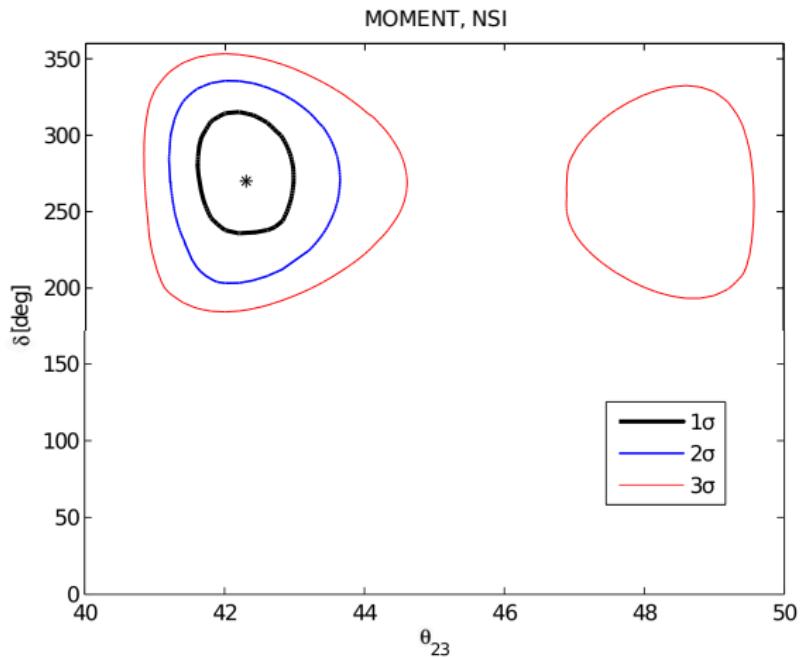
Determination of θ_{23} octant and δ_{CP}



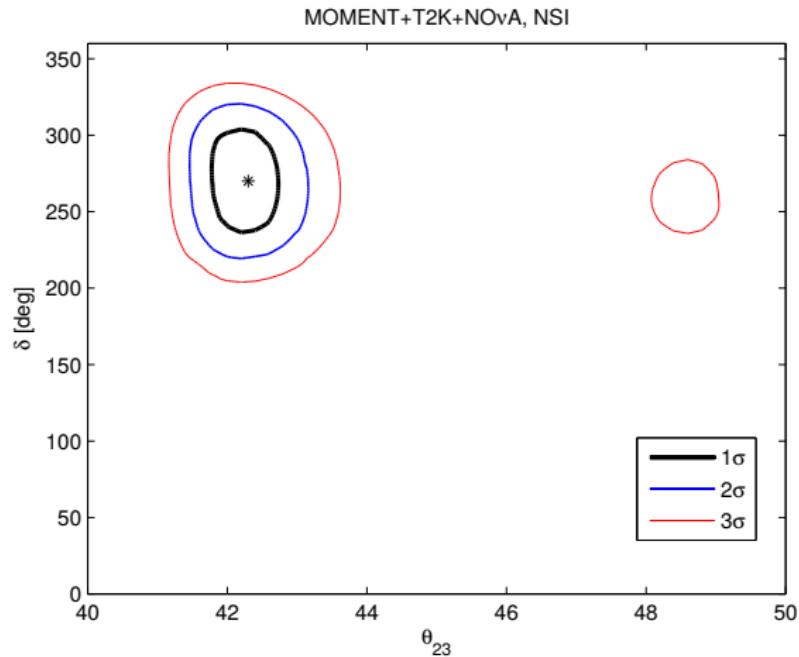
Determination of θ_{23} octant and δ_{CP}



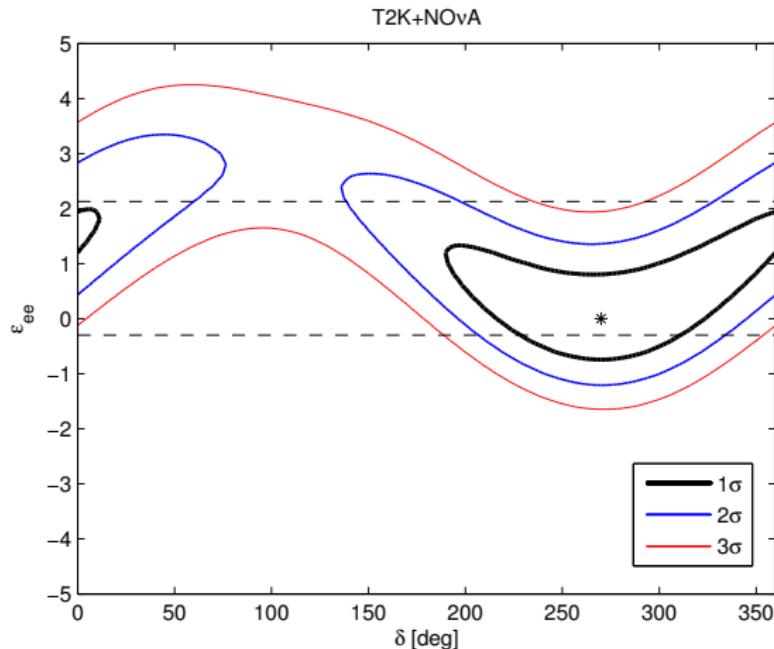
Determination of θ_{23} octant and δ_{CP}



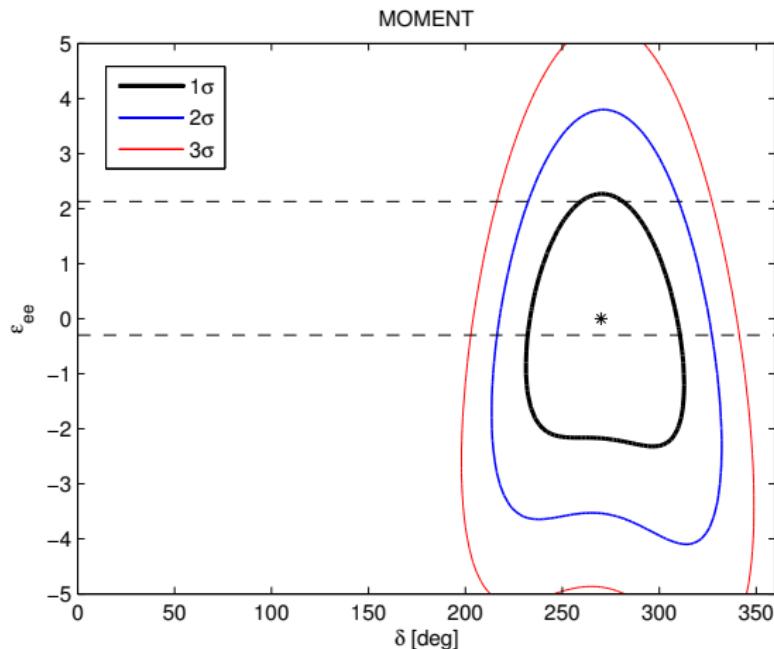
Determination of θ_{23} octant and δ_{CP}



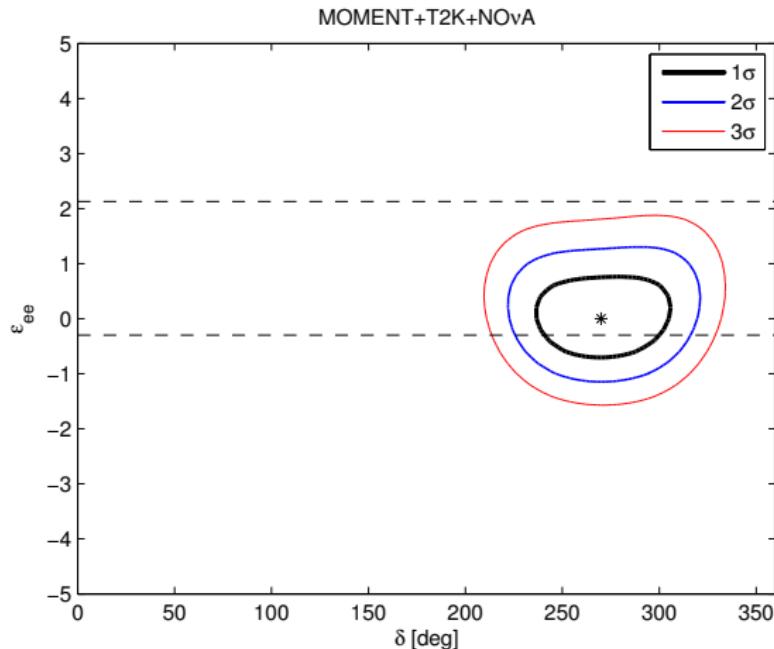
Sensitivity to ϵ_{ee} and δ_{CP}



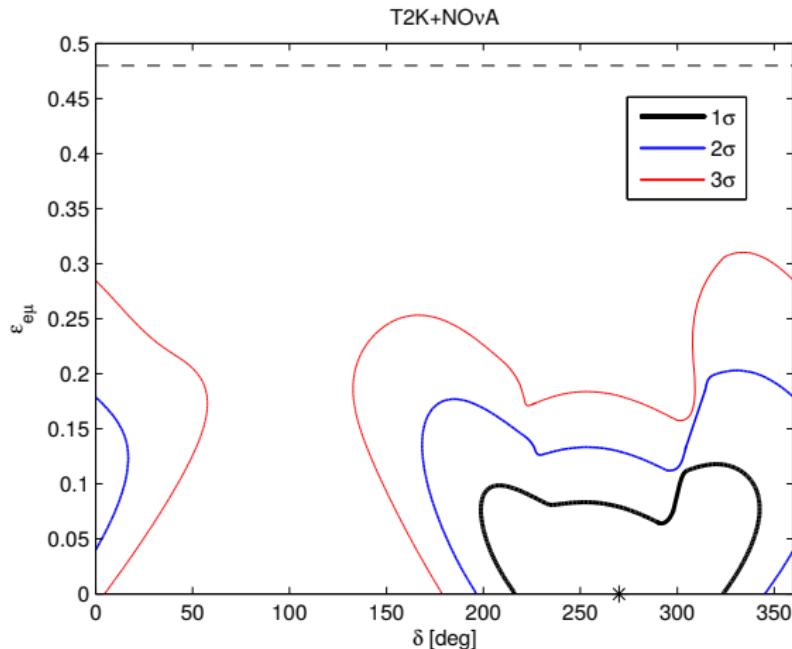
Sensitivity to ϵ_{ee} and δ_{CP}



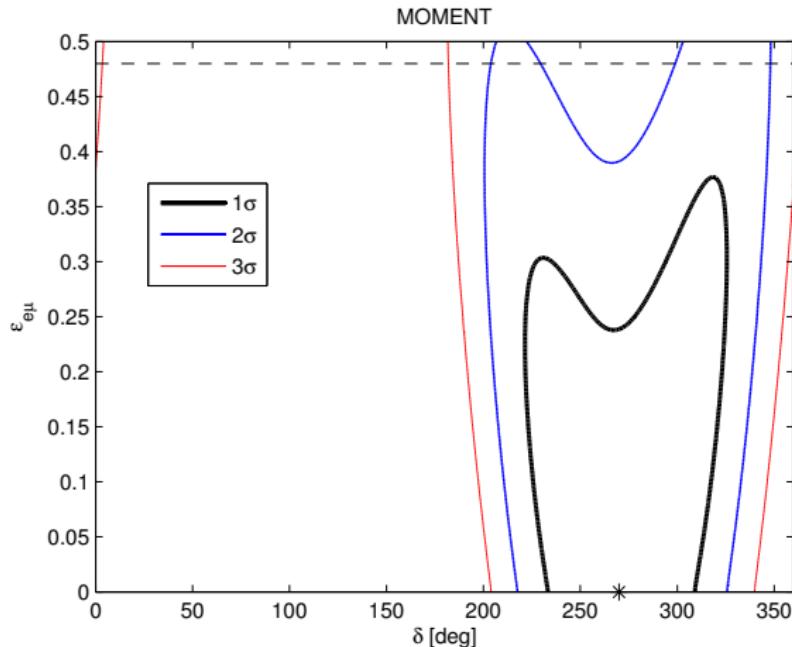
Sensitivity to ϵ_{ee} and δ_{CP}



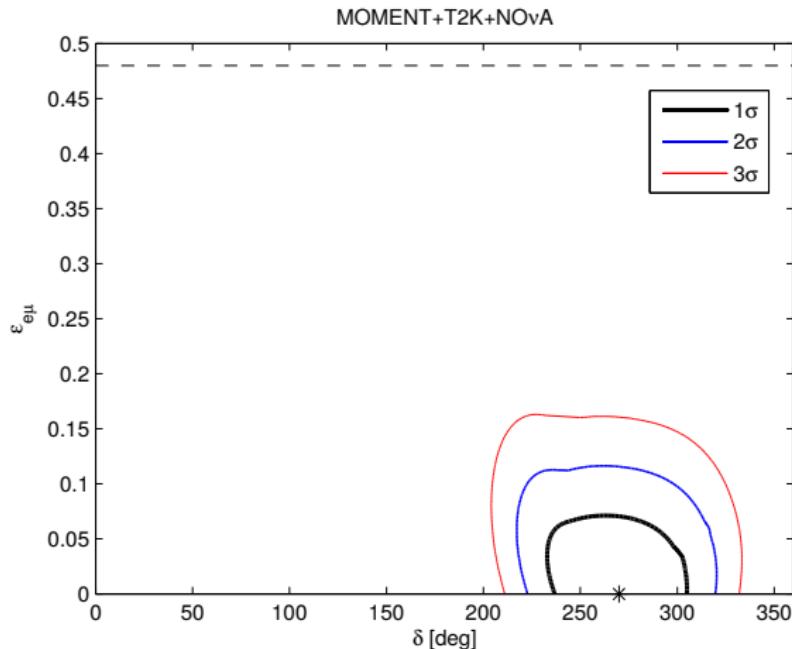
Sensitivity to $\epsilon_{e\mu}$ and δ_{CP}



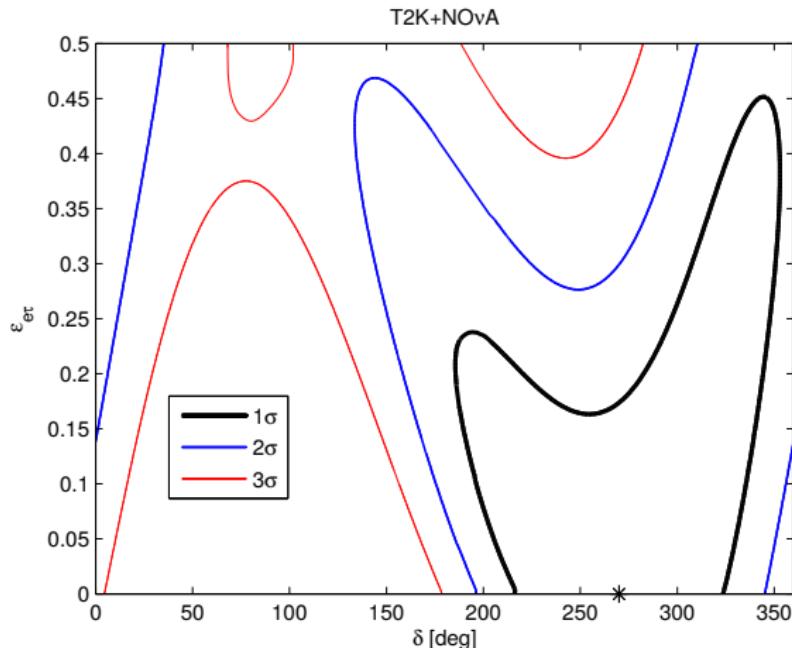
Sensitivity to $\epsilon_{e\mu}$ and δ_{CP}



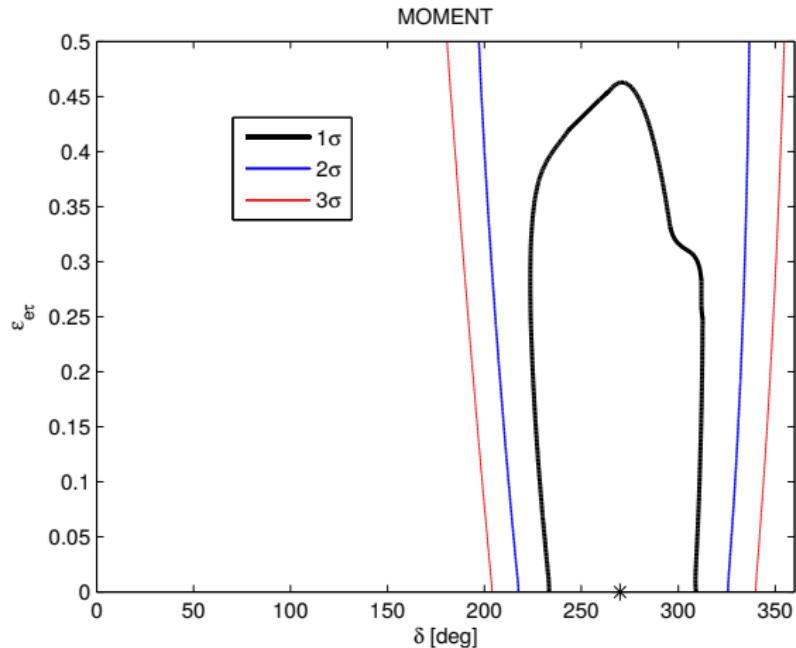
Sensitivity to $\epsilon_{e\mu}$ and δ_{CP}



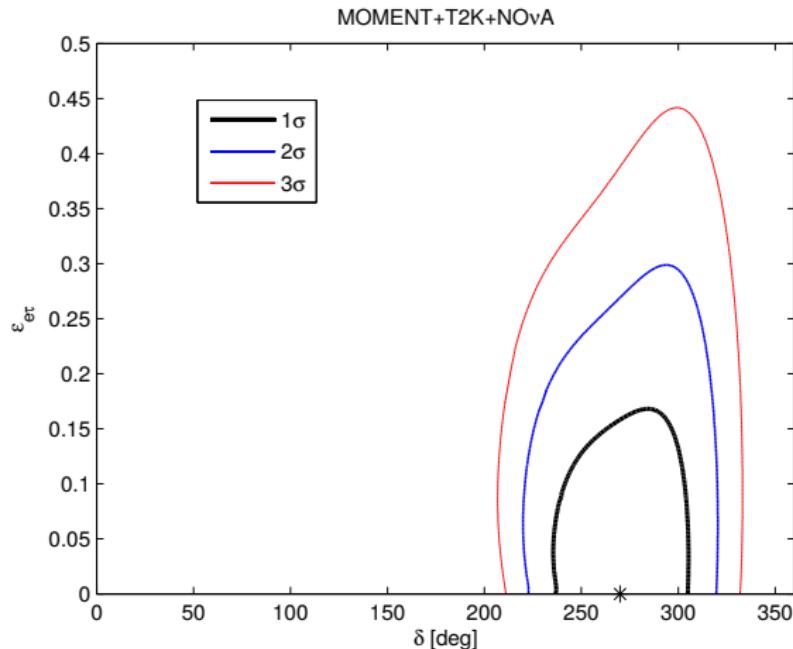
Sensitivity to $\epsilon_{e\tau}$ and δ_{CP}



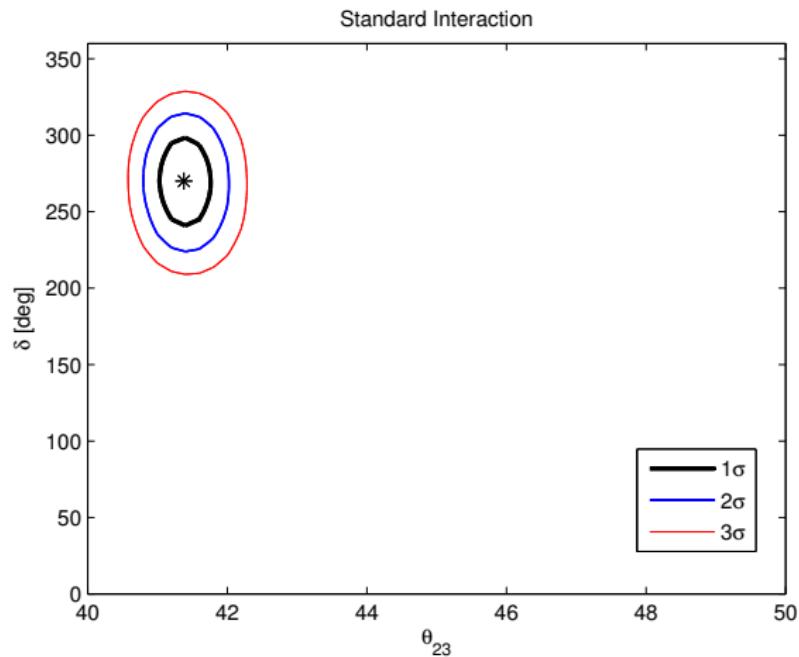
Sensitivity to $\epsilon_{e\tau}$ and δ_{CP}



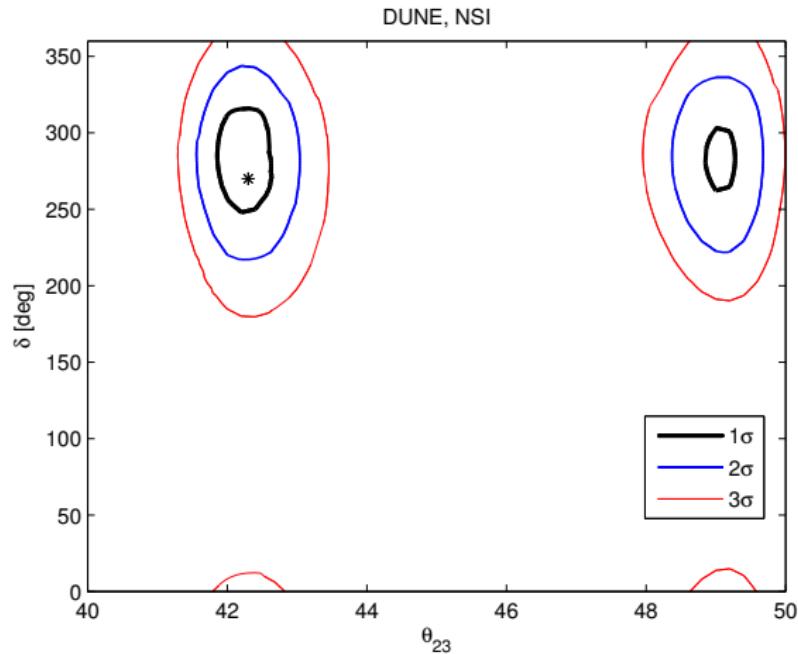
Sensitivity to $\epsilon_{e\tau}$ and δ_{CP}



Determination of θ_{23} octant and δ_{CP} by DUNE

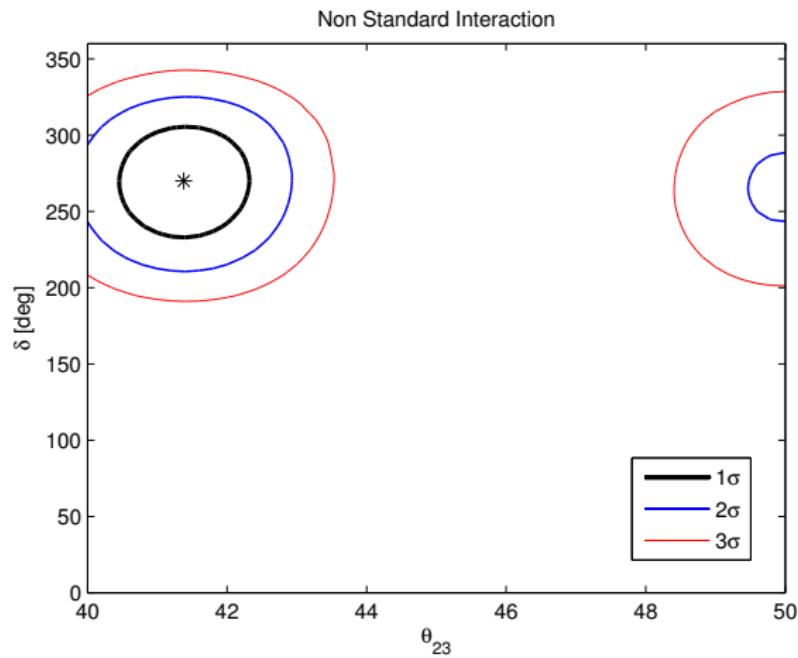


Determination of θ_{23} octant and δ_{CP} by DUNE

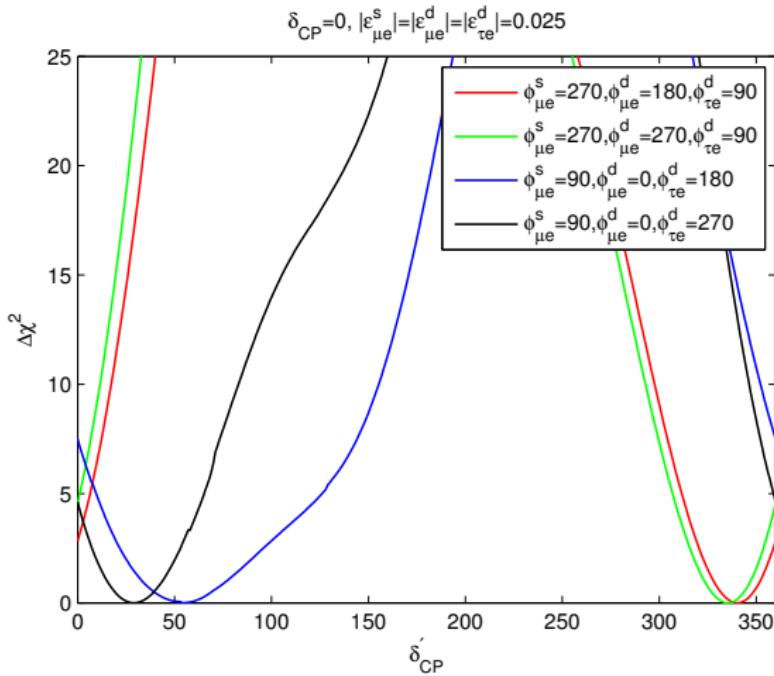


NC NSI

Determination of θ_{23} octant and δ_{CP} by DUNE



Inducing fake CP phase by CC NSI



Constraints on CC NSI by near and far detectors of DUNE

| parameter | Far Detector | Near detector | Current constraints |
|---------------------------|--------------|---------------|---------------------|
| $ \epsilon_{\mu e}^s $ | 0.016 | 0.013 | 0.026 |
| $ \epsilon_{\mu \mu}^s $ | 0.028 | 0.007 | 0.078 |
| $ \epsilon_{\mu \tau}^s $ | 0.074 | - | 0.013 |
| $ \epsilon_{ee}^d $ | 0.057 | 0.005 | 0.041 |
| $ \epsilon_{e \mu}^d $ | 0.049 | - | 0.026 |
| $ \epsilon_{\mu e}^d $ | 0.019 | 0.013 | 0.025 |
| $ \epsilon_{\mu \mu}^d $ | 0.032 | 0.007 | 0.078 |
| $ \epsilon_{\tau e}^d $ | 0.113 | - | 0.041 |
| $ \epsilon_{\tau \mu}^d $ | 0.076 | - | 0.013 |

C. Biggio, M. Blennow and E. Fernandez-Martinez, JHEP **0908** (2009) 090 [arXiv:0907.0097 [hep-ph]].

Summary

Summary

- measurement of oscillation parameters can be affected by NSI
- Effect of NC NSI can be studied by combining a low energy and a high energy neutrino experiment
- CC NSI can induce large amount of fake CP phase in DUNE
- Near detector of DUNE is important to constrain CC NSI

Thank you for your attention.